



Prospects with nanohertz GW astronomy arrays.

Golam Shaifullah

*i do not know what it is about you that closes
and opens; only something in me understands
the voice of your eyes is deeper than all roses*
-- E. E. Cummings

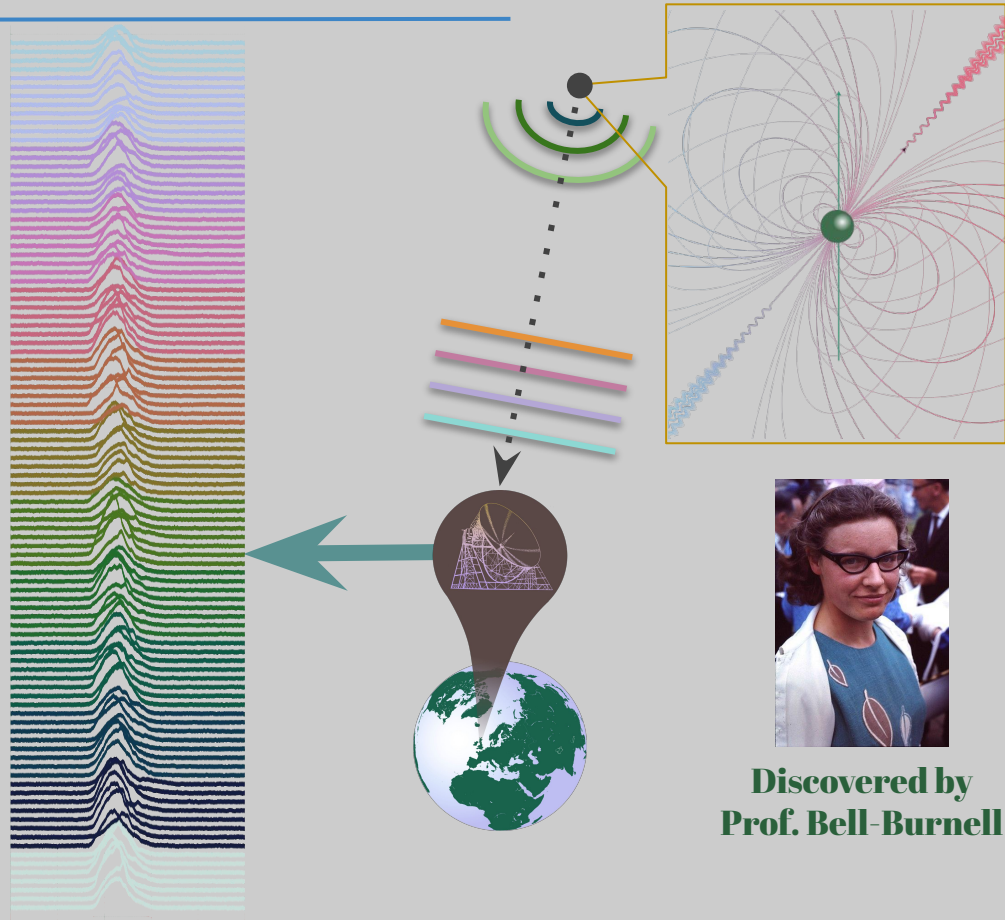


It's not a bird, it's not a plane, definitely not LGM

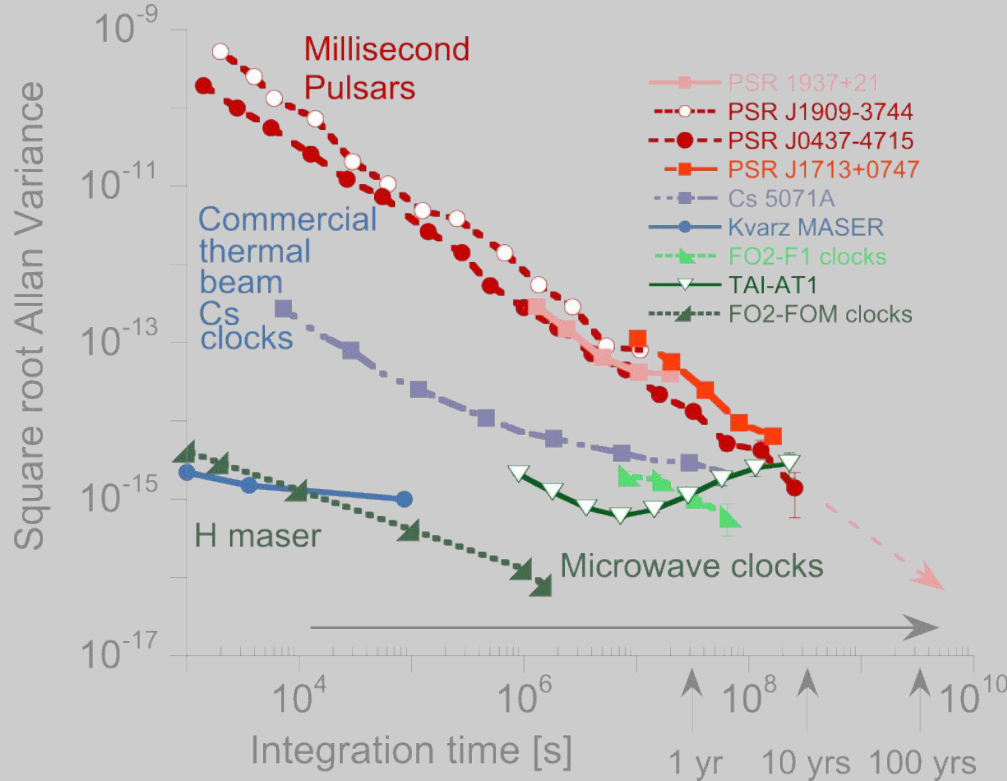
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Pulsars are giant flywheels in space, their compact masses give rise to **incredibly stable rotation**.

On each rotation, the pulsar beam produces a '*pulse*' at Earth, and the photons in that pulse can be assigned a **time-of-arrival (TOA)**.

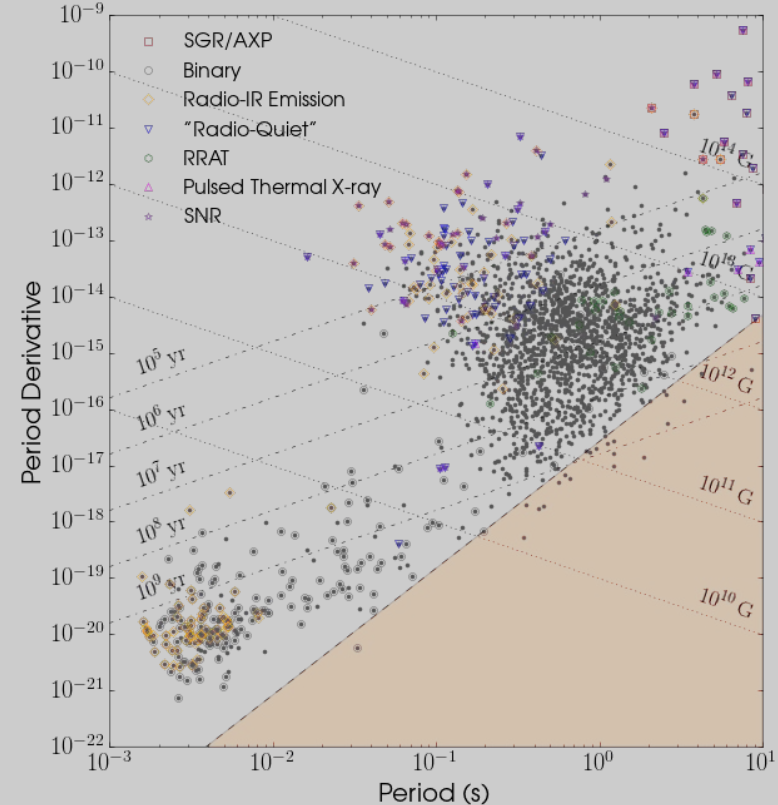


Millisecond pulsars as stable clocks



Adapted from Hartnett & Luiten, 2011

See Shannon et al (2016), Lam et al (2018) & others



psrqpy, (Pitkin, 2011)

Models, models, models

TOAs can be predicted using a model with the following (sets of) parameters:

- **astrometric**,
- **pulsar rotation** and
- **binary** (when applicable).

Apart from these pulsar emission is affected by:

- **Dispersive delays due to the intervening *ionised plasma***
- **Red noise (low frequency) processes**

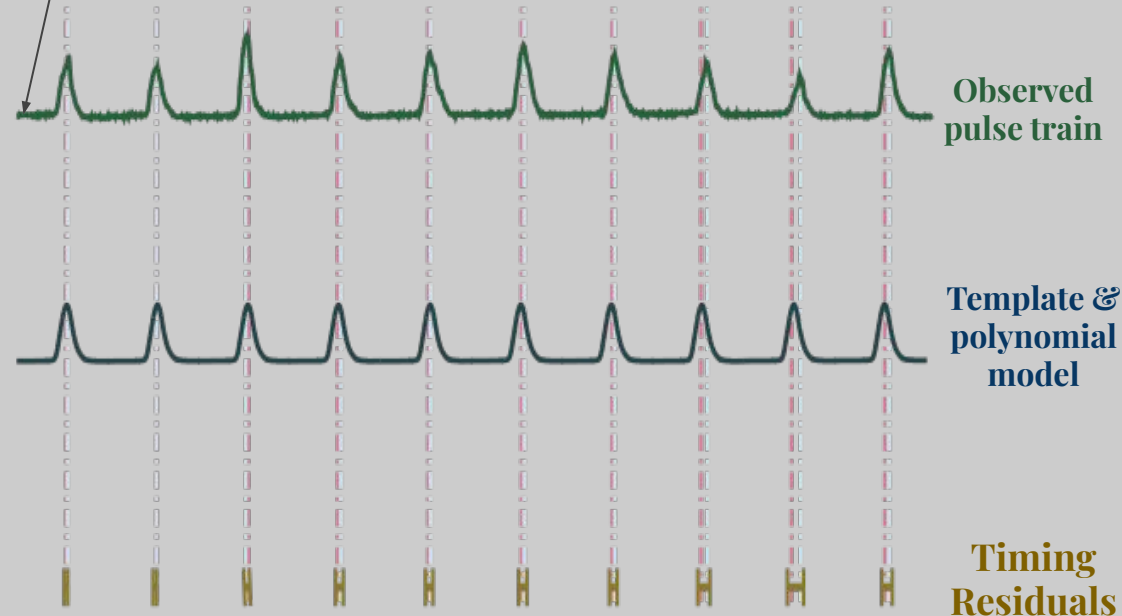


Pulsar timing

However, once we have estimates of those parameters, we can predict very precisely when the next pulse will arrive. Or the one after 20 million rotations.

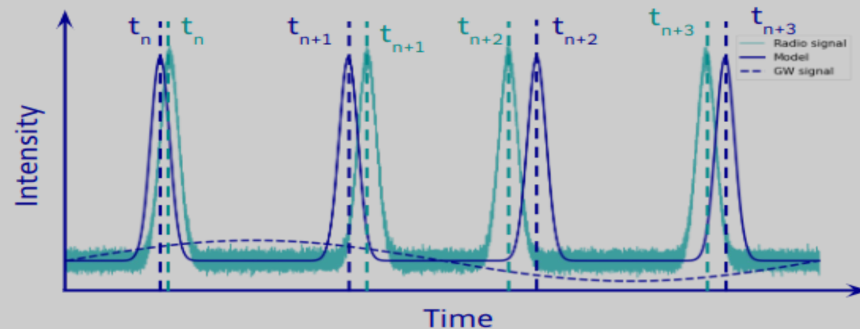
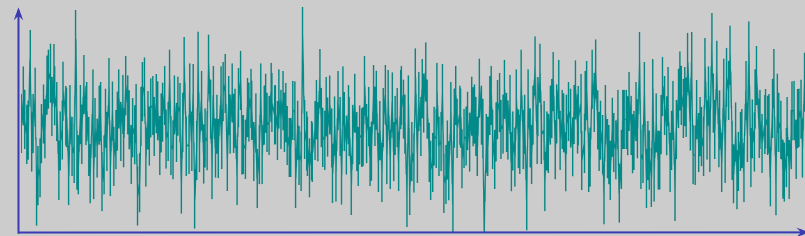
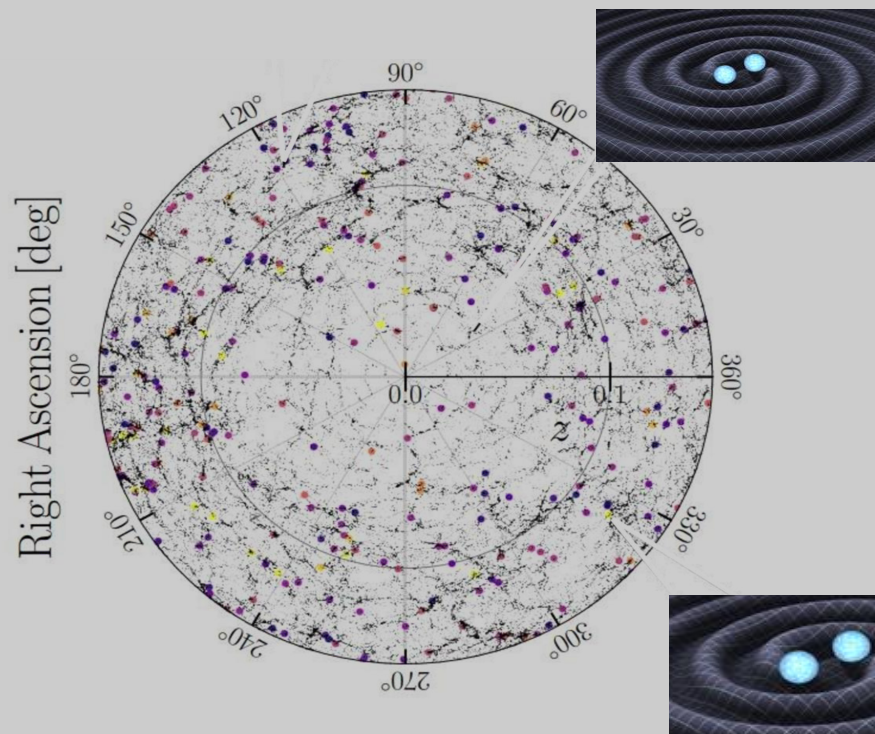
When pulses are averaged this precision quickly tends to tens of microseconds to hundreds of nanoseconds.

Time tagged to a precision of picoseconds

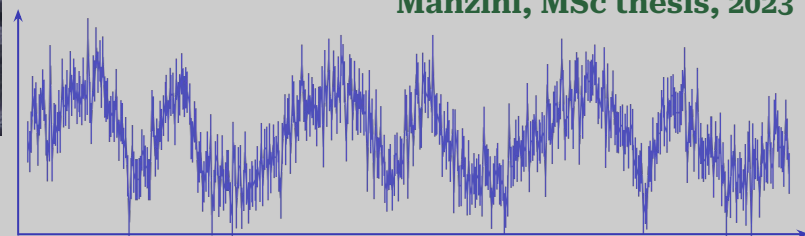


All of the light we cannot see

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Manzini, MSc thesis, 2023

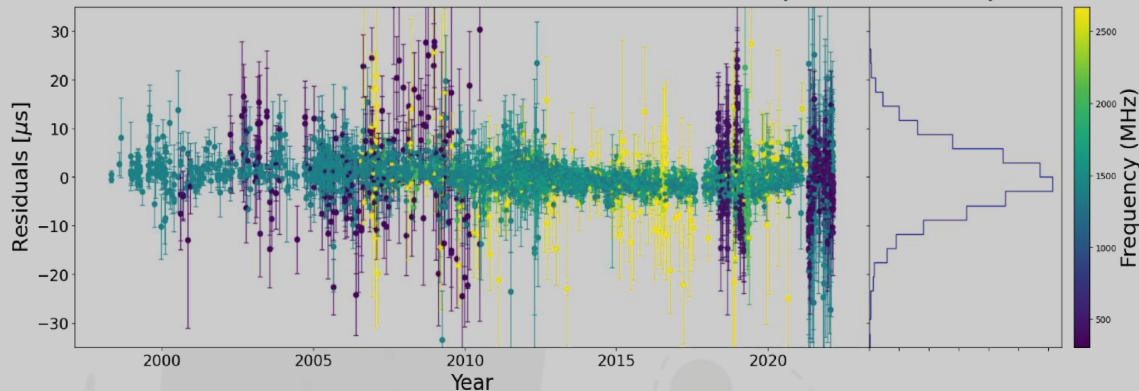


Izquierdo-Villalba et al (2024),
Curylo et al (2023), Bromm & Loeb 2003, Cole et al 2000, Benson (2012)

Pulsar timing arrays

Ferranti, MSc thesis, 2023

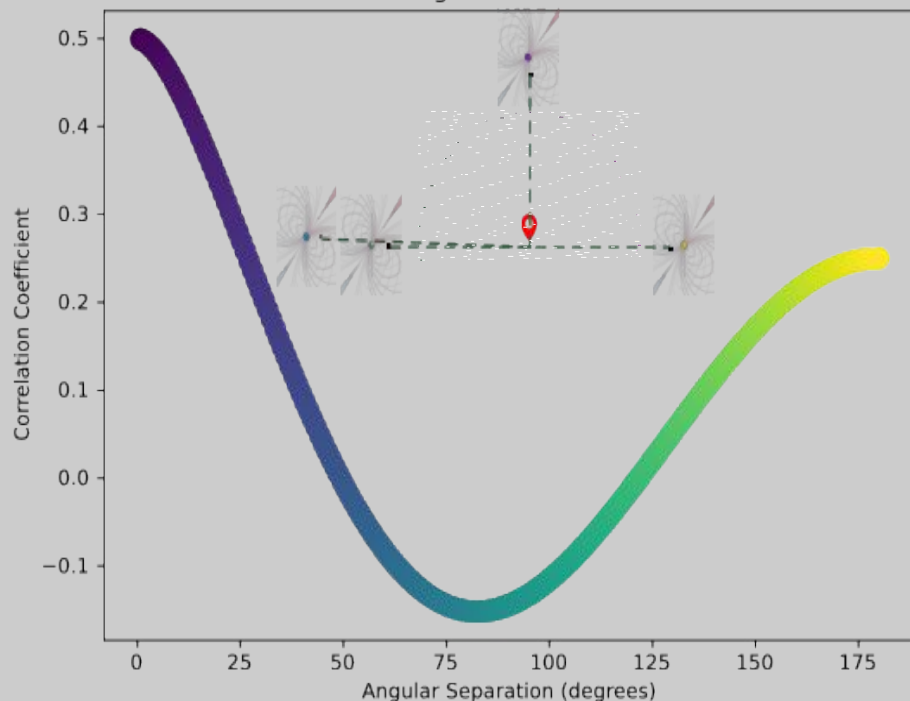
- GWs are expected to induce timing residuals on the order of a few tens of nanoseconds.
- TOA stability scales with number of rotations averaged - use **millisecond pulsars (MSPs)**!
- Single pulsars are 'jittery' and affected by noise, use an **array of MSPs**



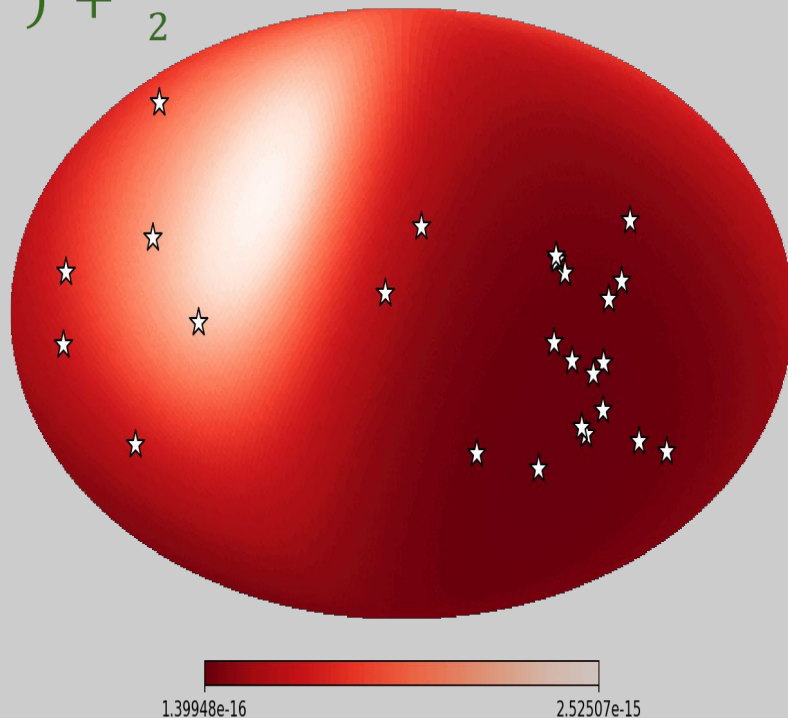
What is the signal PTAs are looking for?

$$\Gamma(\zeta) = \frac{3}{2} \left(\frac{1 - \cos \zeta}{2} \right) \ln \left(\frac{1 - \cos \zeta}{2} \right) - \frac{1}{4} \left(\frac{1 - \cos \zeta}{2} \right) + \frac{1}{2}$$

Hellings & Downs Curve



Sky Sensitivity at 1.00e-06 Hz



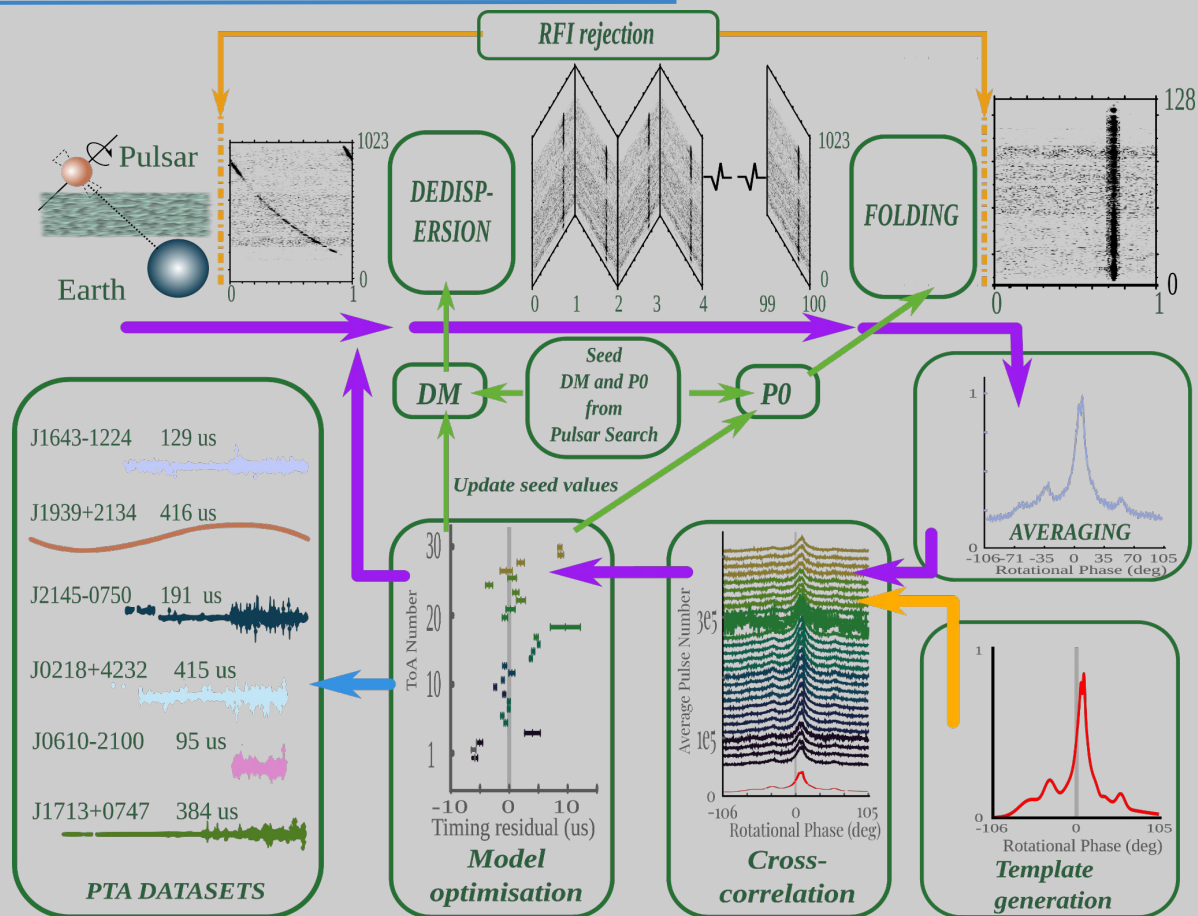
Hellings & Downs, 1983; Also see Romano & Allen, arxiv:2308.05847

FreqBayesTM pulsar timing:

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- Observe a pulsar
- De-disperse
- Stack
- Average
- Make a template
- Cross-correlate
- Line up your TOAs
- Repeat for another 20 - 100 sources
- Sprinkle post-docs for flavour
- Bake for ~30 years, turning it over once or twice a decade.

Figure from Verbiest & Shaifullah, 2018, CQG

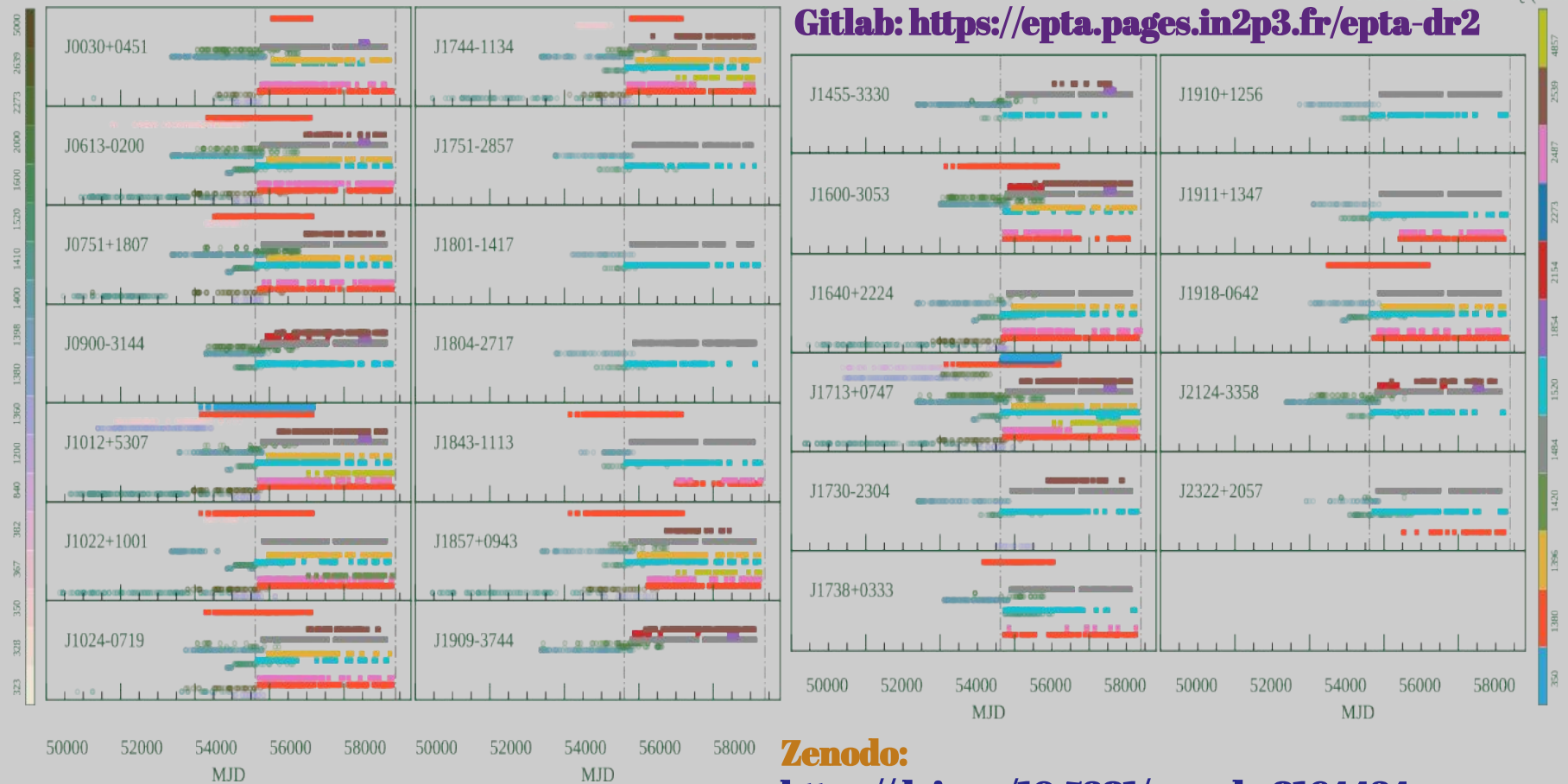


F_c (MHz) ↓ Legacy Backends

EPTA DR2 - Paper I A&A, in press, doi: 10.1051/0004-6361/202346841

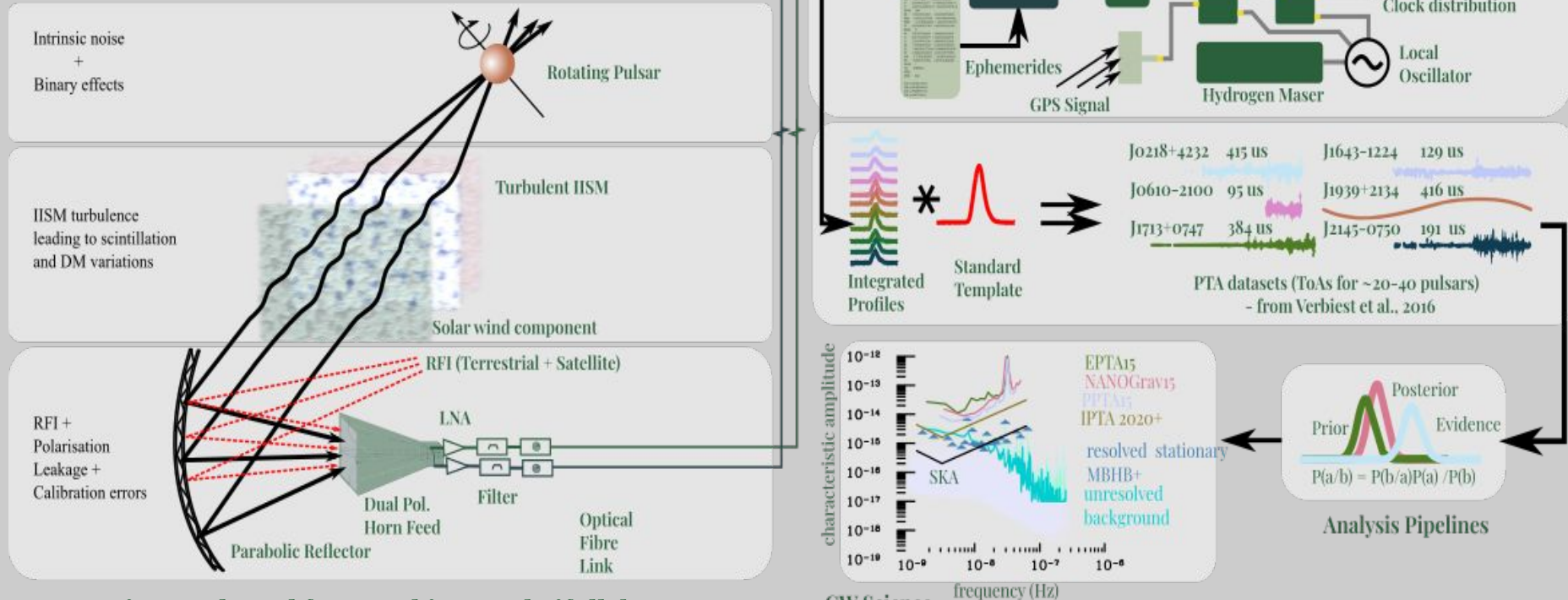
Modern Backends ↓ F_c (MHz)

Gitlab: <https://epta.pages.in2p3.fr/epta-dr2>



PTA noise sources

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• Figure adapted from Verbiest & Shaifullah, 2018, CQG

GW Science
- plot shows part of fig. 33 from Colpi & Sesana, 2017

The detection statistic and search algorithm

- We assume that noise is Gaussian: the likelihood function (likelihood of the signal with given parameters) is

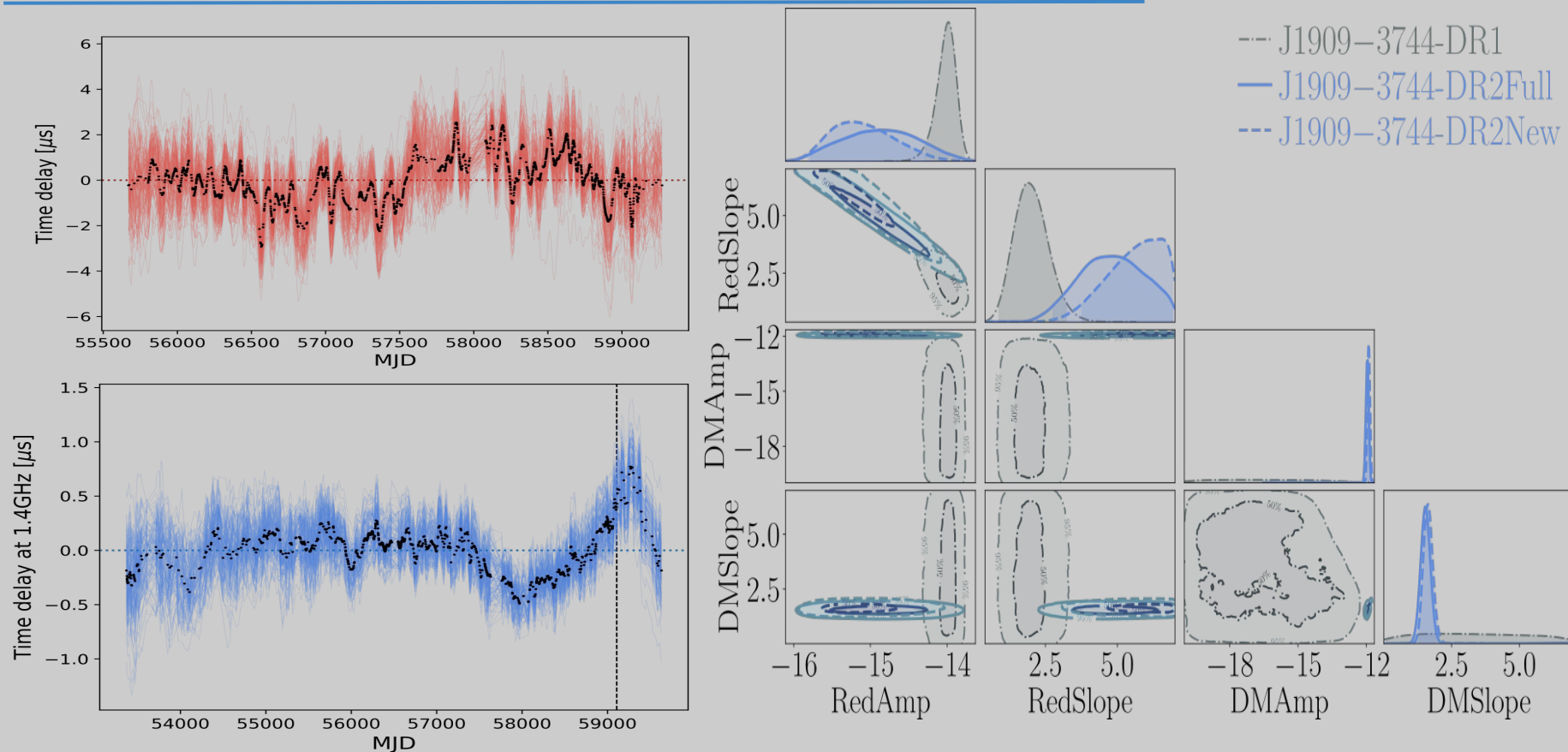
$$P(\vec{\delta t}, \theta) = \{1/\sqrt{(2\pi)^n \det(\mathbf{C})}\} \exp(-\frac{1}{2} (\vec{\delta t} - \vec{s})^T \mathbf{C}^{-1} (\vec{\delta t} - \vec{s}))$$

- δt - concatenated residuals from all pulsars in the array: total size n
- s - is a model of deterministic signals (e.g. - GW signals from individually resolvable SMBHBs)
- C is the noise variance-covariance matrix (size $n \times n$);

$$C_{\alpha i, \beta j} = C^{\text{WN}} \delta_{\alpha\beta} \delta_{ij} + C^{\text{RN}}_{ij} \delta_{\alpha\beta} + C^{\text{DM}}_{ij} \delta_{\alpha\beta} + C^{\text{GW}}_{ij} \delta_{\alpha\beta} + \dots$$

white noise
↑
red (spin) noise
↑
dispersion noise
stochastic GW noise

pulsar index
toa index

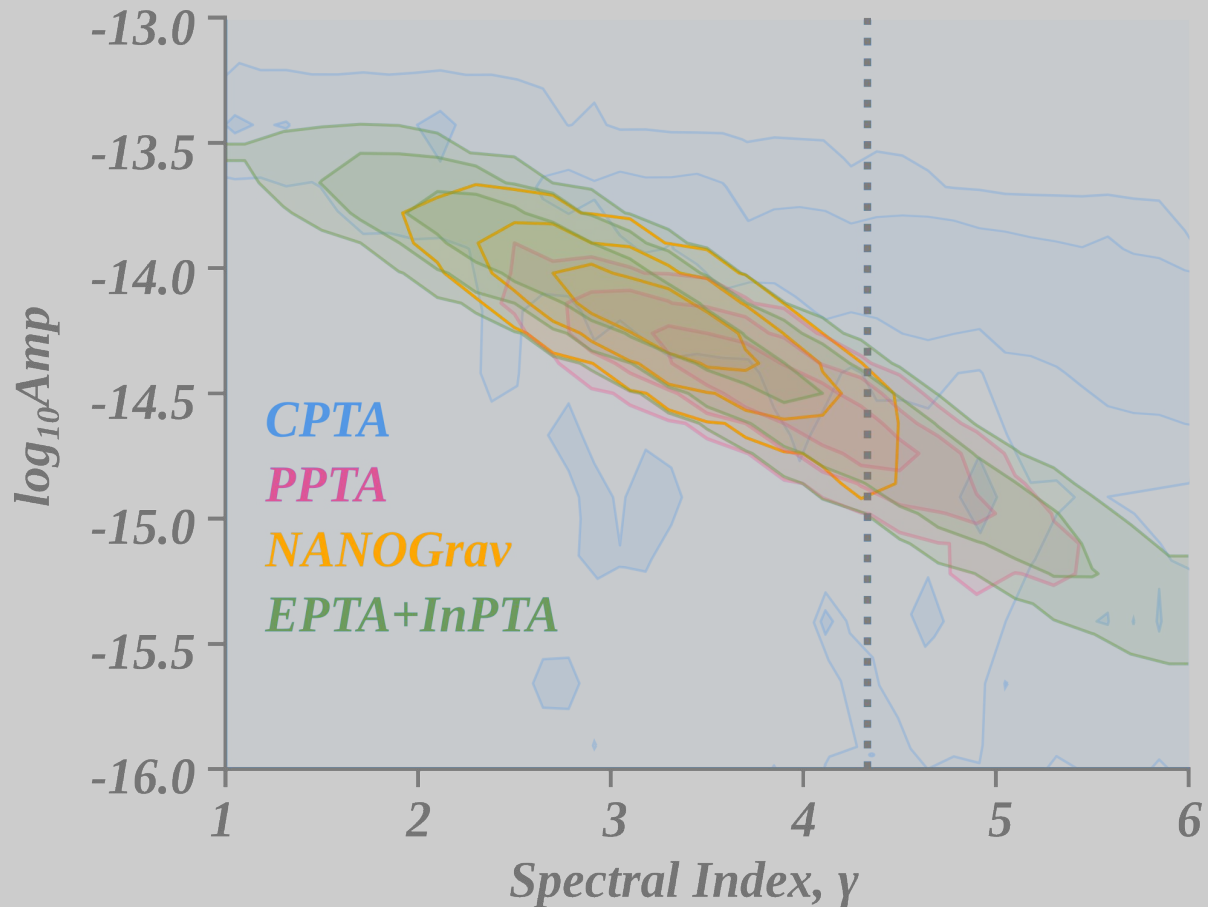


PTAs inching up to the GWB

On June 29, 2023 4 PTAs announced evidence for an HD correlated process in their data.

The significance ranges from **~ 2 to 4.6** σ ; below the **5σ** detection threshold.

Further this amplitude is **loud** ($\sim 2\text{--}3 \times 10^{-15}$) and the spectrum is **flat** (~ 3).



The International Pulsar Timing Array checklist for the detection of nanohertz gravitational waves

Bruce Allen,¹ Sanjeev Dhurandhar,² Yashwant Gupta,³ Maura McLaughlin,⁴
Priyamvada Natarajan,^{5, 6} Ryan M. Shannon,^{7, 8} Eric Thrane,^{9, 10} and Alberto Vecchio¹¹

1 Max Planck Institute for Gravitational Physics, Leibniz Universität Hannover, Callinstrasse 38, D-30167 Hannover, Germany

2 Inter University Centre for Astronomy & Astrophysics, Ganeshkhind, Pune - 411 007, India

3 National Centre for Radio Astrophysics, Pune University Campus, Pune 411007, India

4 West Virginia University Department of Physics and Astronomy, Morgantown, WV, 26501, USA

5 Department of Astronomy, 52 Hillhouse Avenue, New Haven, CT 06511

6 Black Hole Initiative, 20 Garden Street, Cambridge, MA 02138

7 Centre for Astrophysics and Supercomputing, Swinburne University of Technology, Hawthorn, VIC, 3122, Australia

8 OzGrav: The ARC Centre of Excellence for Gravitational Wave Discovery

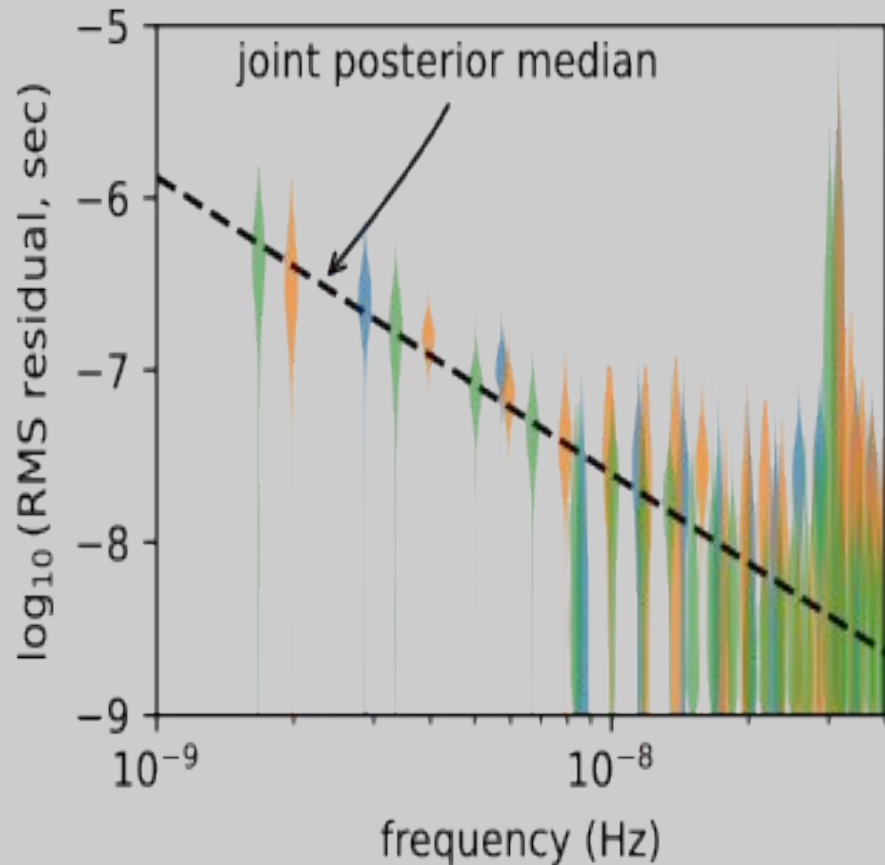
9 School of Physics and Astronomy, Monash University, Clayton VIC 3800, Australia

10 OzGrav: The ARC Centre of Excellence for Gravitational Wave Discovery, Clayton VIC 3800, Australia

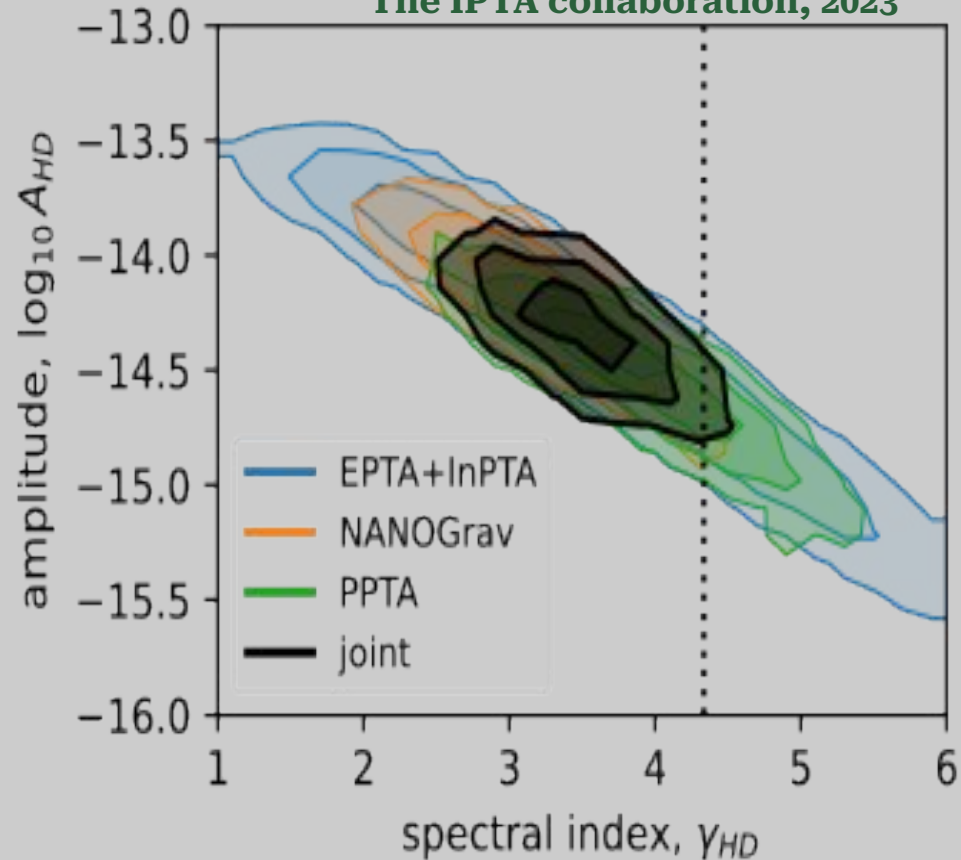
11 School of Physics and Astronomy & Institute for Gravitational Wave Astronomy, University of Birmingham, Birmingham, B15 2TT

“At the present time none of the PTAs have a detection claim.”

Do the PTAs agree?

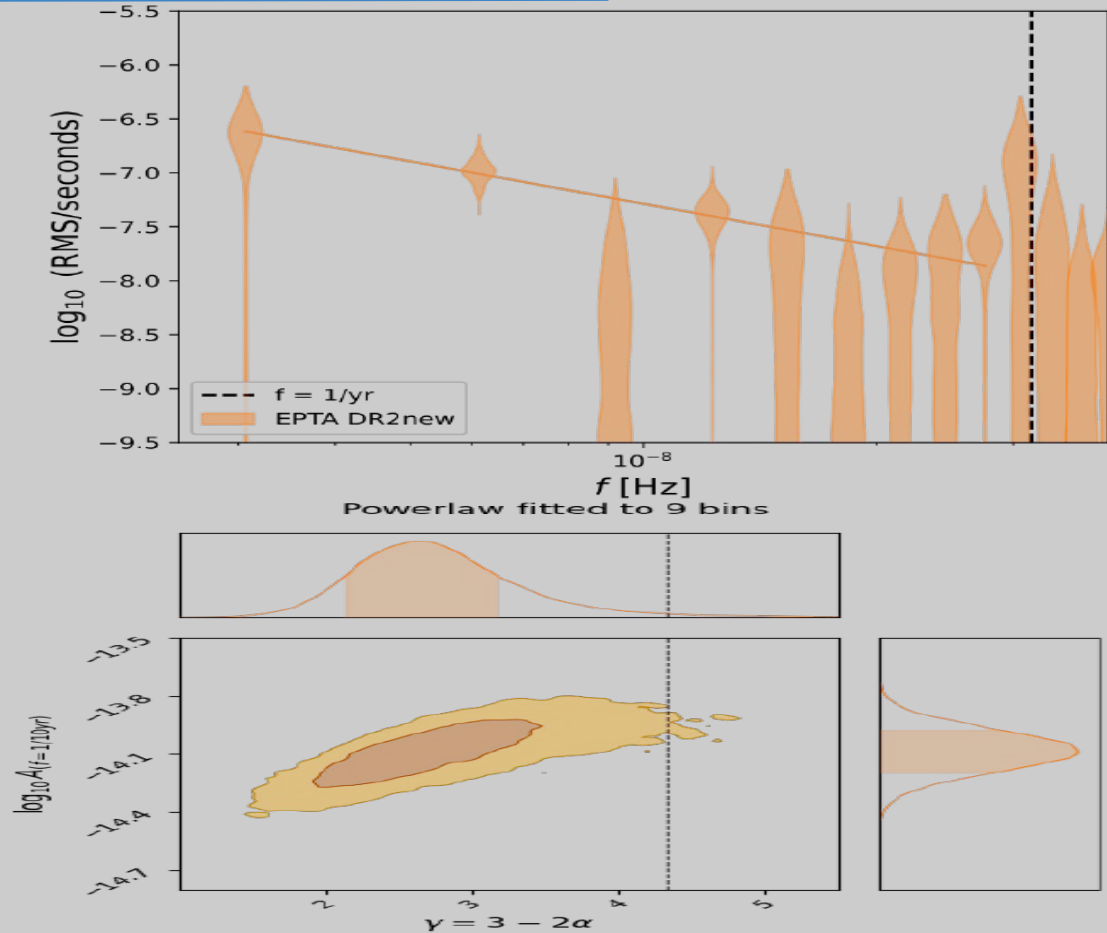


The IPTA collaboration, 2023



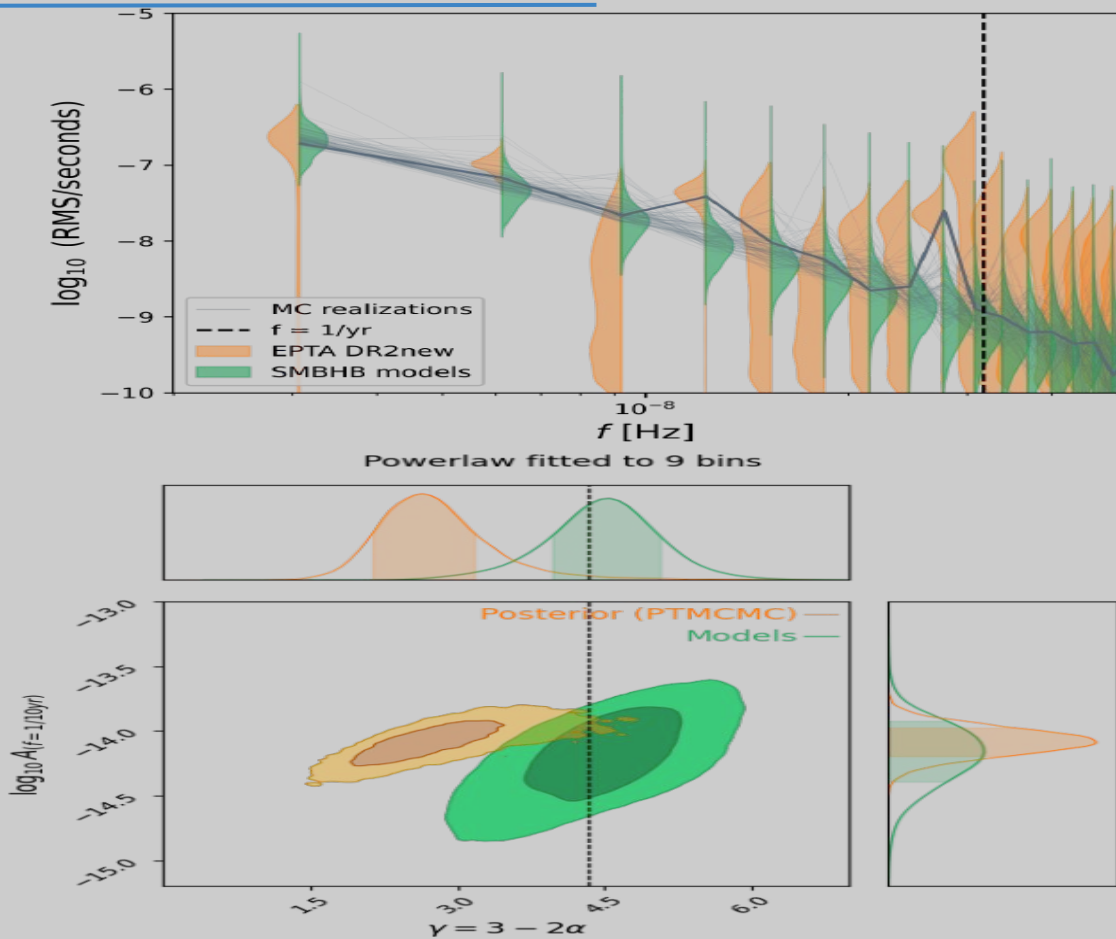
The astrophysical implications

- **The EPTA + InPTA result - a loud background**
- SMBHB generated backgrounds
- Comparisons with Semi-Analytical Models
- Stellar hardening?
- Biased by cosmic variance?
- Inflationary GWB
- Cosmic Strings
- Cosmic turbulence
- Curvature perturbations
- Challenging the ultralight dark matter paradigm



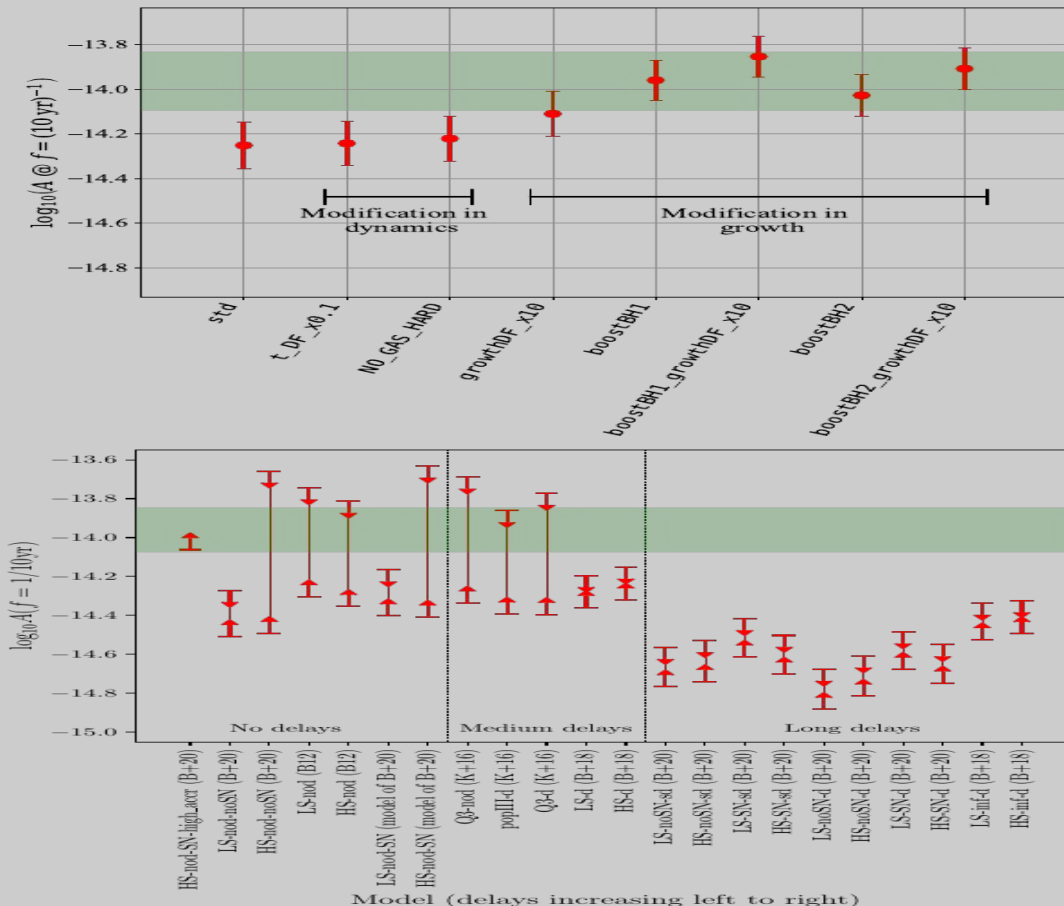
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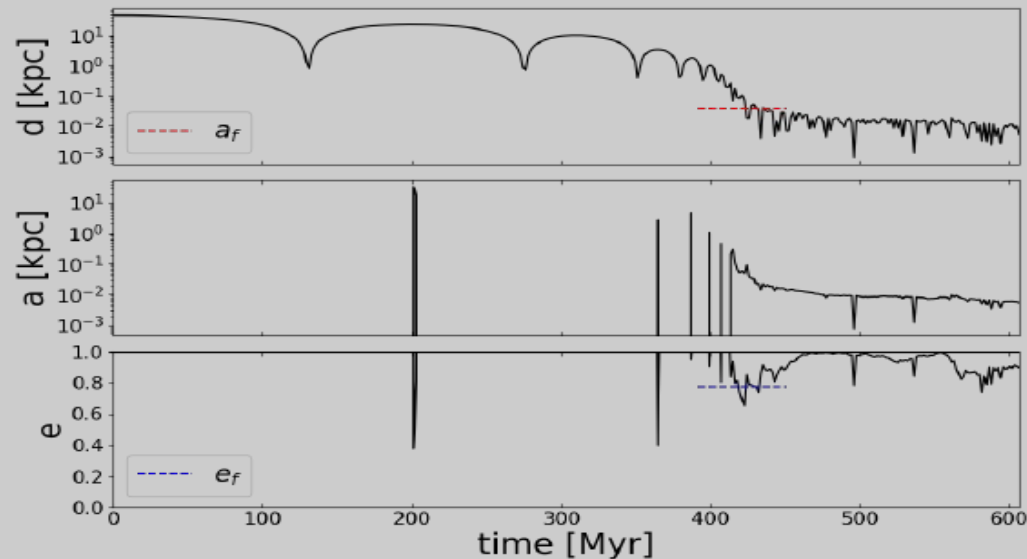
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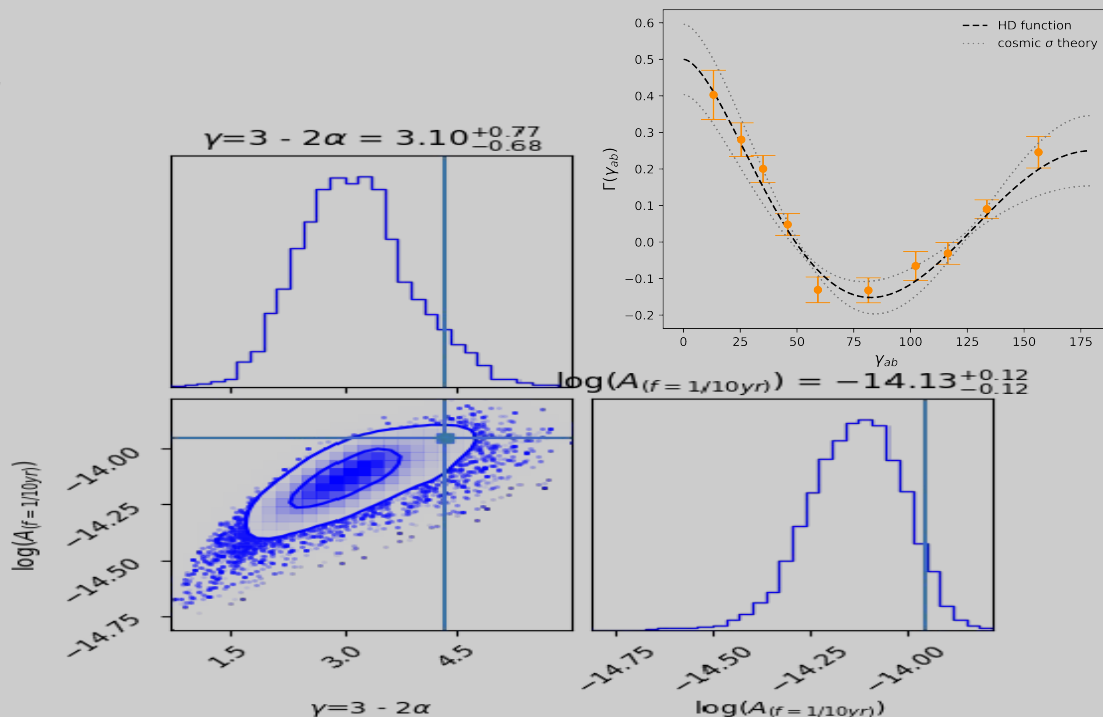
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A GWB generated by stellar hardening-affected SMBHB does NOT explain the PTA result...

The astrophysical implications

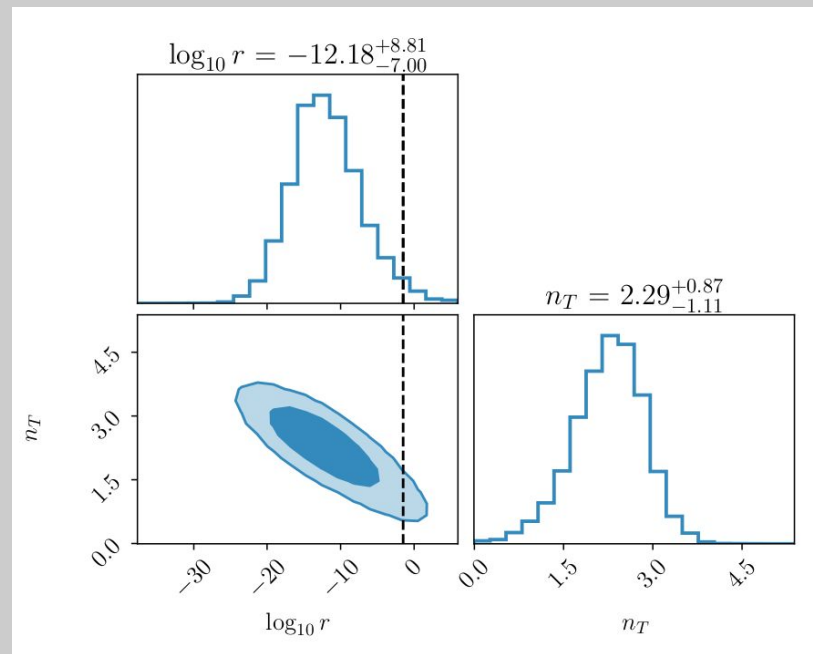
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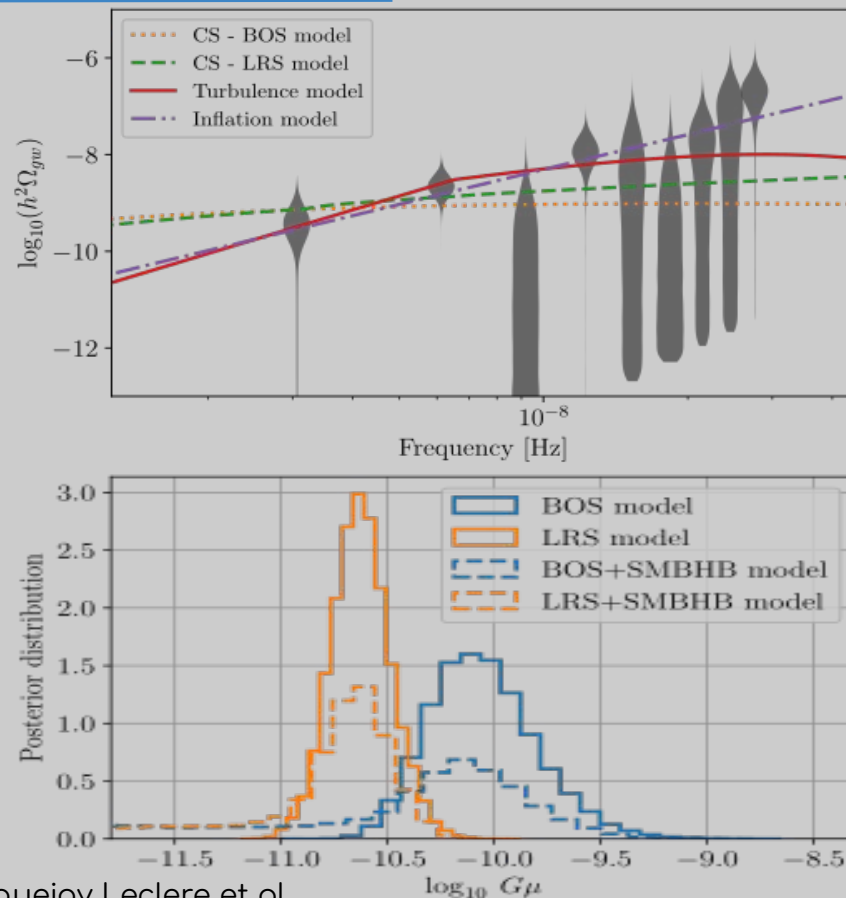
... but a biased consideration of the uncertainties of the Hellings & Downs curve might.

The cosmological implications

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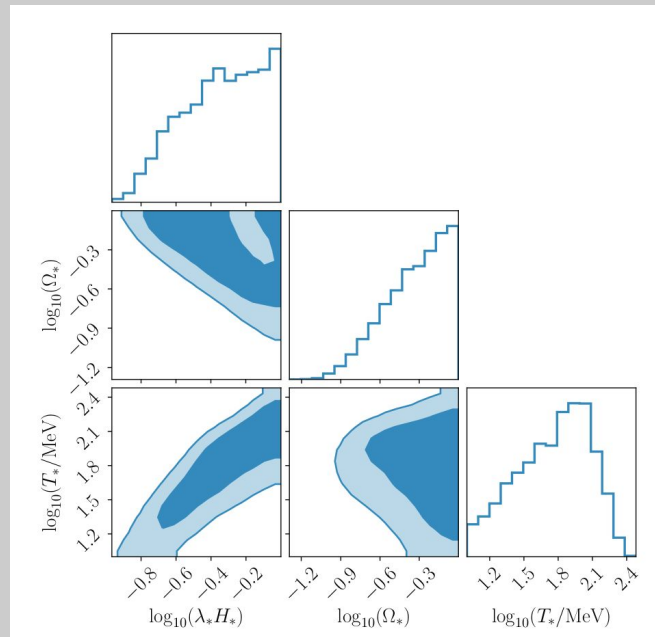


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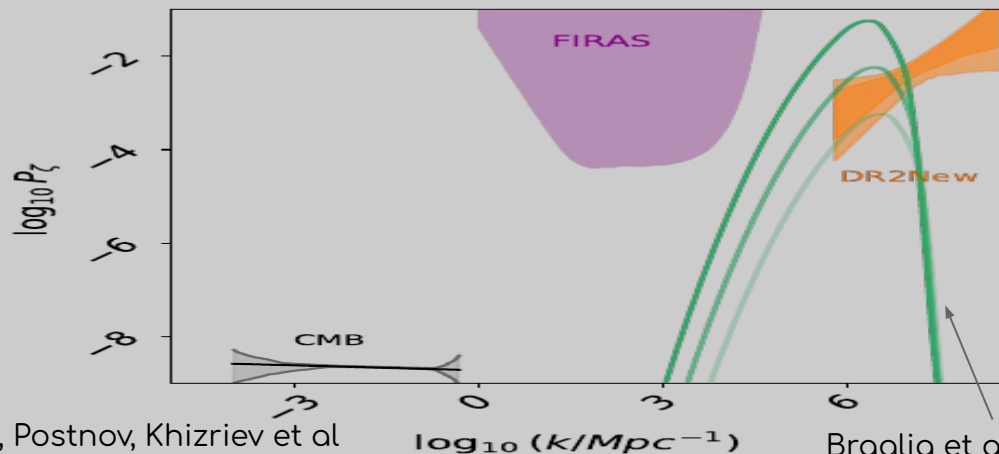
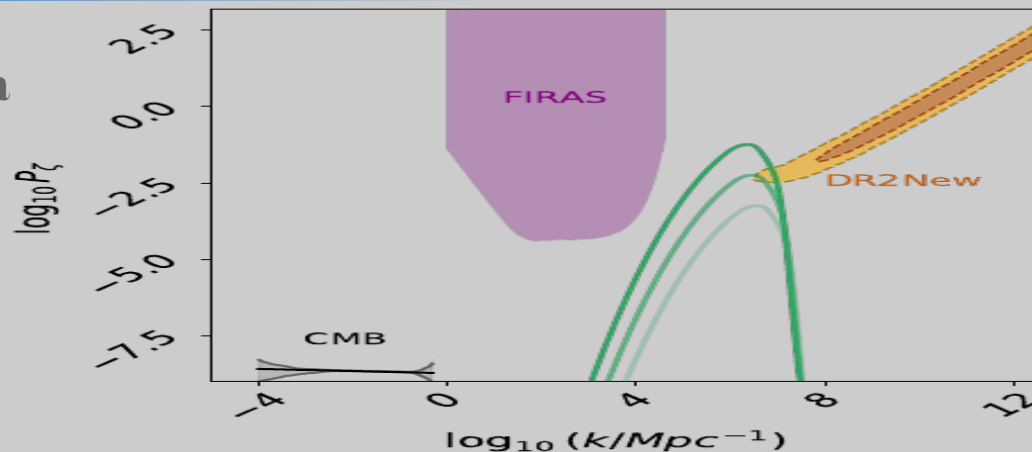


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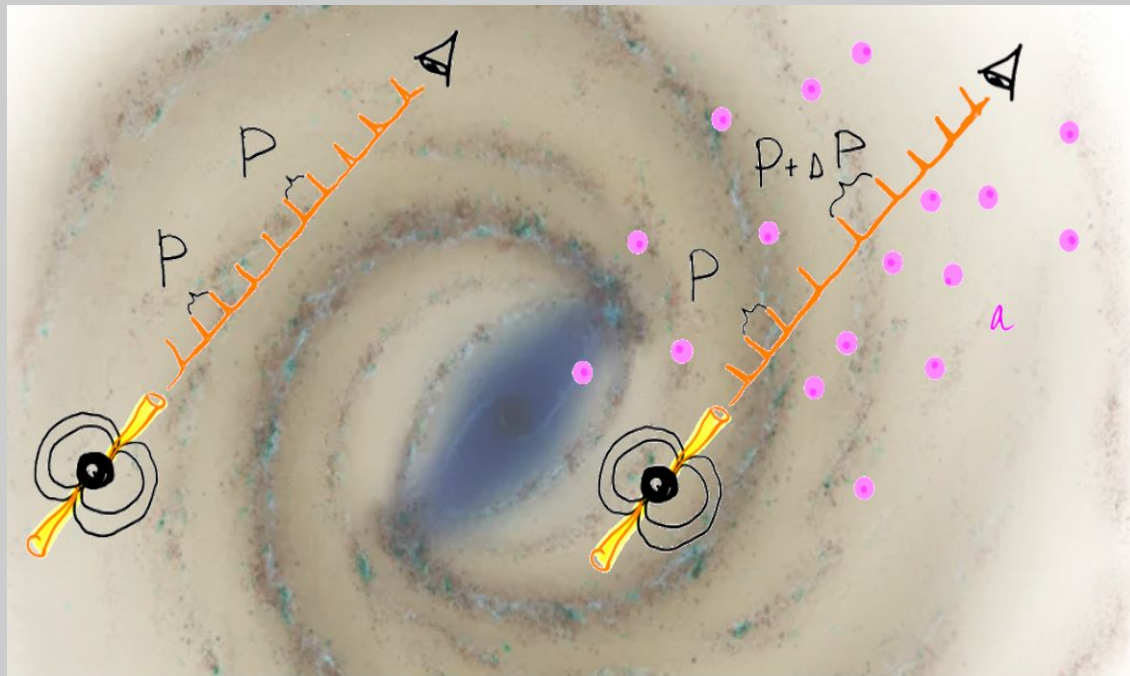


I. ULDM probes through timing data

$$R(t) = r(x_E, t_E) - r(x_p, t_p), \quad r(x_E, t_E) = \frac{\Psi(x_E)}{2\pi f} \kappa(x_E) \sin(2\pi f t_E + \alpha(x_E))$$

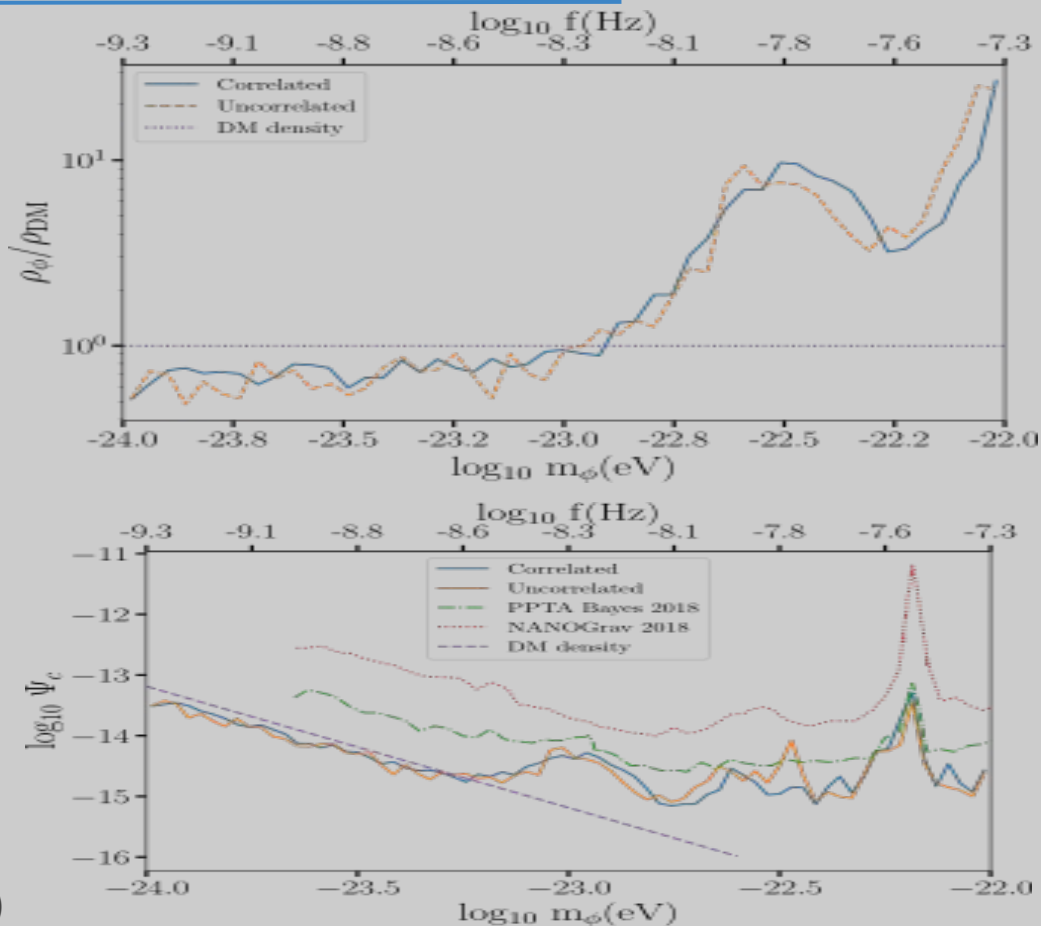
Ultra-light axion dark matter:

1. **Very light axions** with masses ranging between 10^{-23} and 10^{-20} eV
2. **Solve some of the issues of CDM** associated with overproduction of structures at galactic and sub-Galactic scales
3. **Perturb the space-time**, so that the regular flow of pulses deviate from their regular flow

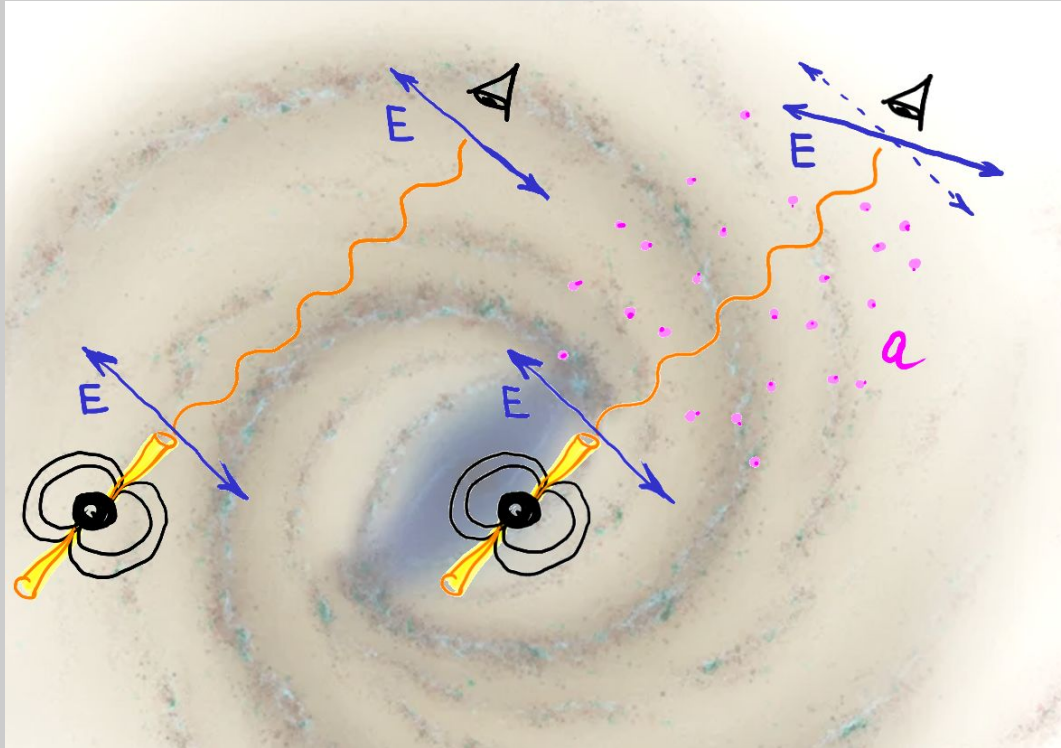


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II. ULDM with pulsar polarimetry



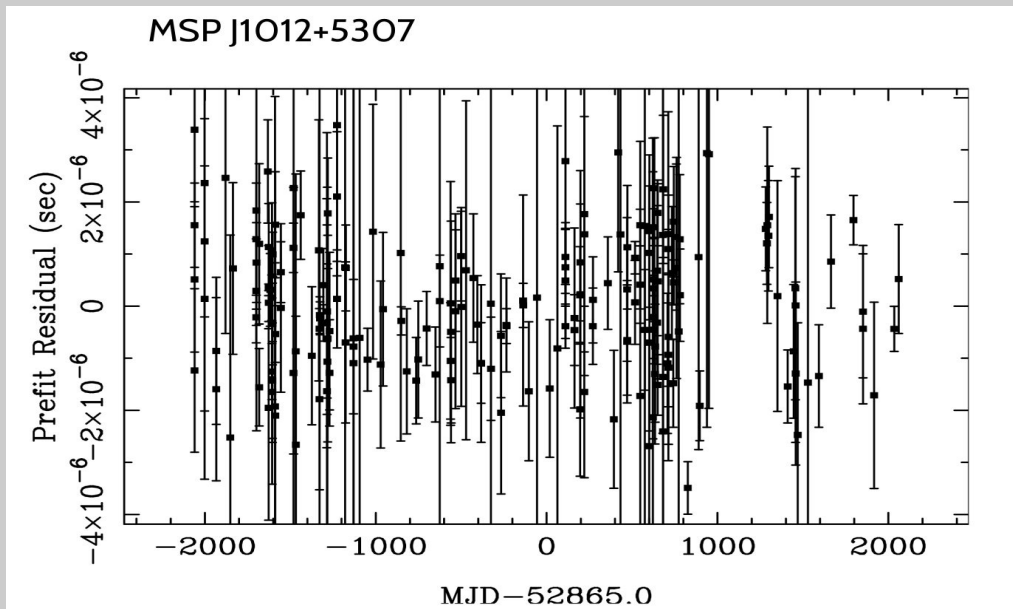
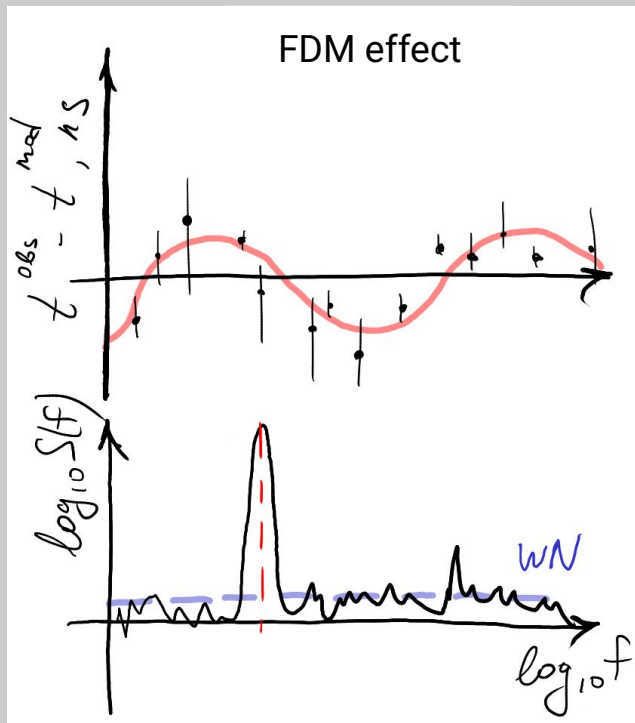
Credit:
NASA/JPL-Caltech

Ultra-light axions in the Milky Way:

1. **Very light axions** with masses ranging between 10^{-23} and 10^{-20} eV
2. When interacting weakly with photons, **rotate the plane** of linearly polarised pulsar light
3. Plane of linear polarisation **oscillates with periods of several years** due to varying pressure

I. ULDM probes through timing data

$$R(t) = r(x_E, t_E) - r(x_p, t_p), \quad r(x_E, t_E) = \frac{\Psi(x_E)}{2\pi f} \kappa(x_E) \sin(2\pi f t_E + \alpha(x_E))$$



II. ULDM with pulsar polarimetry

If we assume non-renormalizable interaction between fuzzy DM particles and photons:

$$\mathcal{L} = \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{g_{a\gamma}}{4}aF_{\mu\nu}\tilde{F}^{\mu\nu} + \frac{1}{2}(\partial_\mu a\partial^\mu a - m_a^2 a^2)$$

$$(\square + m_a^2)a + \frac{g_{a\gamma}}{4}aF_{\mu\nu}\tilde{F}^{\mu\nu} = 0$$

Polarization properties of light are altered

$$\omega_{\pm} = k\sqrt{1 \pm g_{a\gamma}\frac{\partial_0 a}{k}} \simeq k \pm \frac{1}{2}g_{a\gamma}\partial_0 a$$

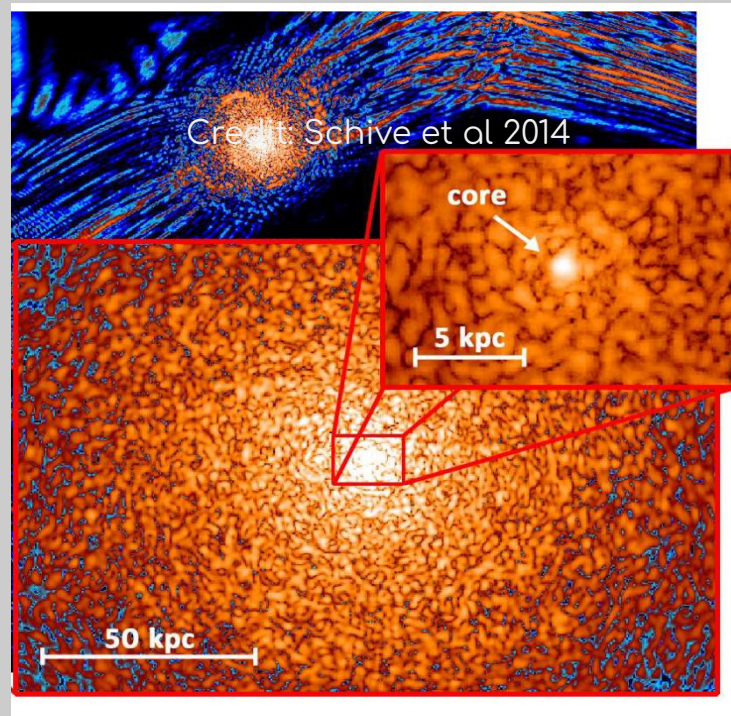
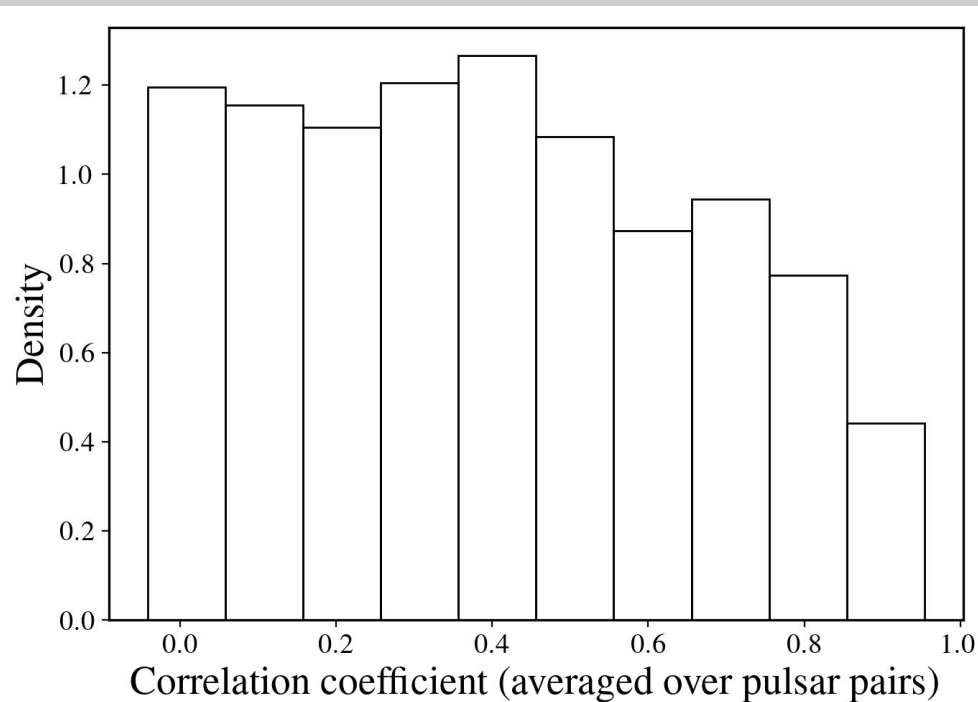
$$\Delta(\text{PA}(t)) = \frac{g_{a\gamma}}{\sqrt{2}m}[\text{p}(t_E, x_E) - \text{p}(t_p, x_p)], \quad \text{p}(t_E, x_E) = \sqrt{\rho_{\text{DM}}\kappa_E} \cos(mt + \phi(x_E))$$



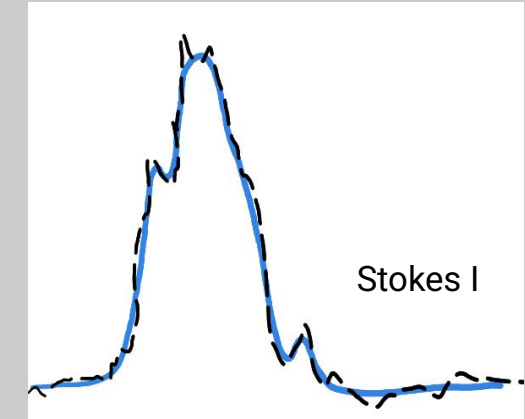
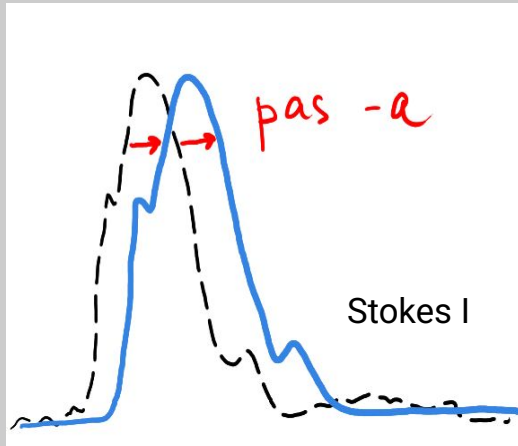
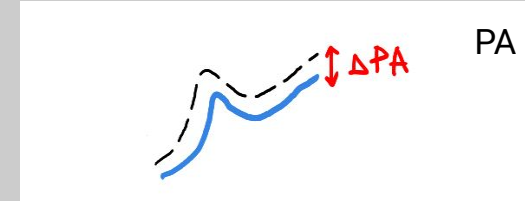
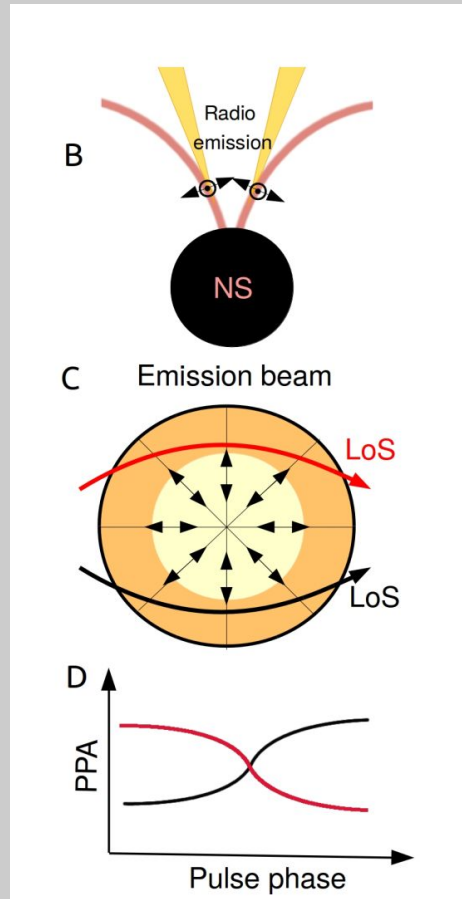
See: Ivanov et al 2018,
Castillo et al 2022

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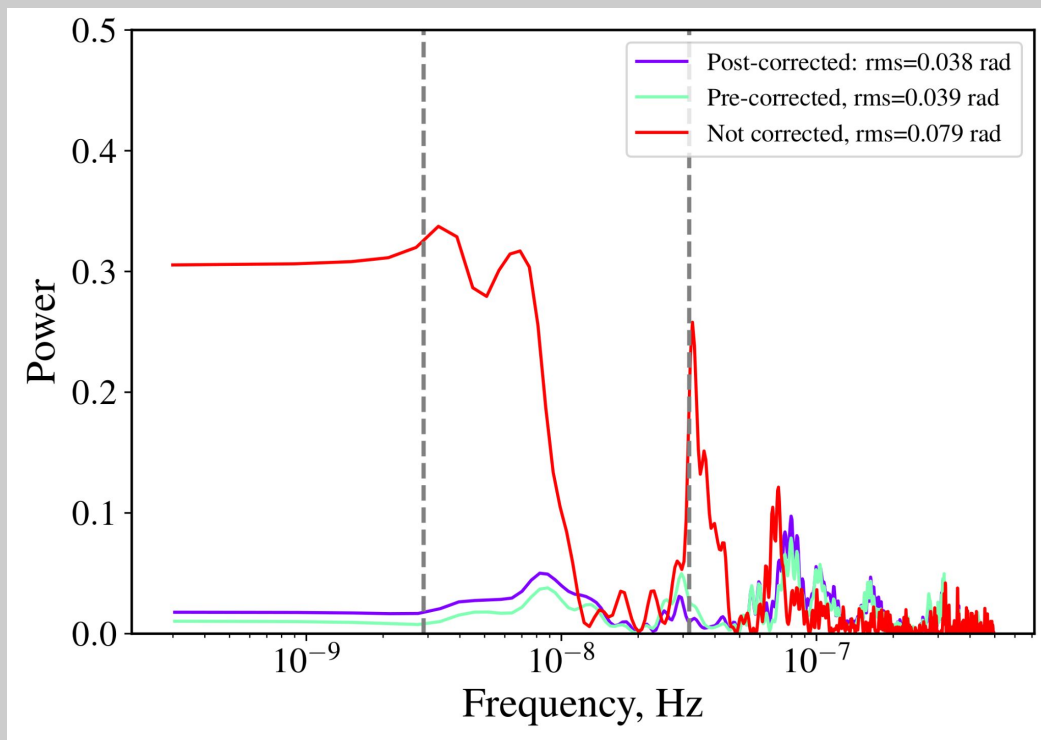


II. ULDM with pulsar polarimetry: data processing



$$\Delta PA = \frac{1}{2} \arcsin \{ (U^{\text{obs}} Q^{\text{tmpl}} - Q^{\text{obs}} U^{\text{tmpl}}) \}$$

II. ULDM with pulsar polarimetry: systematics

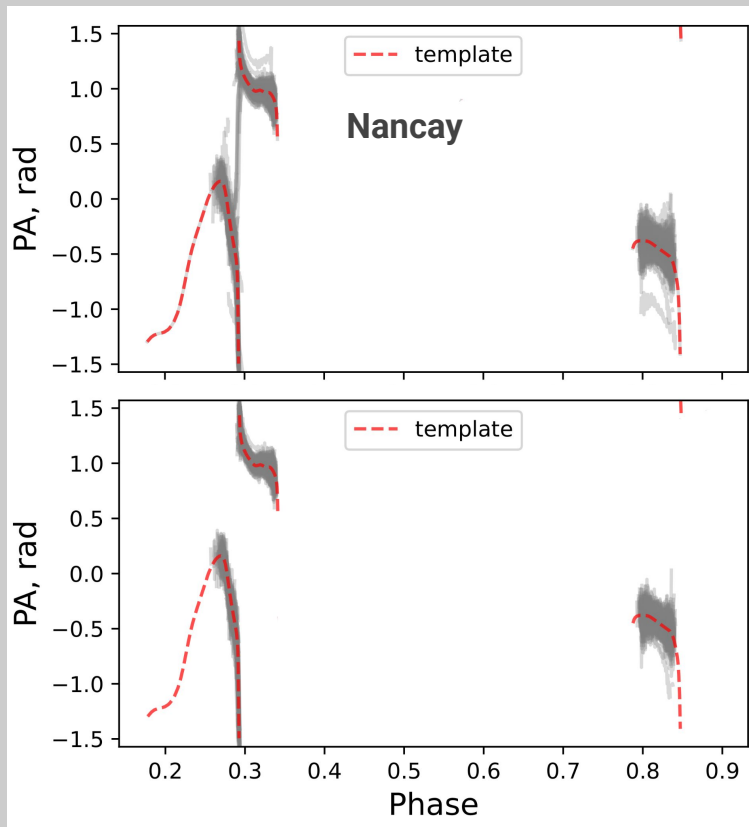


RMextract from Maaijke Mevius:
<https://github.com/lofar-astron/RMextract/tree/master/RMextract>

- i) Ionospheric TEC maps (uqrg)
- +
- ii) Geomagnetic field model (WMM)
- +
- iii) Thin screen approximation

$$\text{RM}_{\text{iono}} = \int n_e \mathbf{B}_{\text{LOS}} d\mathbf{r}$$
$$\text{RM}_{\text{iono}} \sim \text{STEC} \times \mathbf{B}_{\text{IPP}}$$

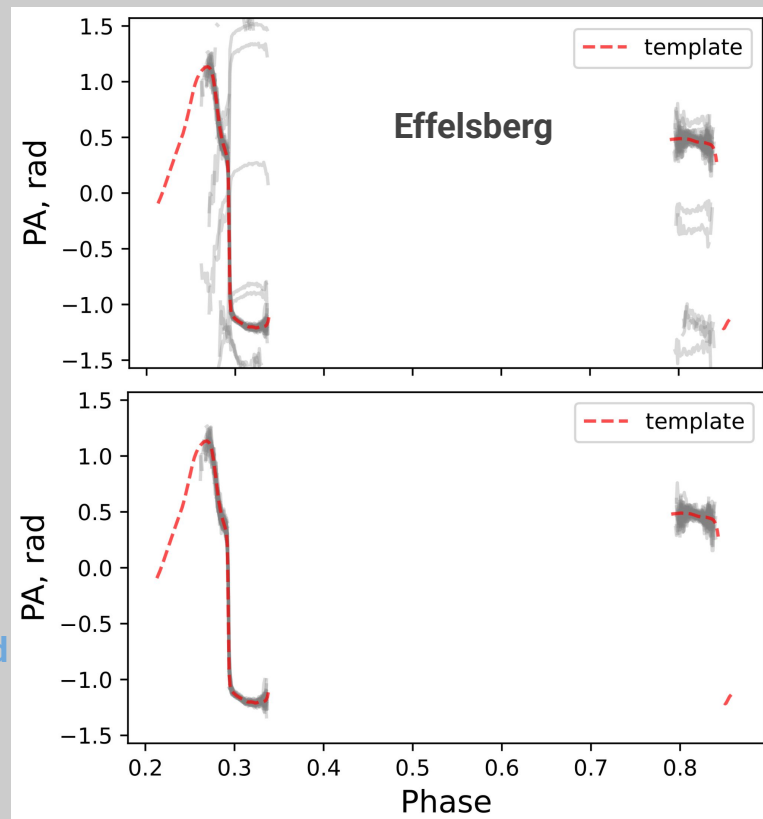
II. ULDM with pulsar polarimetry: challenges



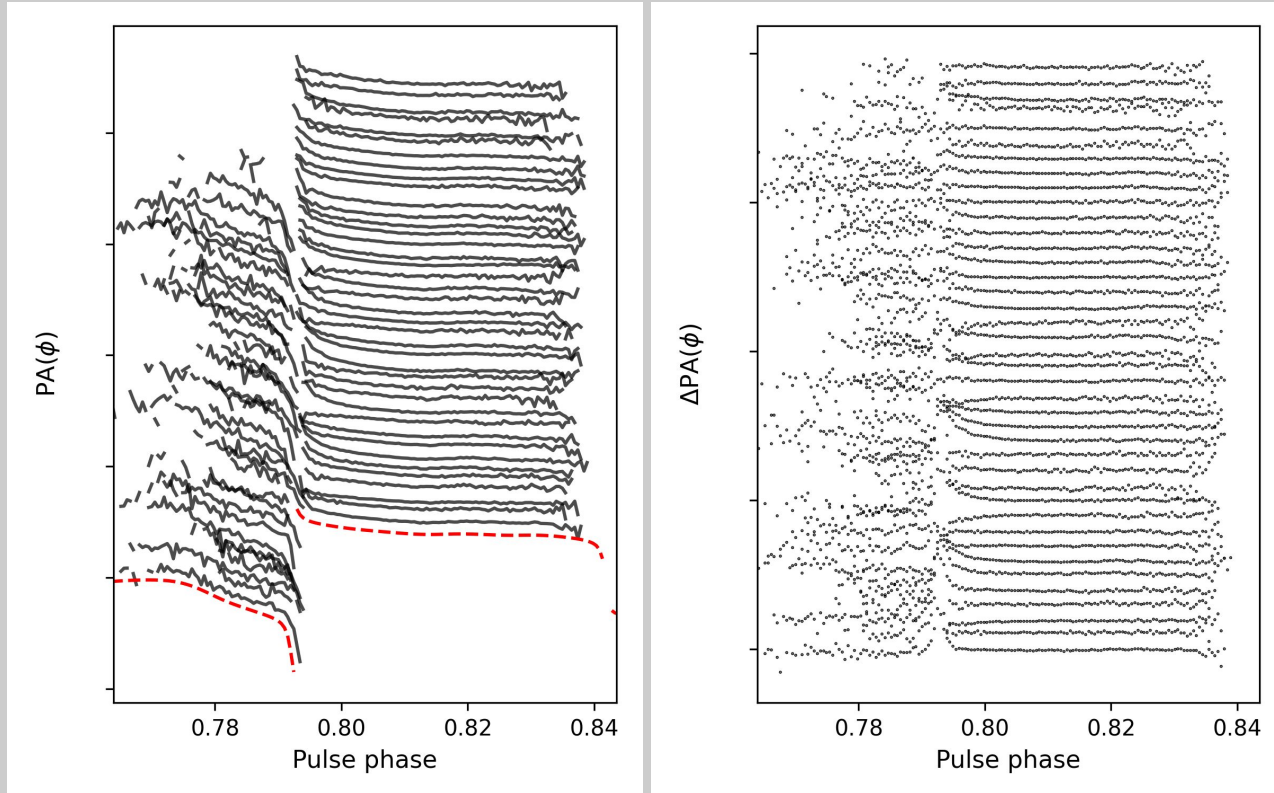
B1937+21

All PA profiles

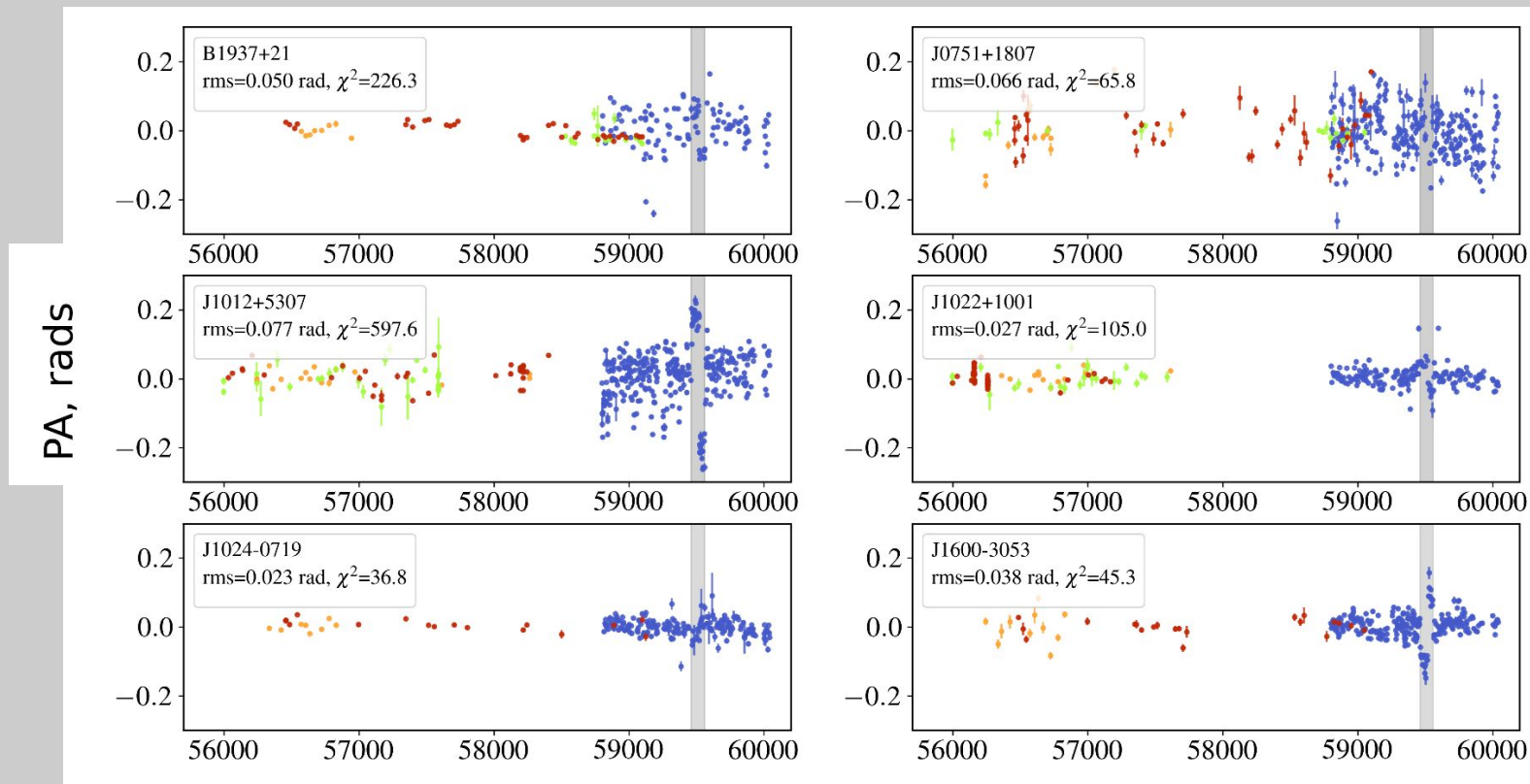
Manually selected
PA profiles



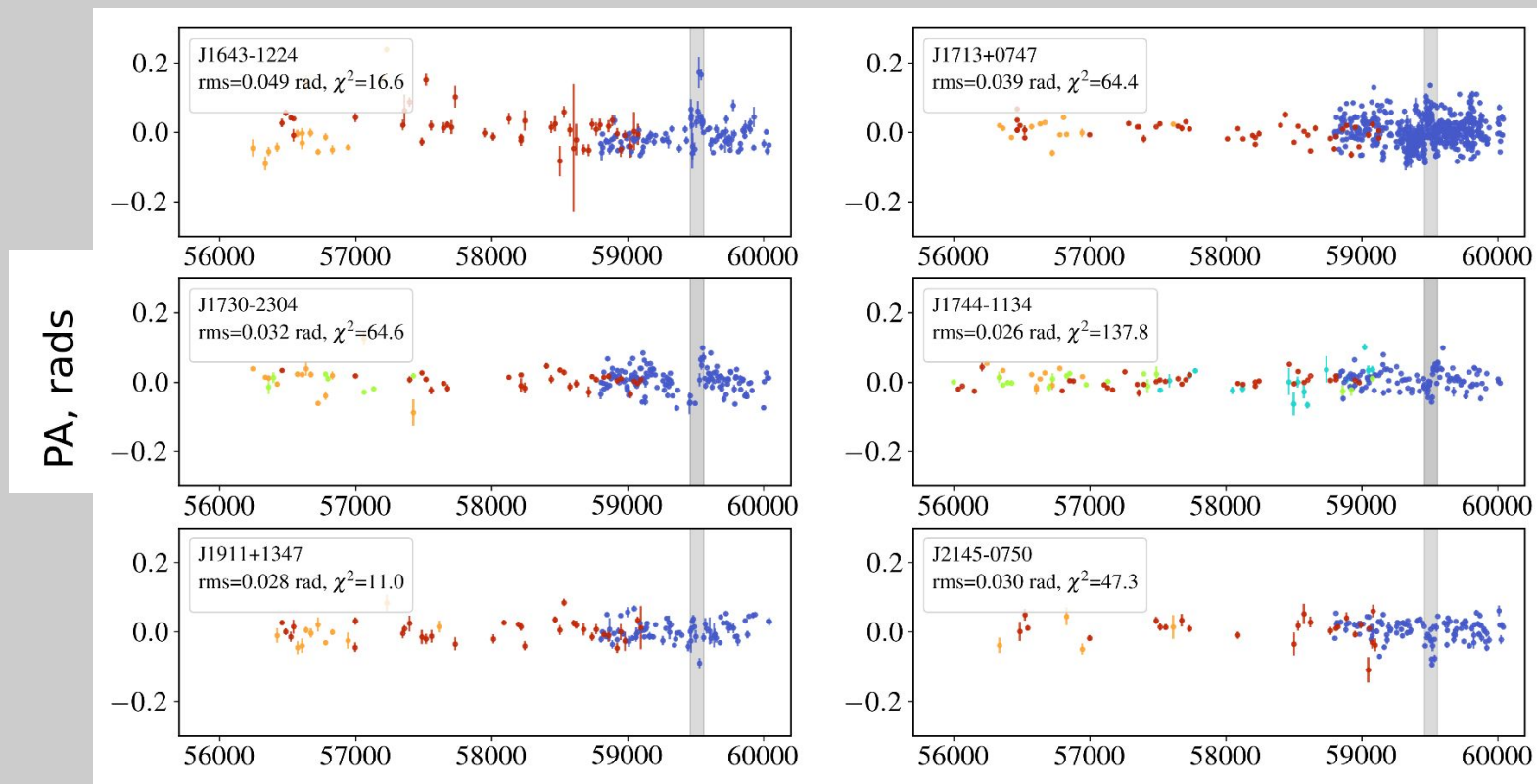
II. ULDM with pulsar polarimetry: challenges



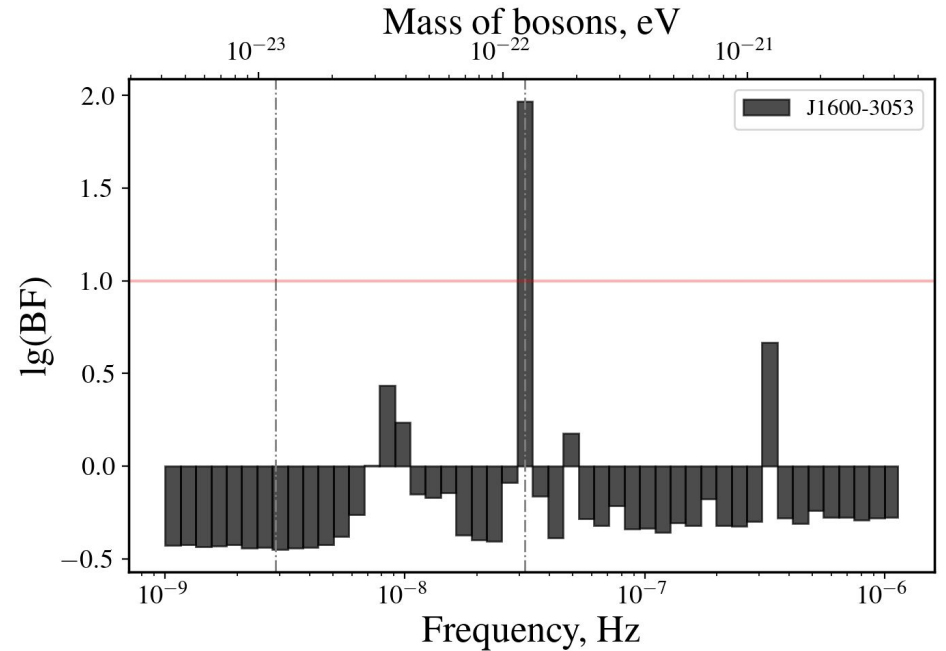
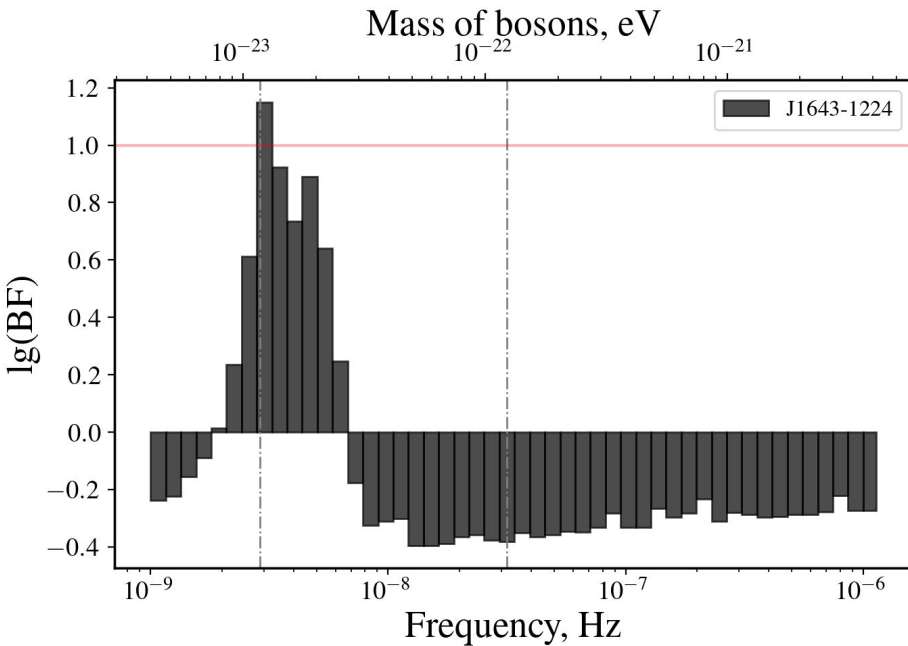
II. ULDM with pulsar polarimetry: dataset



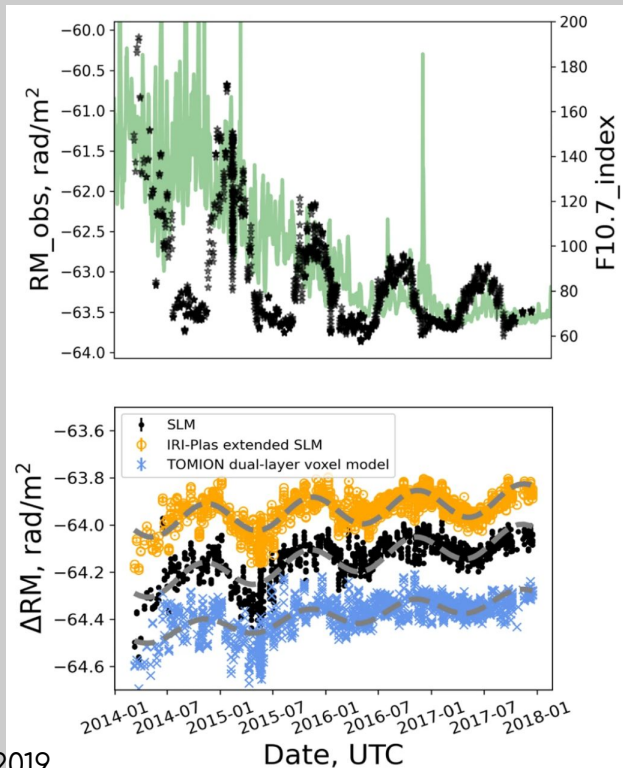
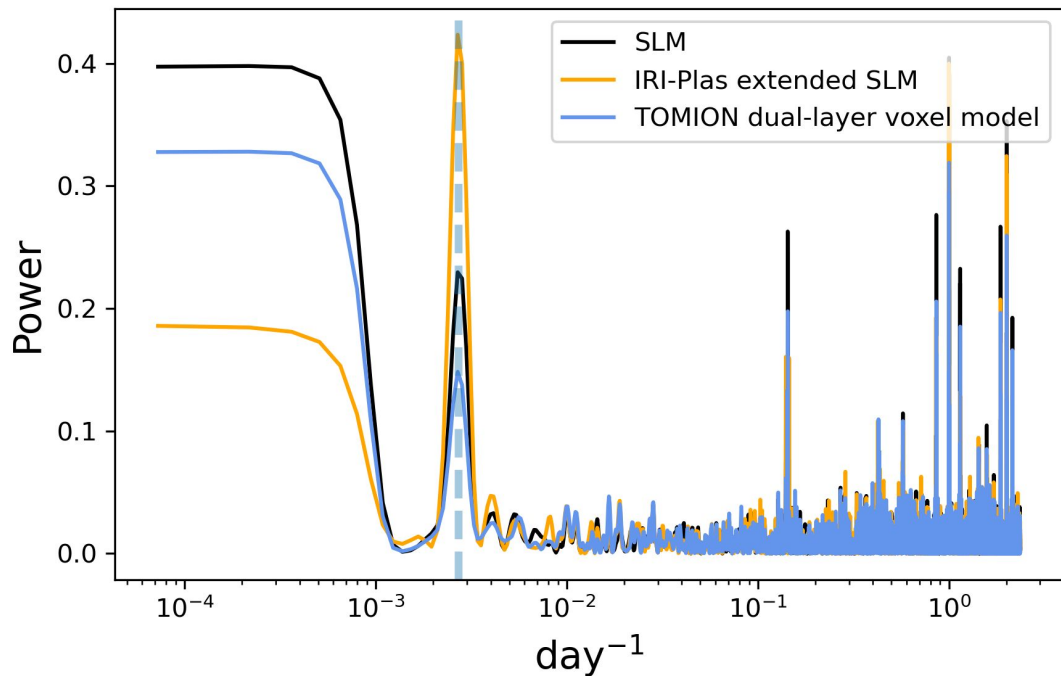
II. ULDM with pulsar polarimetry: dataset



II. ULDM with pulsar polarimetry: back to the ionosphere

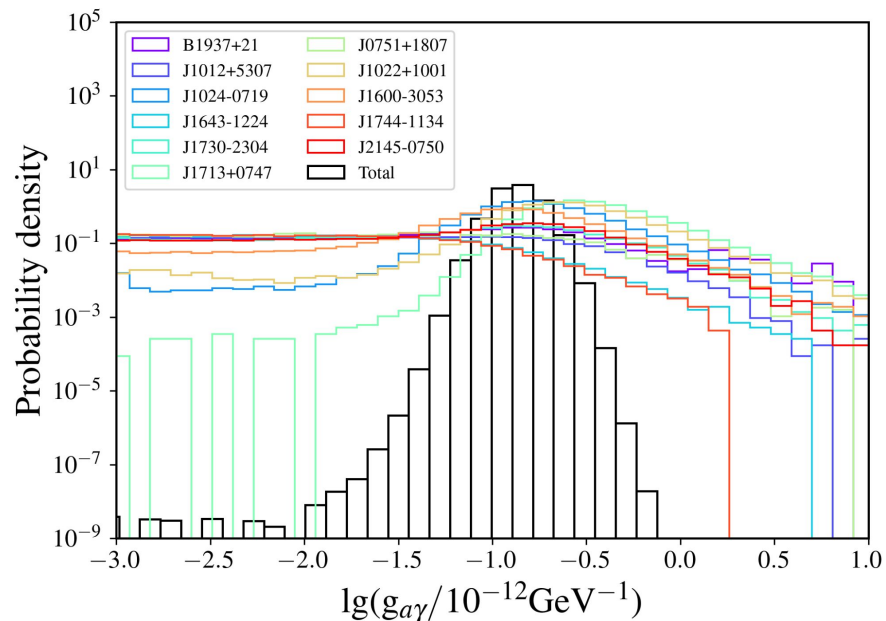
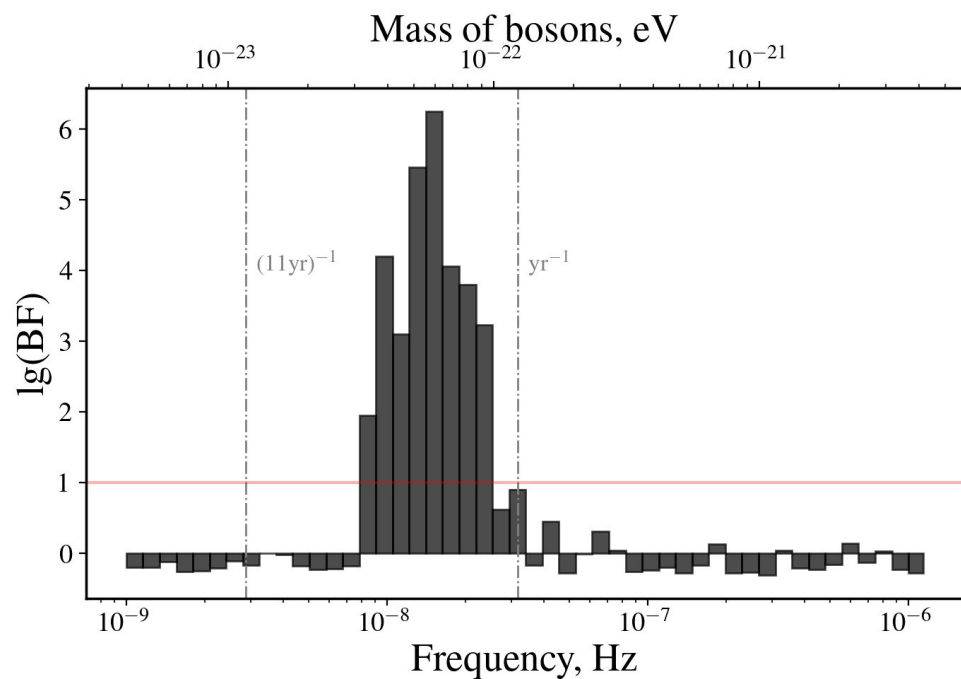


II. ULDM with pulsar polarimetry: back to the ionosphere

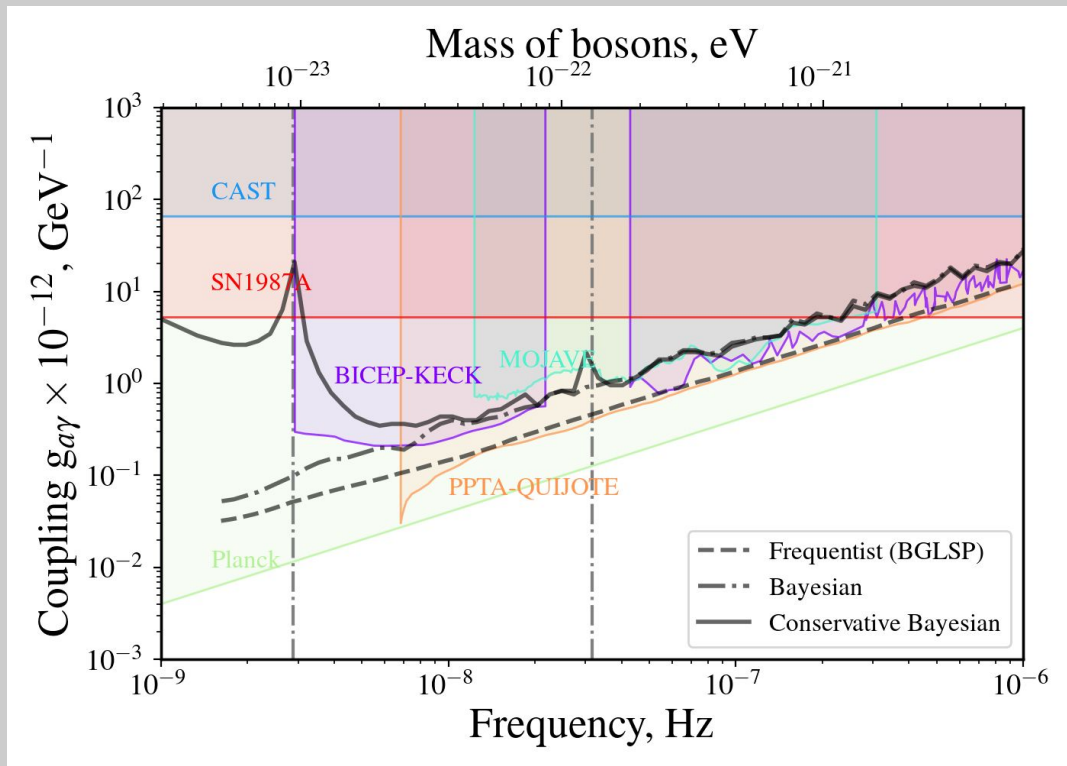


Porayko et al 2019,
Porayko et al 2023

Factorised upper limits and BFs



III. ULDM with pulsar polarimetry: first results



Ultra-light axions in the Milky Way:

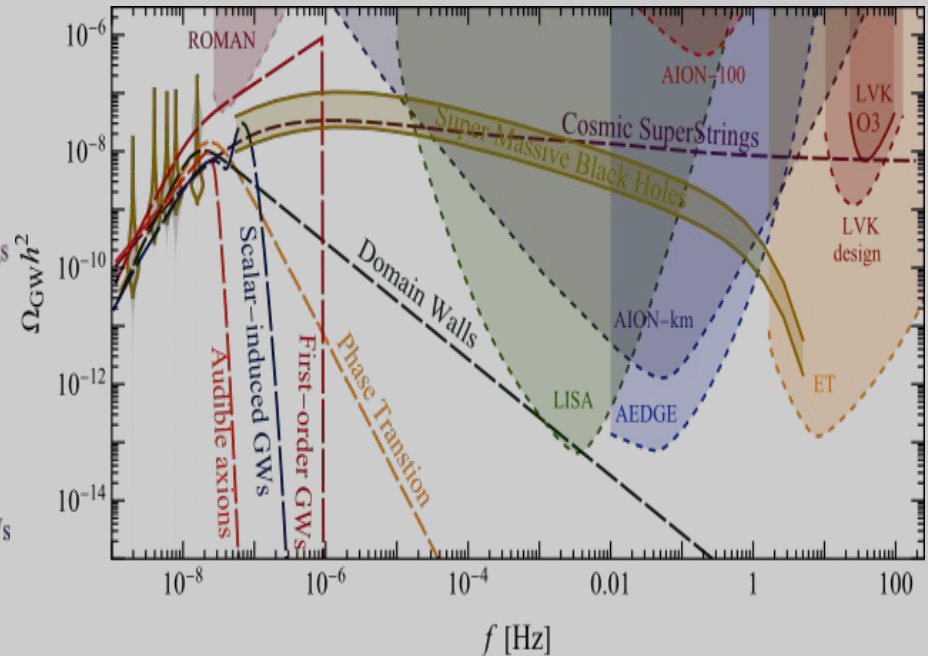
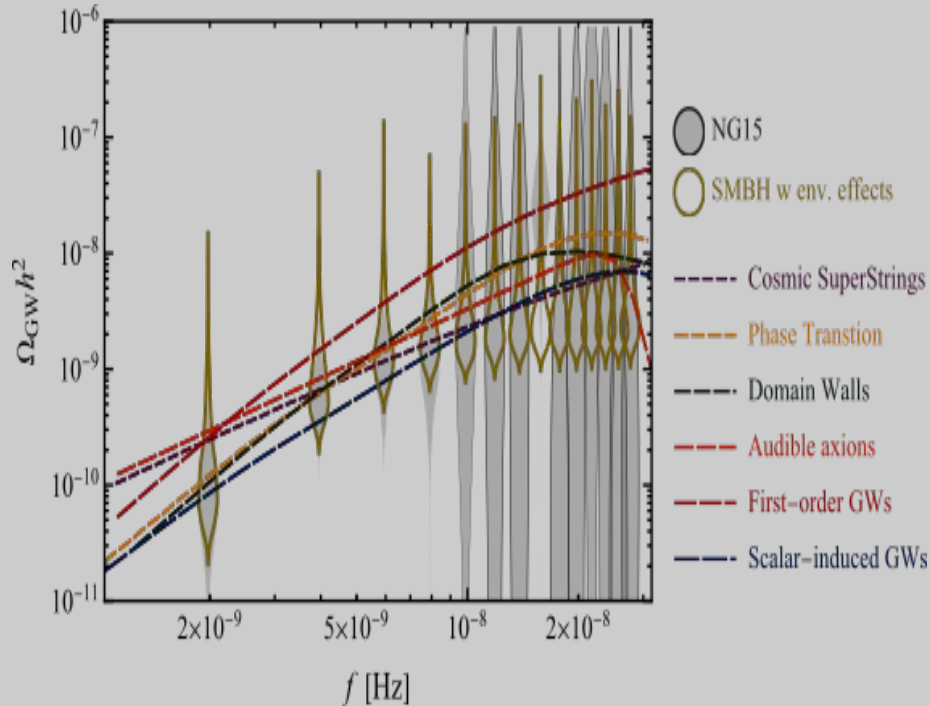
3. The **effect is achromatic**, so can be distinguished from chromatic Faraday rotation

4. **Terrestrial ionosphere** is the main source of noise, when searching for ultra-light axions in pulsar polarimetry

5. We plan to incorporate **low-frequency data** from LOFAR(2.0) to independently mitigate ionospheric Faraday rotation

The true picture & a wider landscape

So we are tantalisingly close to *detecting* the GWB, but what is the true nature of this signal?



Ellis et al (2023; 2308.08546)

IPTA DR3 dimensions

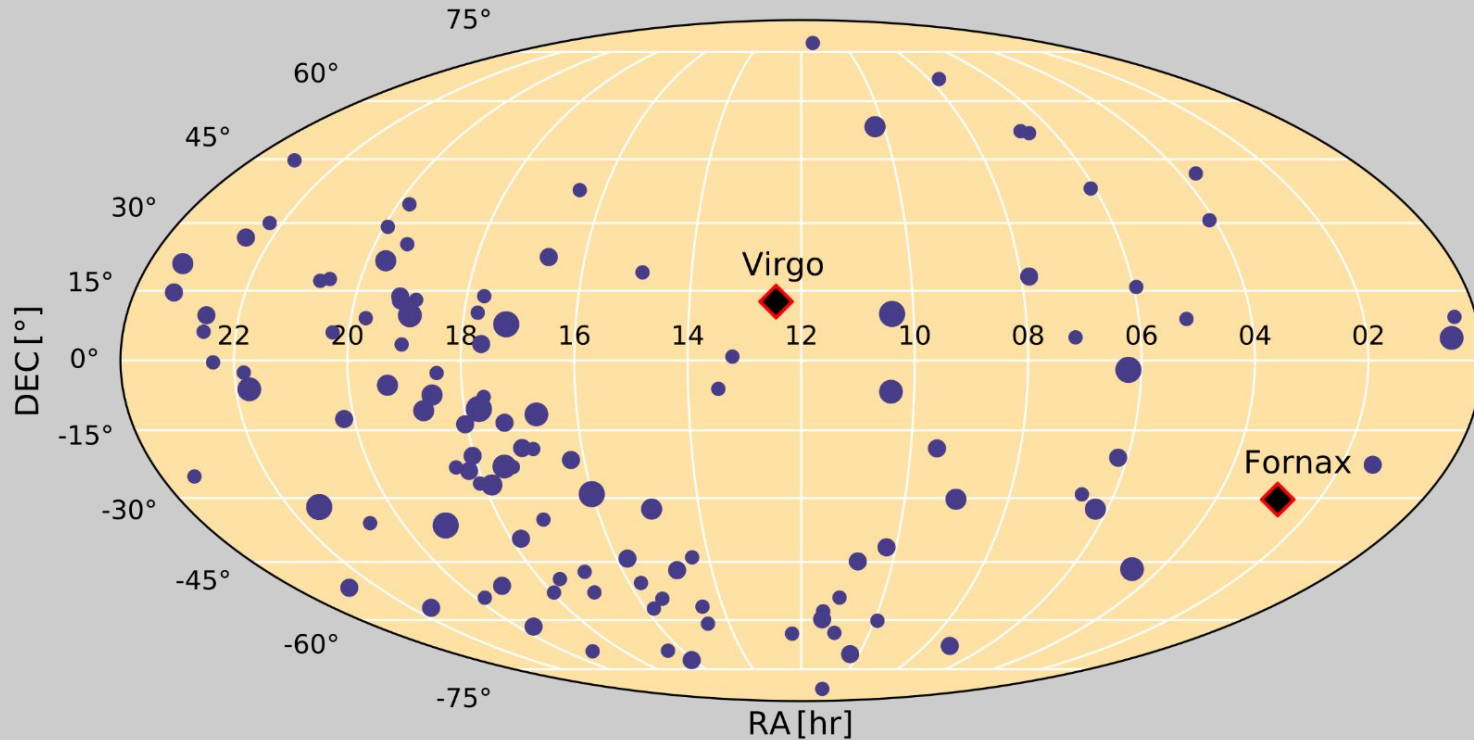


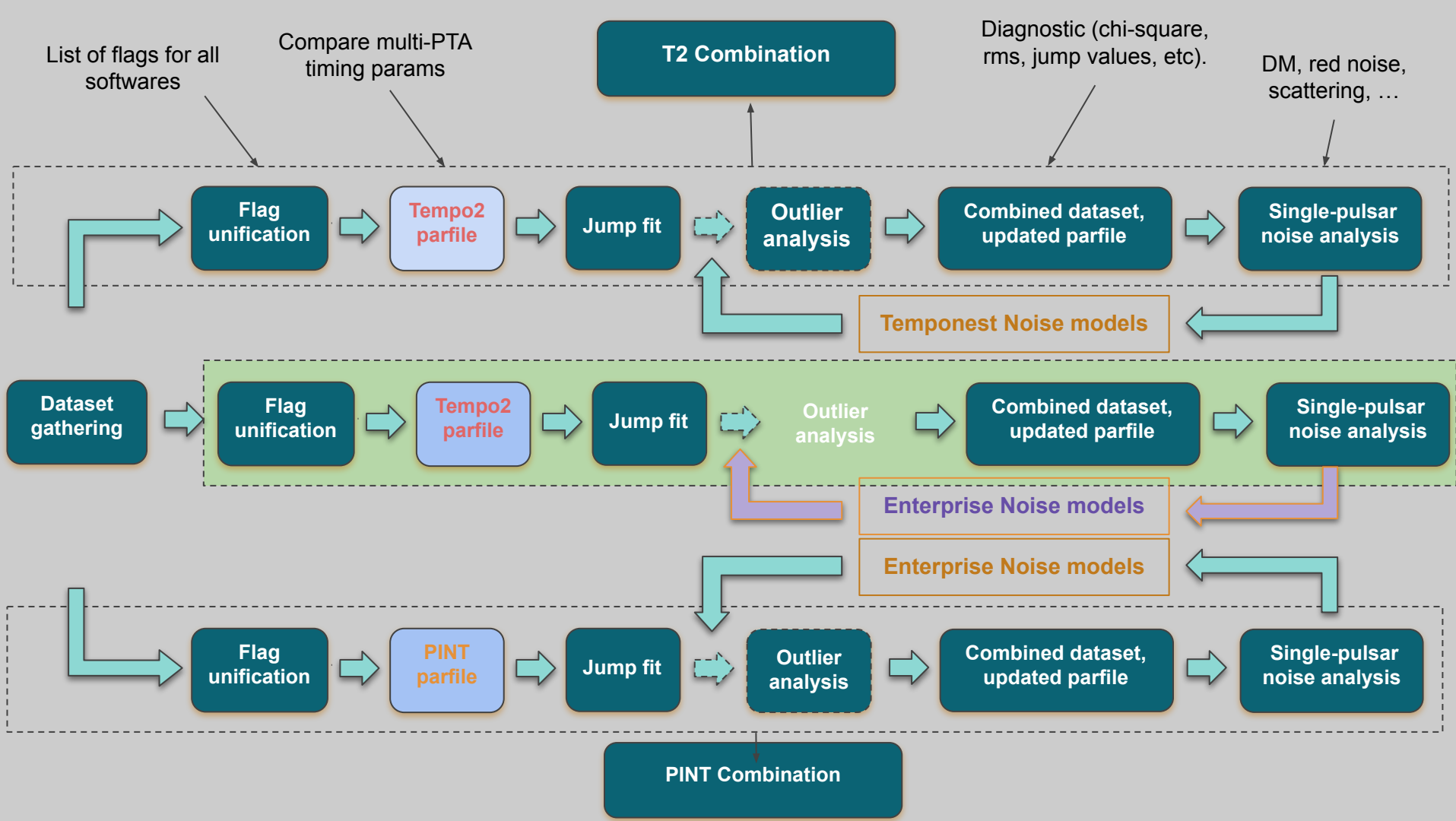
PTA	Dataset	PSRs	Tspan (years)	$f_{\text{GW,low}}$ (nHz)	f_{radio} (MHz)
EPTA	DR2 / DR3	25 / +35	24.5	1.29	283 - 5107
	LOFAR + NENUFAR	17	9.6	-	30 - 190
NANOGrav	15-yr	68	15.9	1.99	302 - 3988
	CHIME	11	2.5		400 - 800
PPTA	DR3	24	18.1	1.75	704 - 4032
InPTA	DR1	14	3.5	9.05	300 - 1460
MeerKAT	DR2	88	4.5	7.04	856 - 1712
IPTA	DR3	121	~25/40	1.29/0.79	30 - 5107

IPTA DR3 dimensions

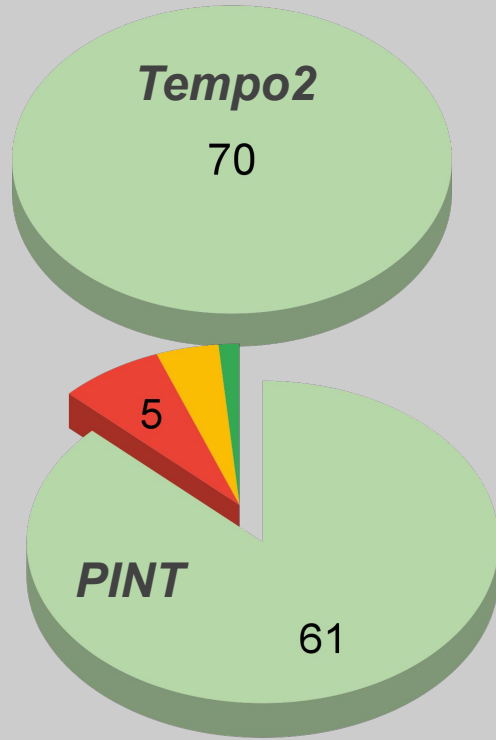


- In total **121** pulsars in full DR3;
 - The biggest / most sensitive PTA dataset ever made !!

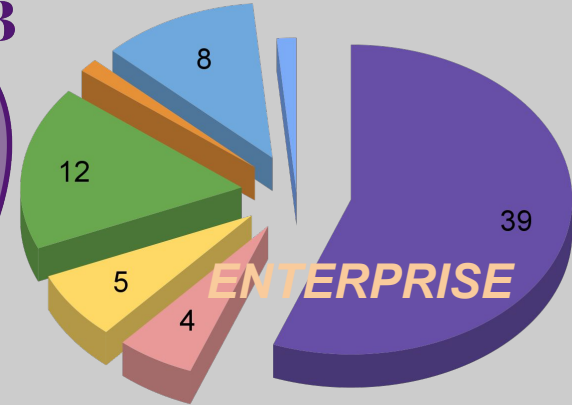
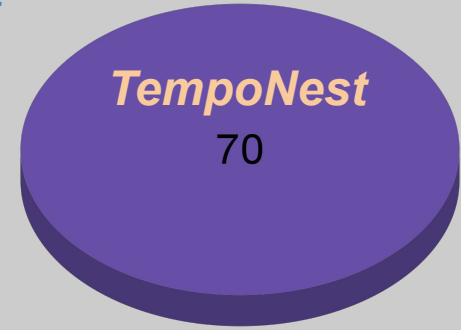
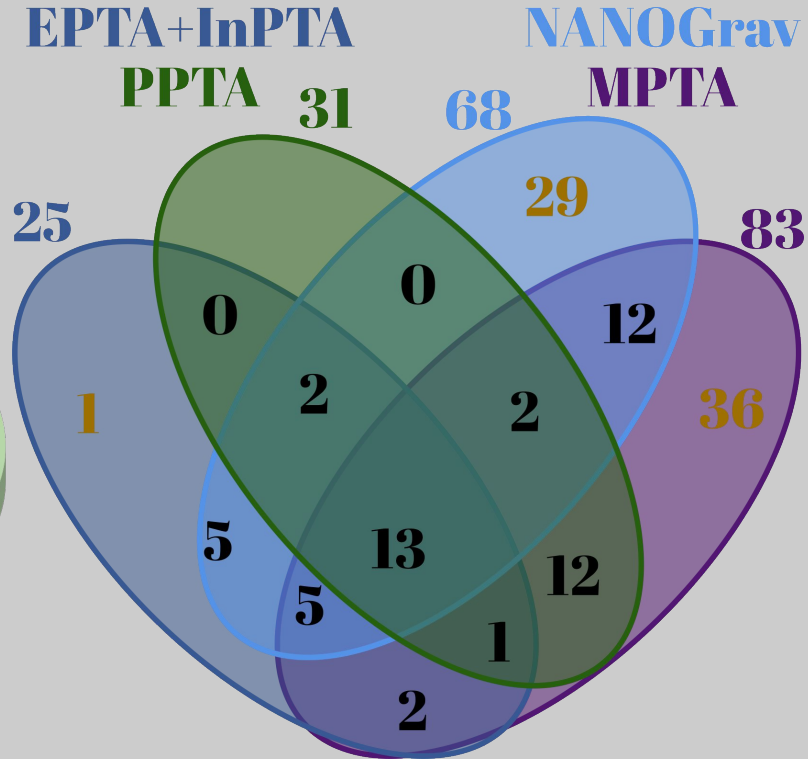




Current Status

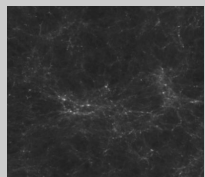


● Done ● NA ● In progress ● Updated par file



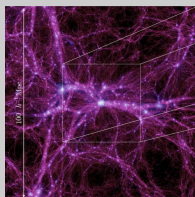
● Done ● Missing Tim ● waiting on config file ● NA
● config not found ● Running ● tim par not copied

DARK MATTER MERGER TREES



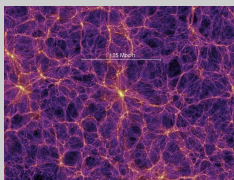
$$L_{\text{box}} = 30 \text{ Mpc} / h$$

$$M_{\text{halo}} \sim 10^6 M_{\text{sun}}$$



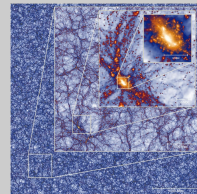
$$L_{\text{box}} = 100 \text{ Mpc} / h$$

$$M_{\text{halo}} \sim 10^8 M_{\text{sun}}$$



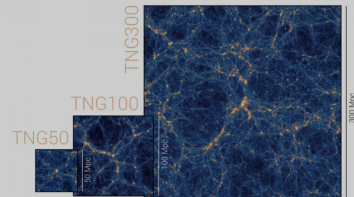
$$L_{\text{box}} = 500 \text{ Mpc} / h$$

$$M_{\text{halo}} \sim 10^{10} M_{\text{sun}}$$



$$L_{\text{box}} = 3 \text{ Gpc} / h$$

$$M_{\text{halo}} \sim 10^{11} M_{\text{sun}}$$



$$L_{\text{box}} = 50 \text{ Mpc}, 100 \text{ Mpc}, 300 \text{ Mpc}$$

$$M_{\text{halo}} \sim 10^7 M_{\text{sun}}, 10^8 M_{\text{sun}}, 10^{10} M_{\text{sun}}$$

BARYONIC PHYSICS

Galaxy physics

Massive black hole physics

Galaxy merger

~ kpc

Galaxy nucleus

~ pc

~ 10⁻² pc

BH merger

PAIRING

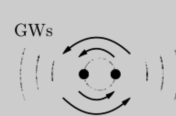
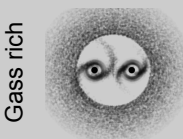
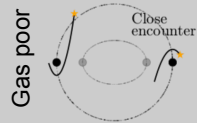
$$t_{\text{dyn}}^{\text{BH}} = 19 f(\varepsilon) \left(\frac{r_0}{4 \text{ kpc}} \right)^2 \left(\frac{\sigma}{200 \text{ km/s}} \right) \left(\frac{10^8 M_{\odot}}{M_{\text{BH}}} \right) \frac{1}{\Lambda} [\text{Gyr}]$$

TRACK SELF-CONSISTENTLY VIA NUMERICAL INTEGRATION THE

- BINARY SEPARATION (a_{BH})

- BINARY ECCENTRICITY (e_{BH})

HARDENING



GRAVITATIONAL WAVE INSPIRAL

$$\frac{da_{\text{BH}}}{dt} = \left(\frac{da_{\text{BH}}}{dt} \right)_{\text{Hard}} + \left(\frac{da_{\text{BH}}}{dt} \right)_{\text{GW}} = - \frac{GH \rho_{\text{inf}}}{\sigma_{\text{inf}}} a_{\text{BH}}^2 - \frac{64 G^3 (M_{\text{BH}_1} + M_{\text{BH}_2})^3 F(e)}{5 c^5 (1+q)^2 a_{\text{BH}}^3}$$

$$\frac{de}{dt} = a_{\text{BH}} \frac{G \rho_{\text{inf}} H K}{\sigma_{\text{inf}}} - \frac{304}{15} \frac{G^3 q (M_{\text{BH}_1} + M_{\text{BH}_2})^3}{c^5 (1+q)^2 a_{\text{BH}}^4 (1-e^2)^{5/2}} \left(e + \frac{121}{304} e^3 \right)$$

$$\frac{da_{\text{BH}}}{dt} = - \frac{2 \dot{M}}{\mu} \sqrt{\frac{\delta}{1-e^2}} a_{\text{BH}}$$

Triplets Bonetti et al. 2018

Gas accretion $\dot{M}_{\text{BH}_1} = \dot{M}_{\text{BH}_2} (0.1 + 0.9q)$

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Jishnu Suresh



Lucia Papalini

LAAC



Jessica Steinlechner



Mikhail Korobko



PTA



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