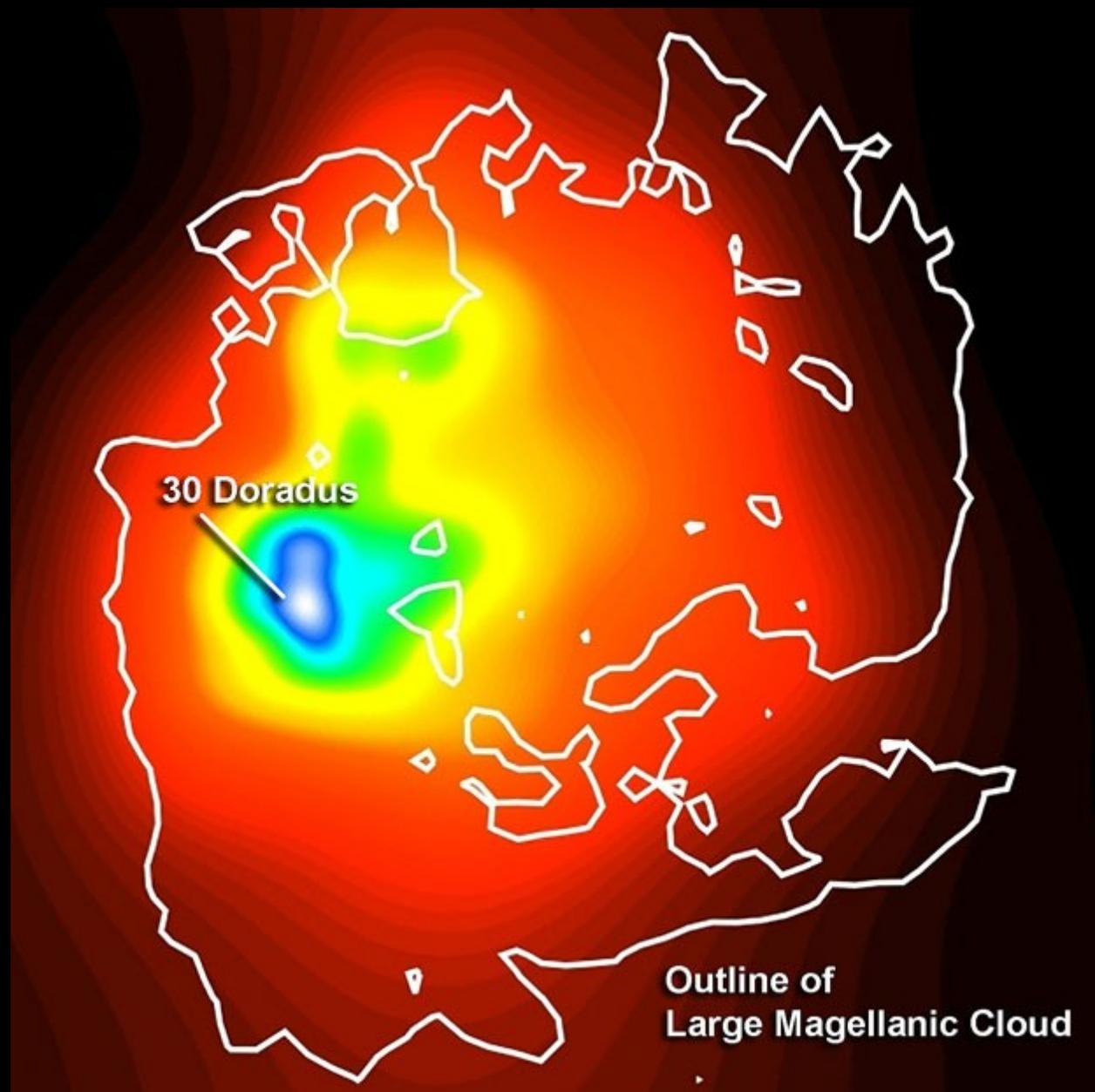


Cosmic Rays in other galaxies





30 Doradus

Outline of
Large Magellanic Cloud

Summary of the lesson

Cosmic Rays

CR and the Galaxy

γ Rays prod. mech.

The MilkyWay in γ Rays

Diffuse gas

Molecular Clouds

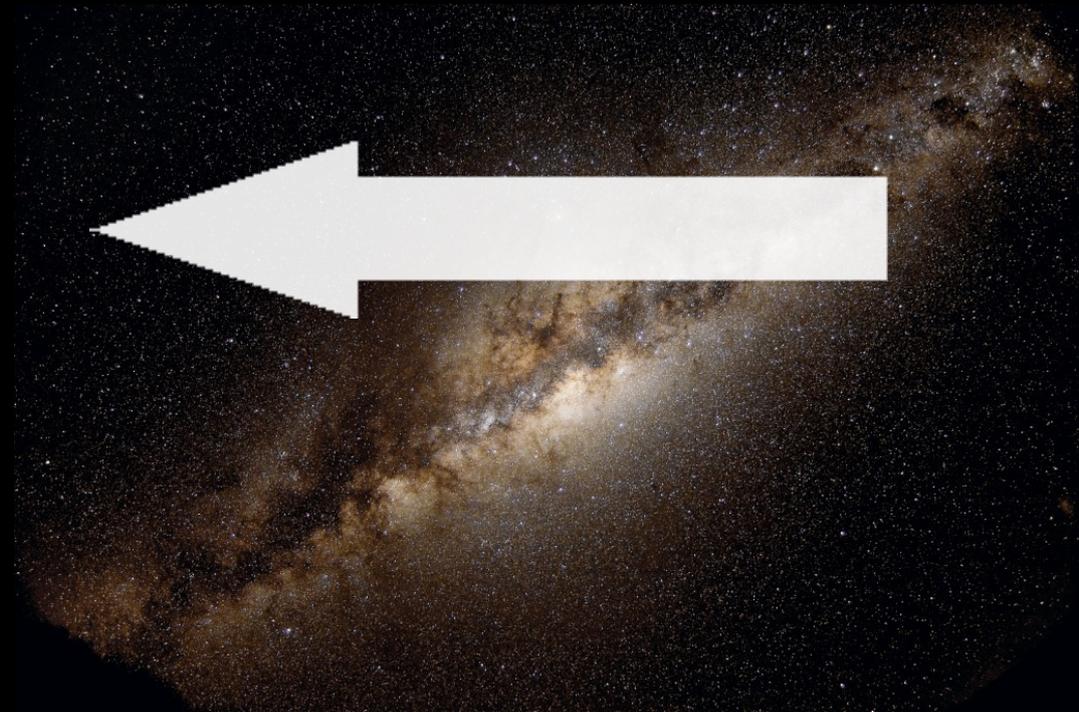
InterStellar Radiation Field

Supernova Remnants

Evolution

MC associations

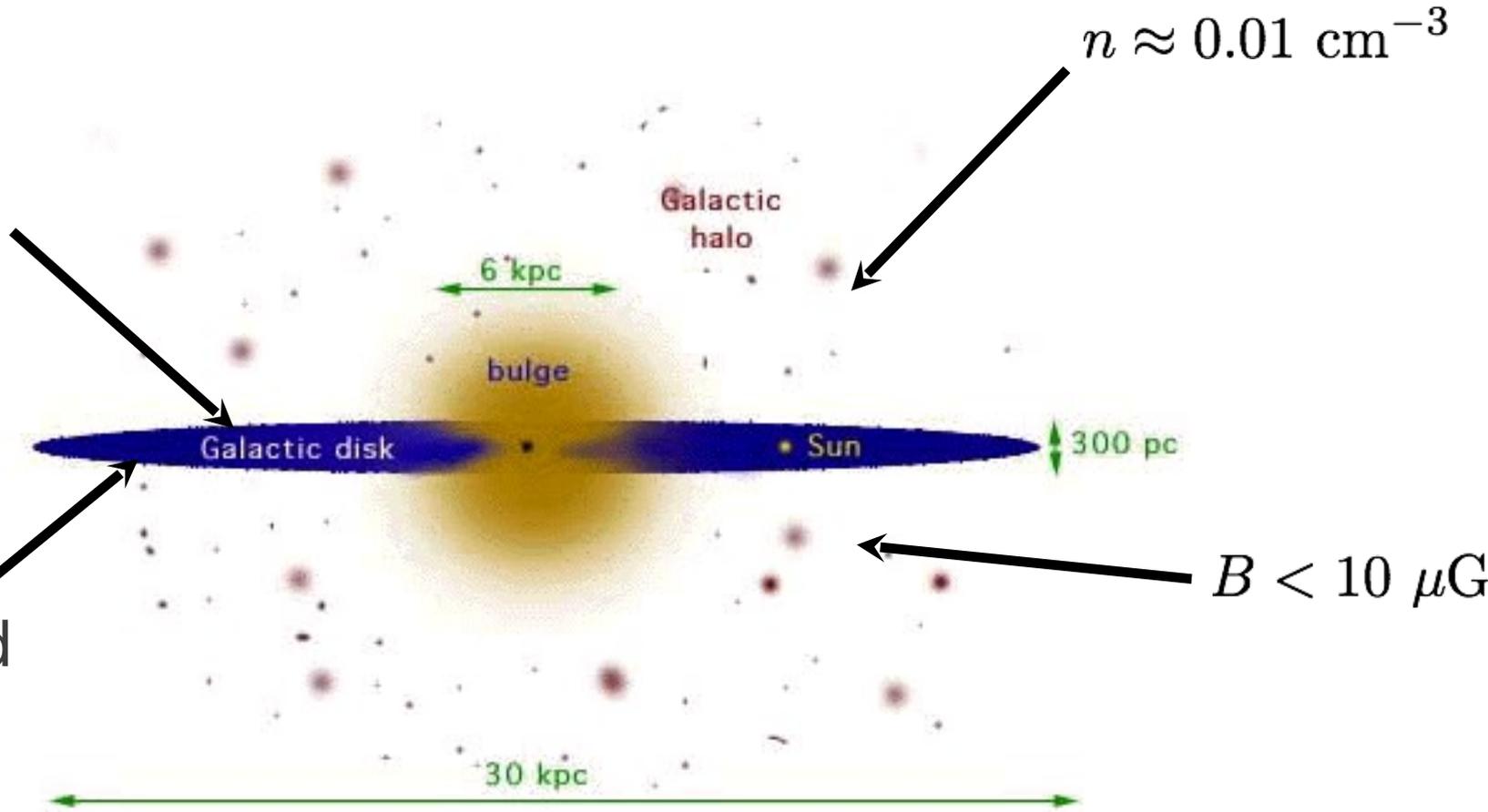
Gamma-ray from SNR



the Milky Way

gas density
 $n \approx 1 \text{ cm}^{-3}$

$n \approx 0.01 \text{ cm}^{-3}$

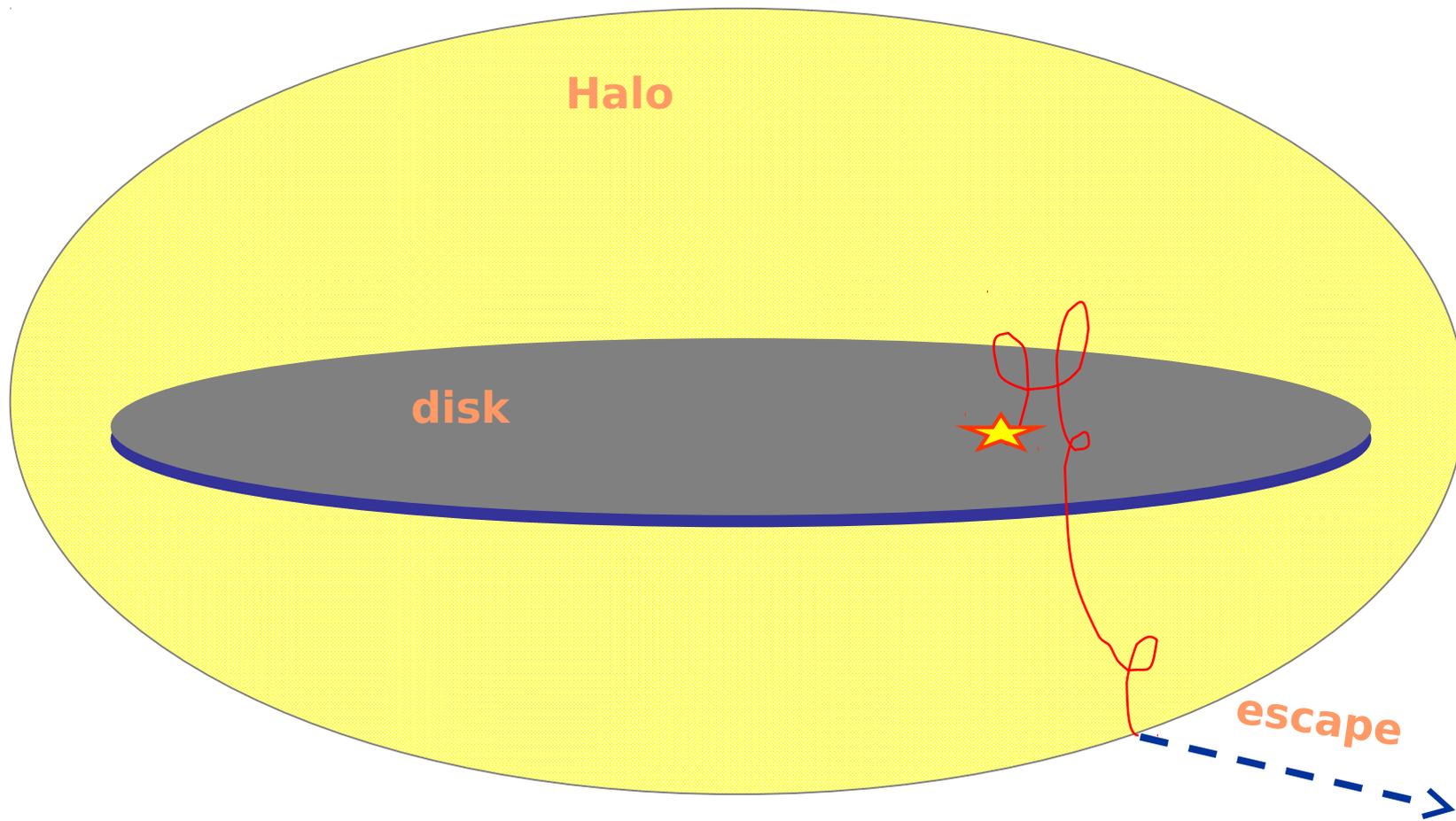


magnetic field

$B \approx 3 \mu\text{G}$

$B < 10 \mu\text{G}$

CR Propagation



Cosmic Ray power in the Galaxy

$$w_{CR} \sim 1 \text{ eV/cm}^3 \longrightarrow \mathcal{E}_{CR} = w_{CR} V_{disk}$$

Cosmic Ray power in the Galaxy

$$w_{CR} \sim 1 \text{ eV/cm}^3 \longrightarrow \mathcal{E}_{CR} = w_{CR} V_{disk}$$

$$P_{CR} = \frac{\mathcal{E}_{CR}}{t_{disk}} = \frac{w_{CR} V_{disk}}{t_{disk}} = 3 \times 10^{40} \text{ erg/s}$$

The Supernovae power in the Galaxy

A **supernova** releases $\sim 10^{51}$ ergs in form of kinetic energy.
In the Galaxy the observed supernova rate is of the order of $1/30 - 1/100 \text{ yr}^{-1}$.

$$P_{SN} = 3 \times 10^{41} \text{ erg/s}$$

The Supernovae power in the Galaxy

A **supernova** releases $\sim 10^{51}$ ergs in form of kinetic energy.
In the Galaxy the observed supernova rate is of the order of $1/30 - 1/100 \text{ yr}^{-1}$.

$$P_{SN} = 3 \times 10^{41} \text{ erg/s}$$

$$P_{CR} = 3 \times 10^{40} \text{ erg/s}$$

SuperNovae alone could maintain the CR population provided that about **10%** of their kinetic energy is **somehow converted into CRs**

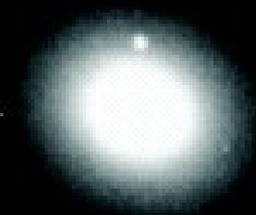
Ginzburg e Syrovatskii (1964)

SuperNova Remnants



M101

Spiral galaxies:
Host all types of SNe
Stellar population:
a mix of old and young



M49

Elliptical galaxies:
Only Type Ia
Stellar population: old

The Galaxy: +/- 2 supernovae per century

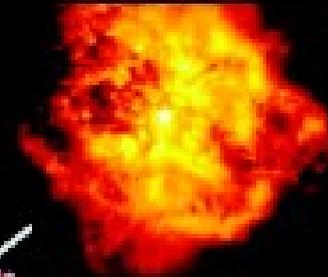


SuperNova Remnants

Massive stars $> 7 M_{\text{sun}}$



Very Massive stars $> 22 M_{\text{sun}}$
Wolf-Rayet Stars



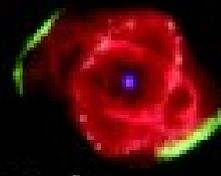
Main Seq.

Red Super Giant Phase



Supernova!

Stars $< 7 M_{\text{sun}}$



Red giant

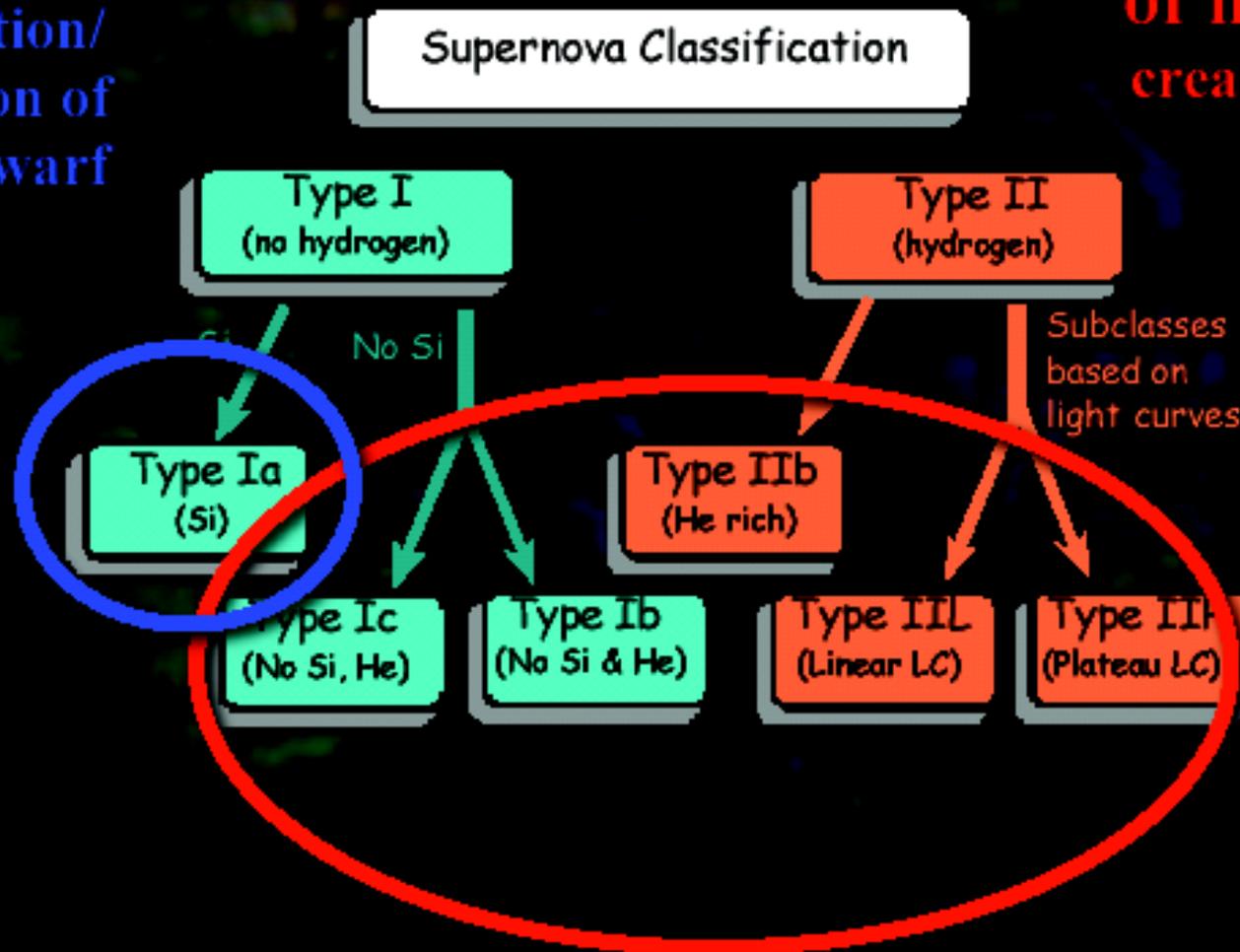
Planetary Nebulae White dwarf



SuperNova Remnants

Thermonuclear explosion
deflagration/
detonation of
White Dwarf

Core collapse
of massive star
creation of NS/BH



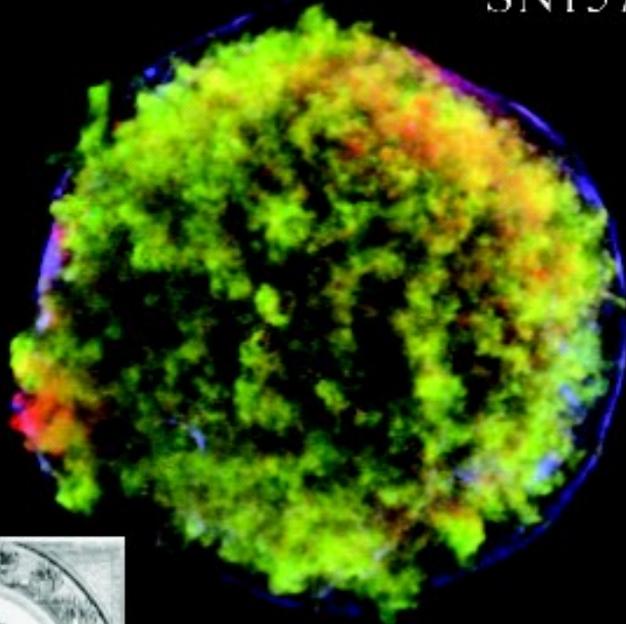
SuperNova Remnants

- Type Ia:
- The whole star is disrupted by the explosion
 - The source of energy is nuclear fusion, predominantly the burning of C/O into ^{56}Ni
 - Most of the energy is in the form of heat (10^{51} erg)
- Type II/Ibc:
- The core of the star collapses into a neutron star
 - The source of energy is therefore gravity ($\sim GM^2/R_{\text{NS}} \sim 10^{53}$ erg)
 - Most of the energy released as neutrinos!
 - Only 1% converted to heat/kinetic energy!
 - Nuclear fusion: by-product/not source of the explosion



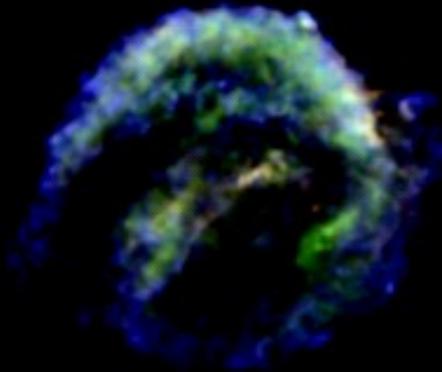
SuperNova Remnants

SN1572



Chandra

SN1604

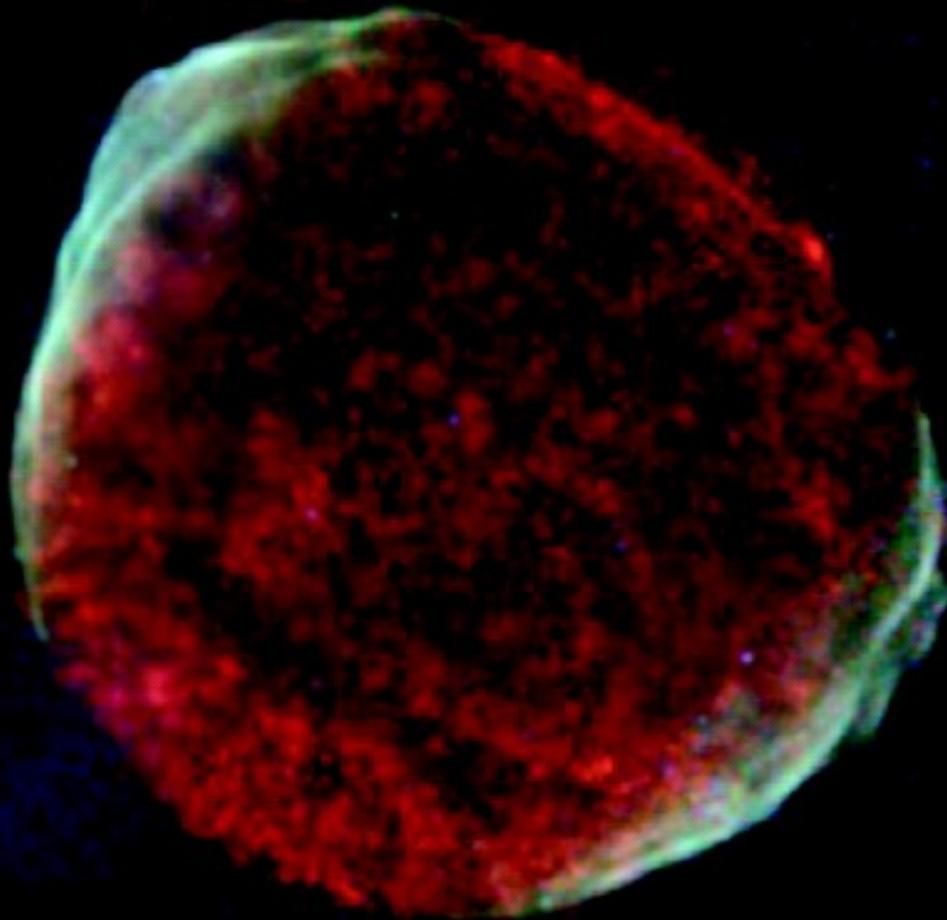


Tycho Brahe



Johannes Kepler

SuperNova Remnants



Chandra

Size: 30'

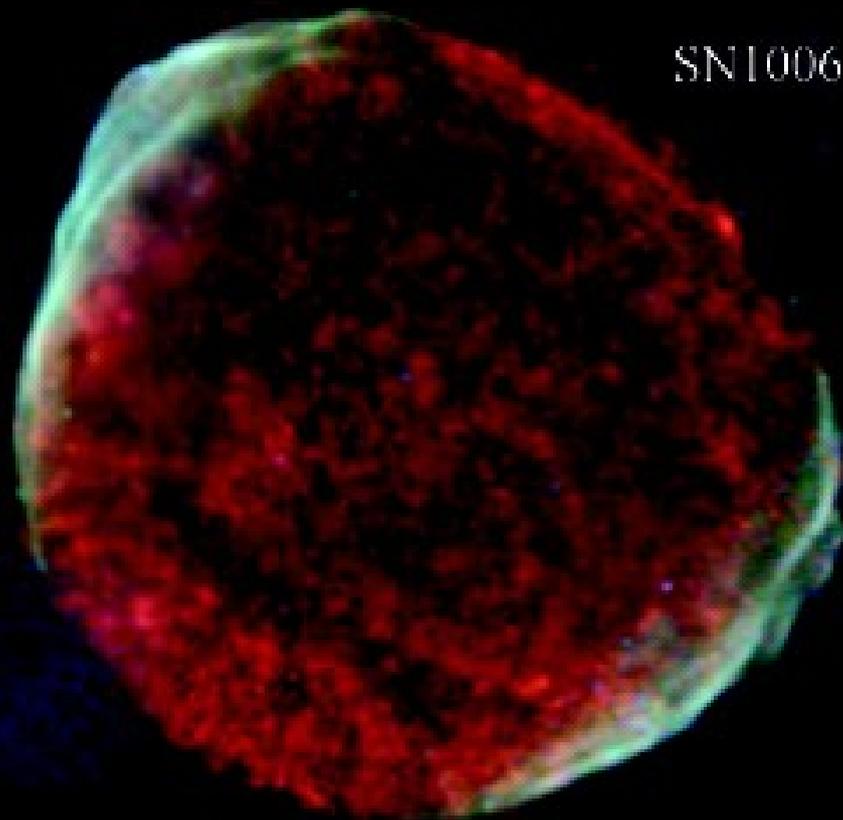
The size of the full Moon

Historical brightness

-6 mag, perhaps -9 mag

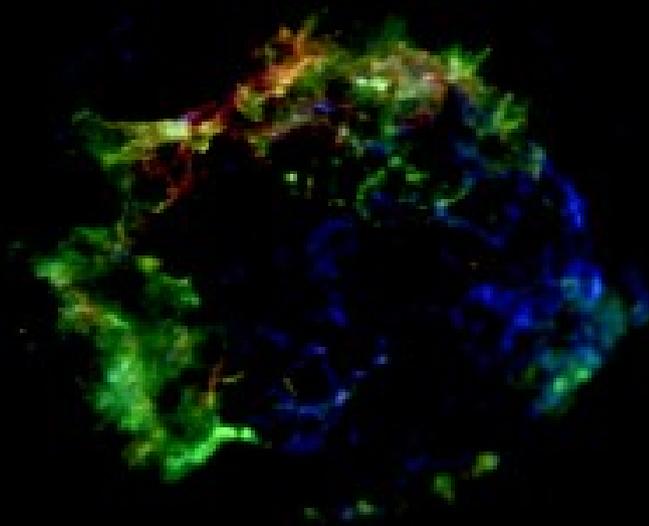
SuperNova Remnants

Shell Types: Shell of shock heated gas



SN1006

Type Ia supernova
Shell Type SNR

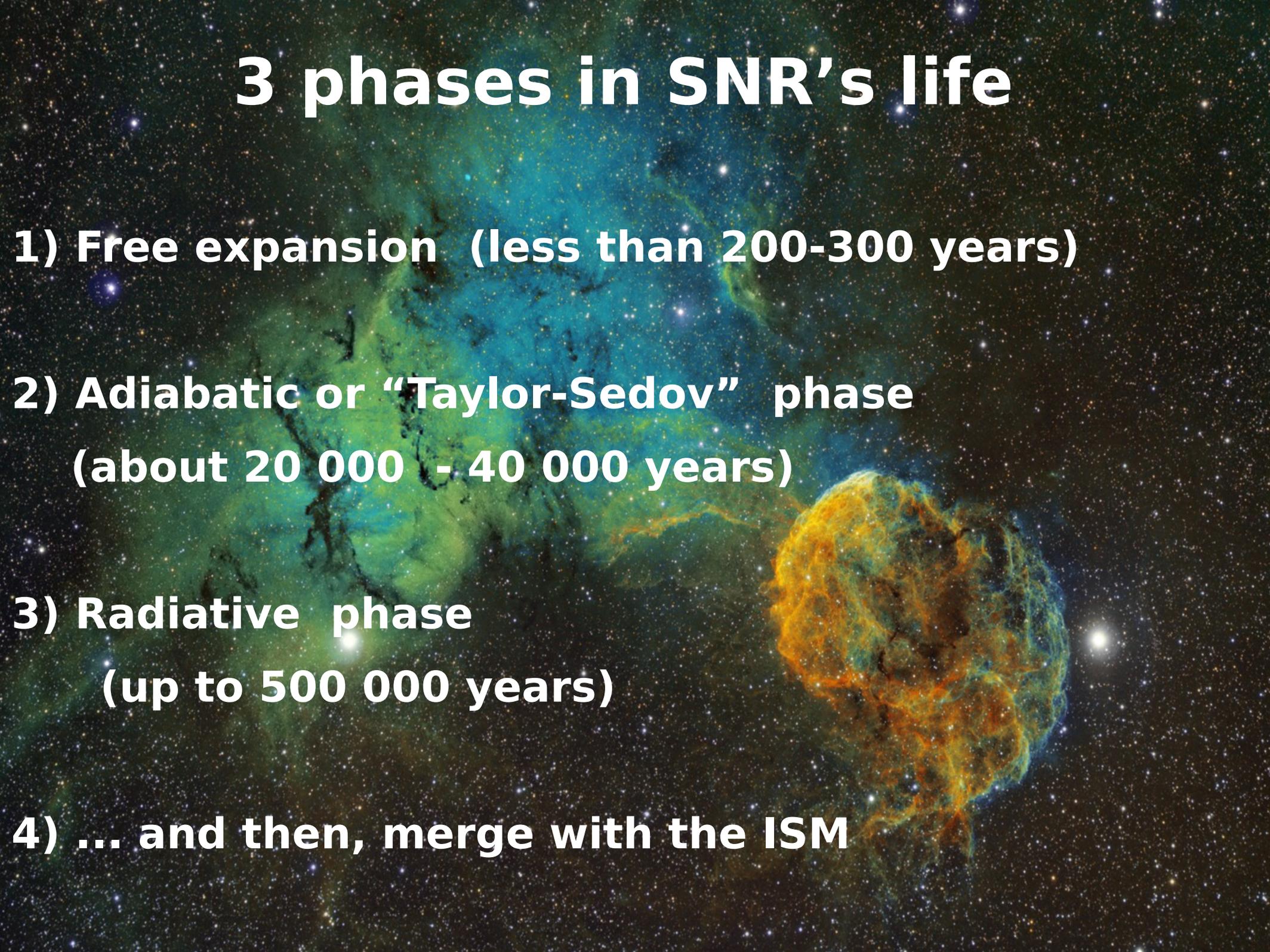


Cas A

O VIII
Si XIII
Continuum

Core-collapse supernova (Type Ib)
Shell Type SNR

3 phases in SNR's life



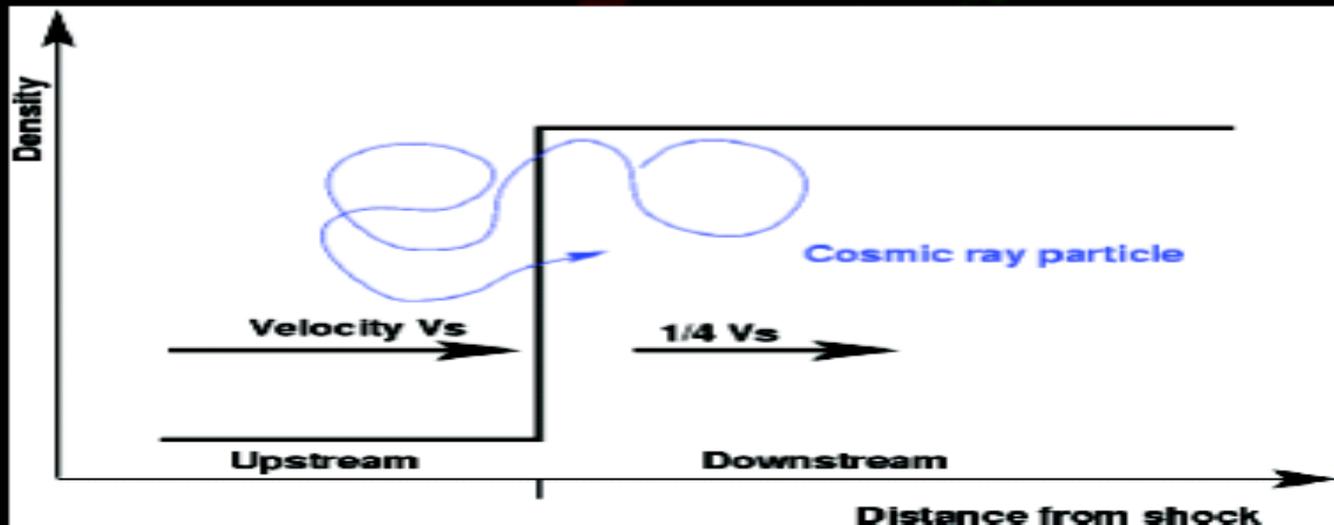
1) Free expansion (less than 200-300 years)

**2) Adiabatic or “Taylor-Sedov” phase
(about 20 000 - 40 000 years)**

**3) Radiative phase
(up to 500 000 years)**

4) ... and then, merge with the ISM

Shocks



- Conservation laws: mass, momentum and energy conservation:
Use system in which shock is at rest

$$\rho_1 v_1 = \rho_2 v_2$$

$$(\rho_1 v_1) v_1 + p = (\rho_2 v_2) v_2 + p$$

$$(1/2 \rho_1 v_1^2 + u) v_1 = (1/2 \rho_2 v_2^2 + u) v_2$$

internal energy $u = p/(\gamma - 1)$, $\gamma = 5/3$ for monatomic gas

- Simplification: heat sinks (cosmic ray acceleration!), magnetic fields, and radiation losses not taken into account.

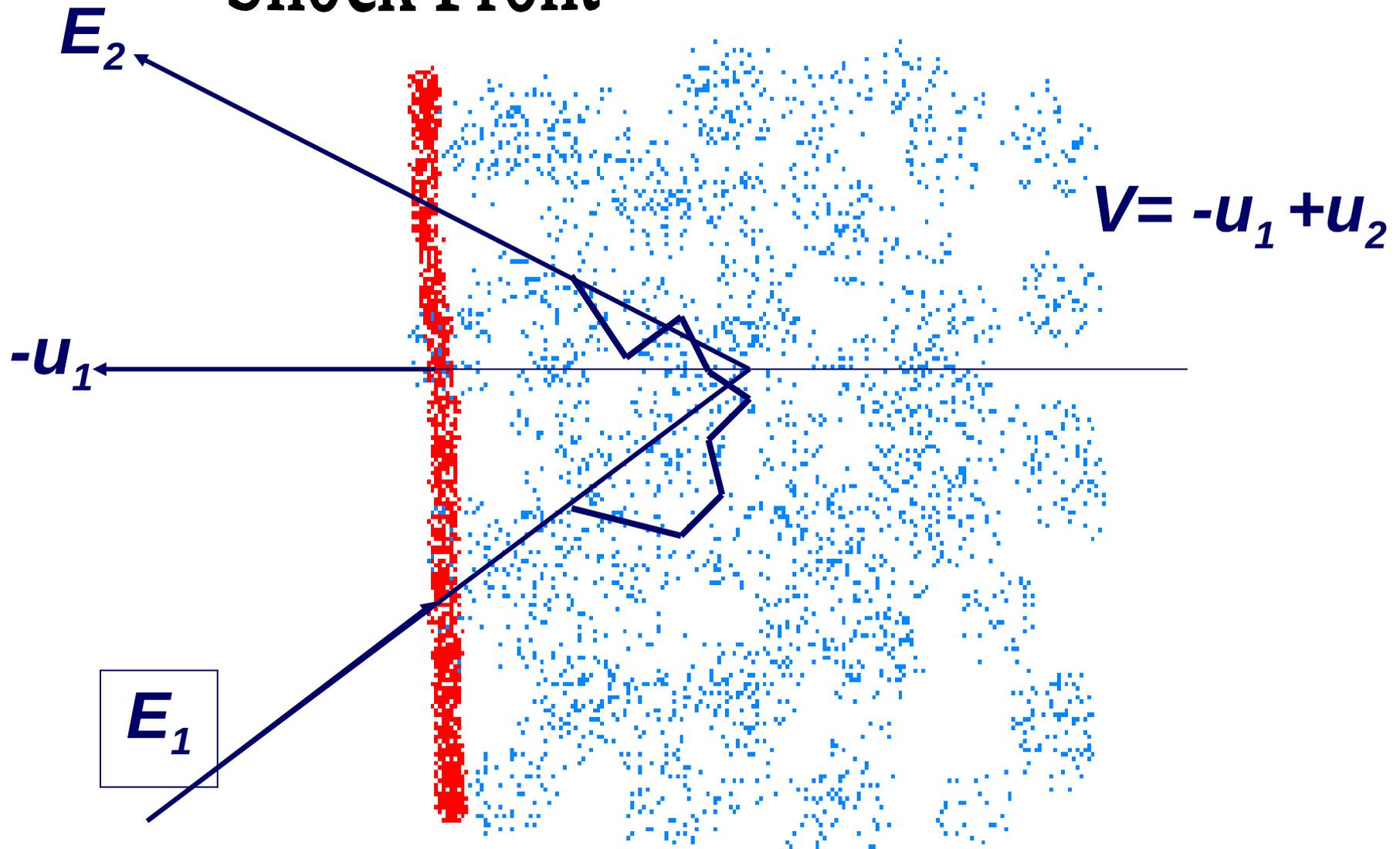
- For strong shocks ($M \rightarrow \infty$) one finds:

$$\rho_2 / \rho_1 = (\gamma + 1) / (\gamma - 1) = 4, \text{ implying } v_2 = 1/4 v_s$$

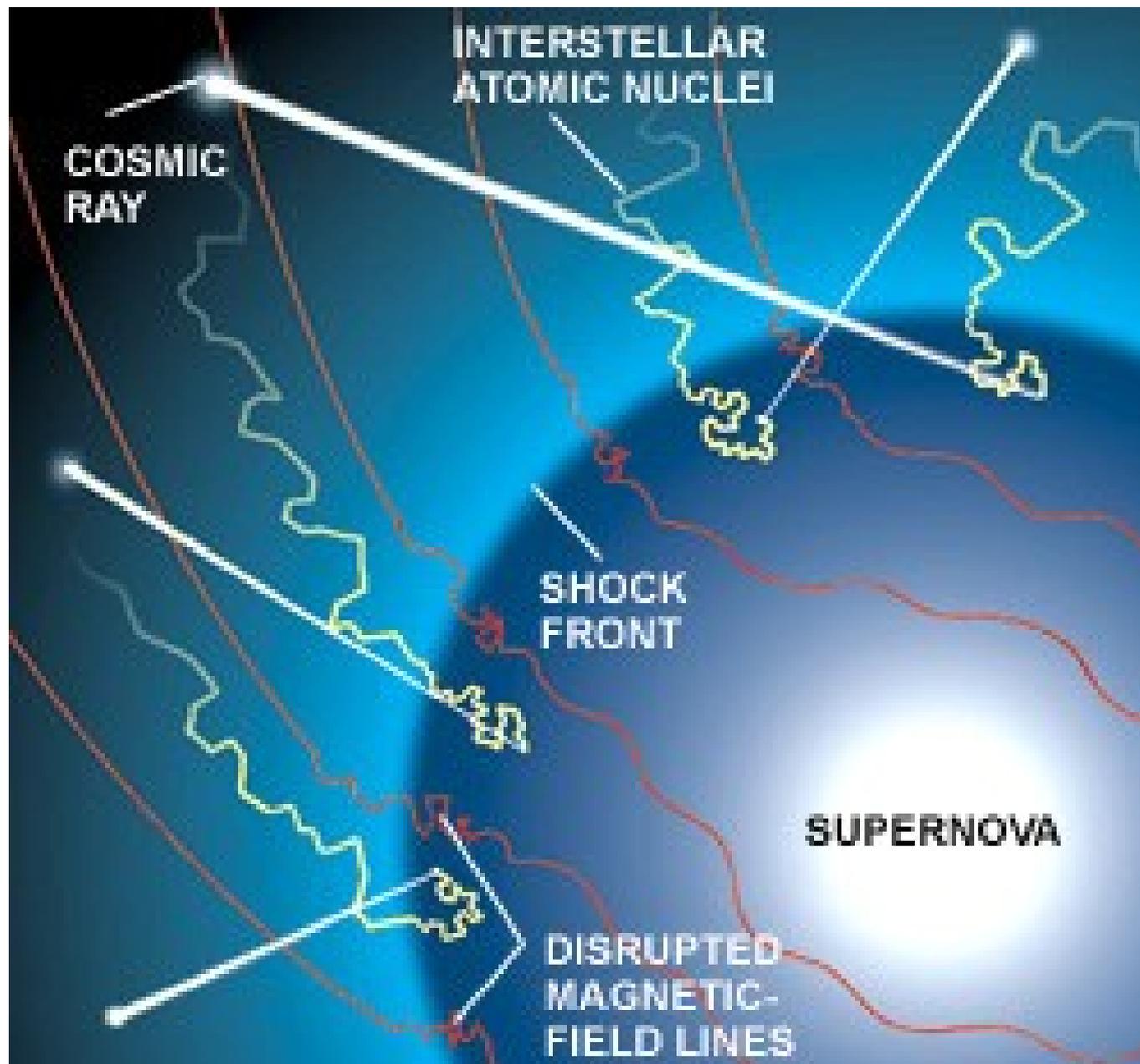
$$kT_2 = 2(\gamma - 1)(\gamma + 1)^{-2} m v_s^2 = 3/16 m v_s^2, \text{ with } m \text{ particle mass}$$

Accelerazione di Fermi del primo ordine

Shock Front

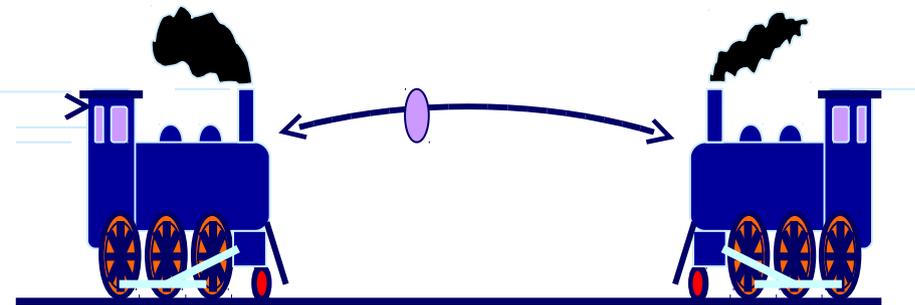


Meccanismo di accelerazione del primo ordine di Fermi



The Fermi acceleration

toy model



The resulting energy spectrum of many particles undergoing this process (assuming that they do not influence the structure of the shock) turns out to be a power law:

$$\frac{dN(\epsilon)}{d\epsilon} \propto \epsilon^{-p} \quad p \geq 2$$

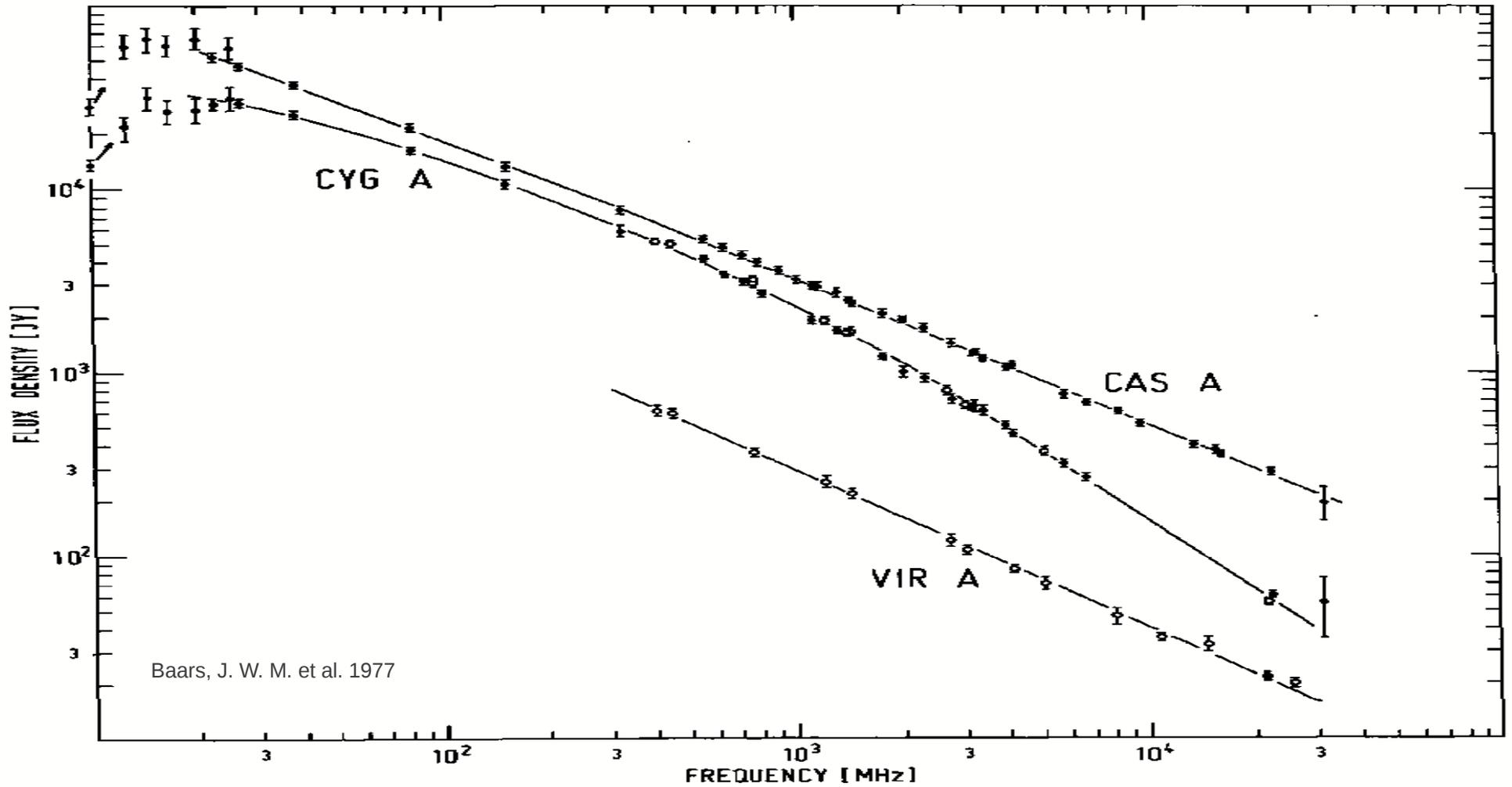
where the spectral index depends, for non-relativistic shocks, only on the compression ratio of the shock

 		 
Urto	Velocità	>
n.	nel lab.	
0	0	
1	+2v	
2	-6v	
3	+14v	
....		

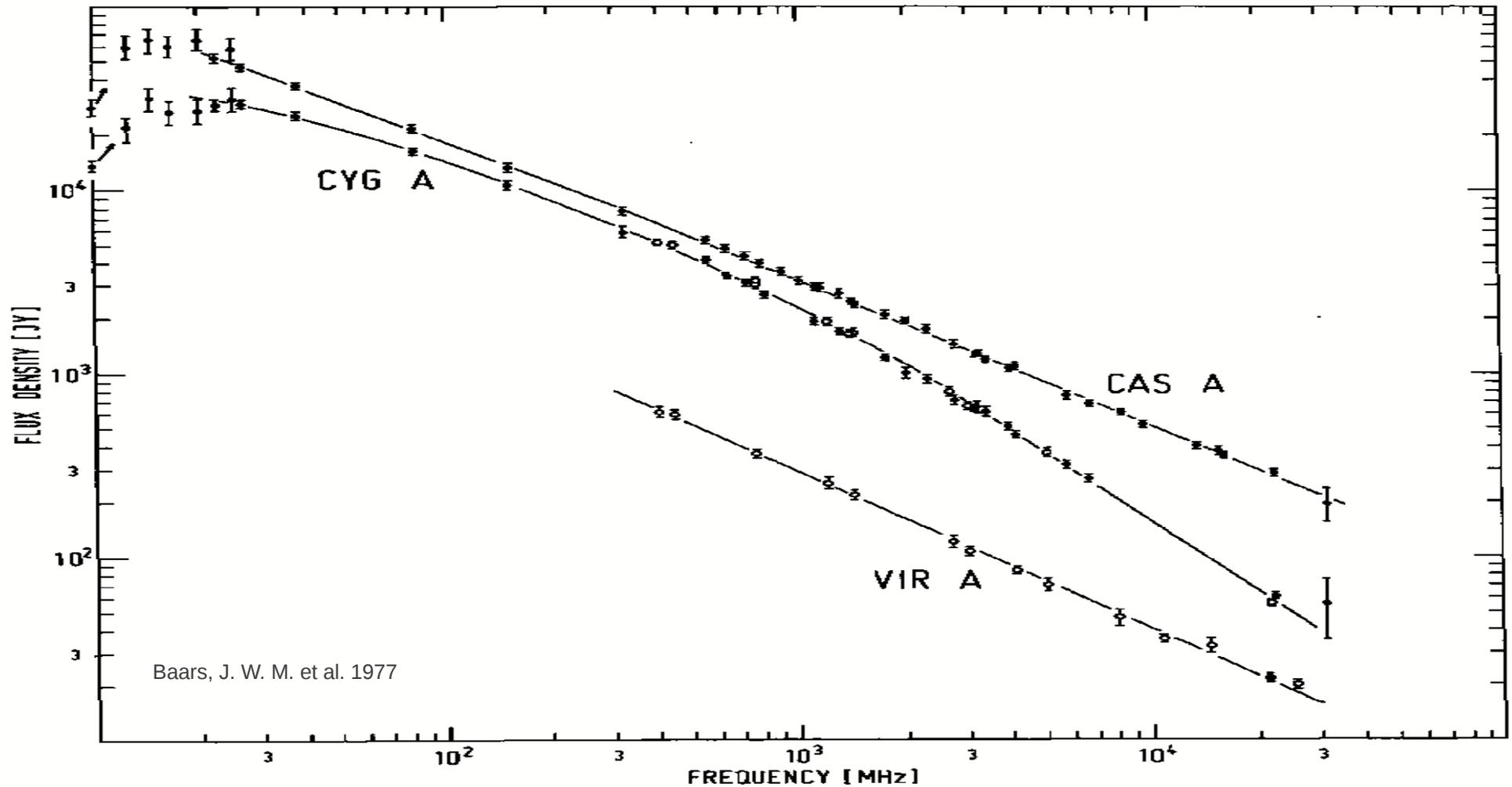
SNR Cas A (2nd brightest radio source!) Age ~300 yr



SNR Cas A (2nd brightest radio source!) Age ~300 yr



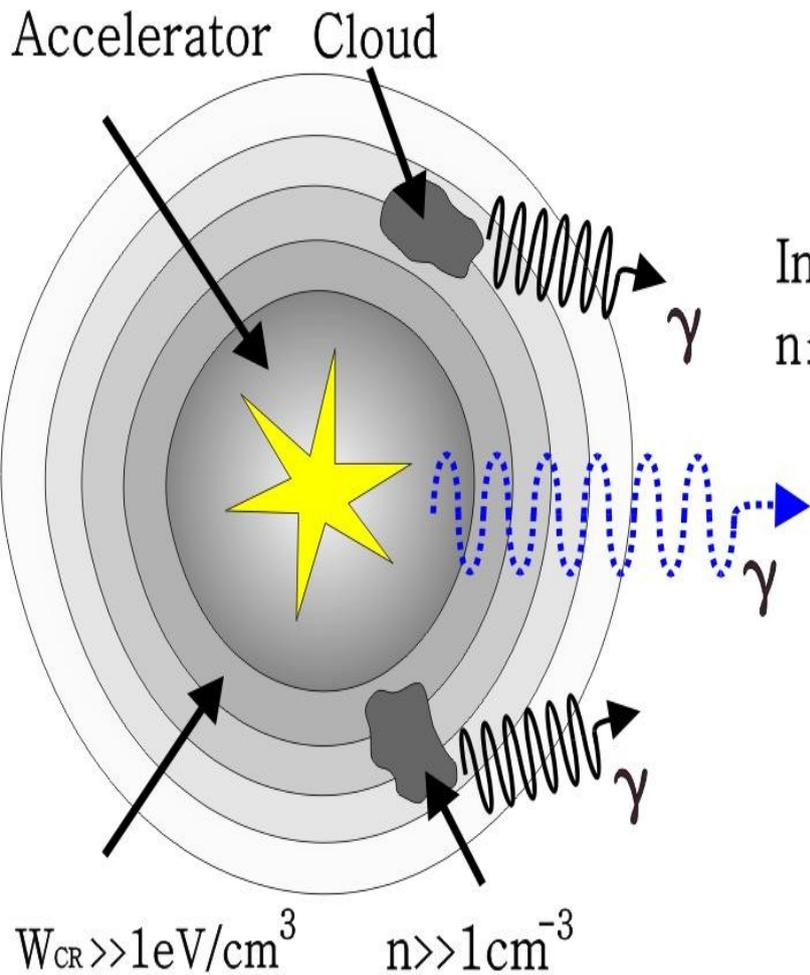
SNR Cas A (2nd brightest radio source!) Age ~300 yr



Featurless pl spectrum --> Non-thermal emission
--> Sync. emission from ultrarel. electrons

$$F \sim E^{-p} \quad p = 1-5 - 2.5 \quad E_{\max} \sim 1-10 \text{ TeV}$$

SNRs in gamma-rays



Neutral π decay :
protons + ISM nuclei $\rightarrow \pi^0$
 $\rightarrow \gamma$ rays

But also

Inverse Compton :
electrons + ISRF photons
 $\rightarrow \gamma$ rays

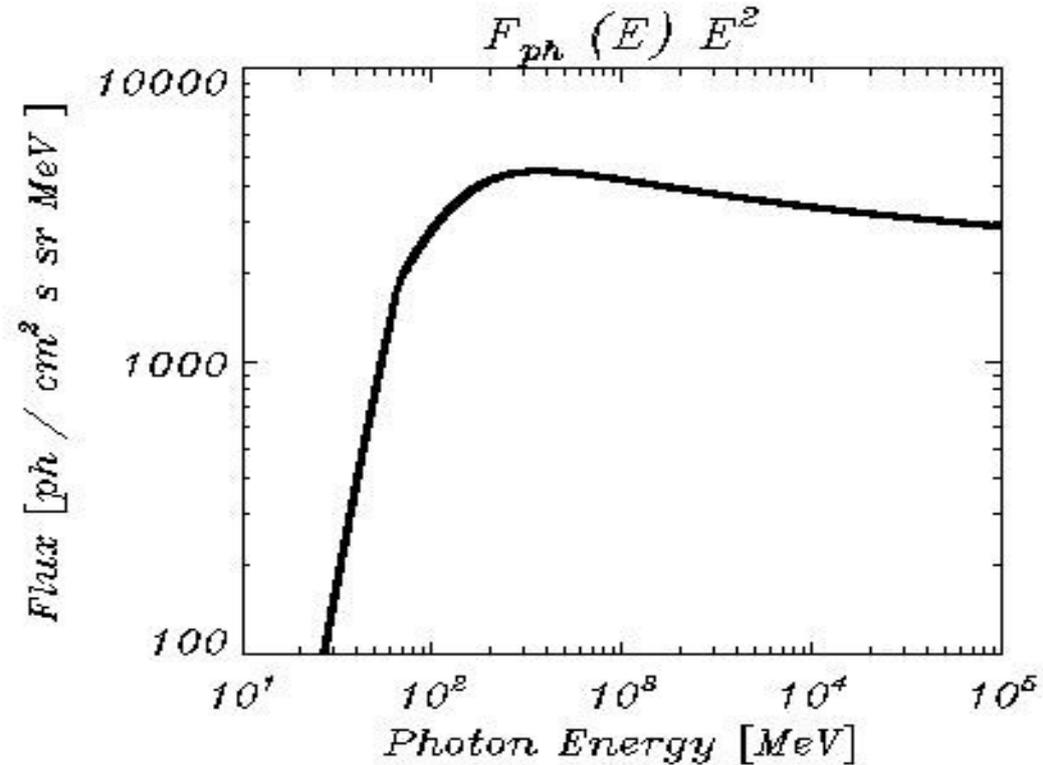
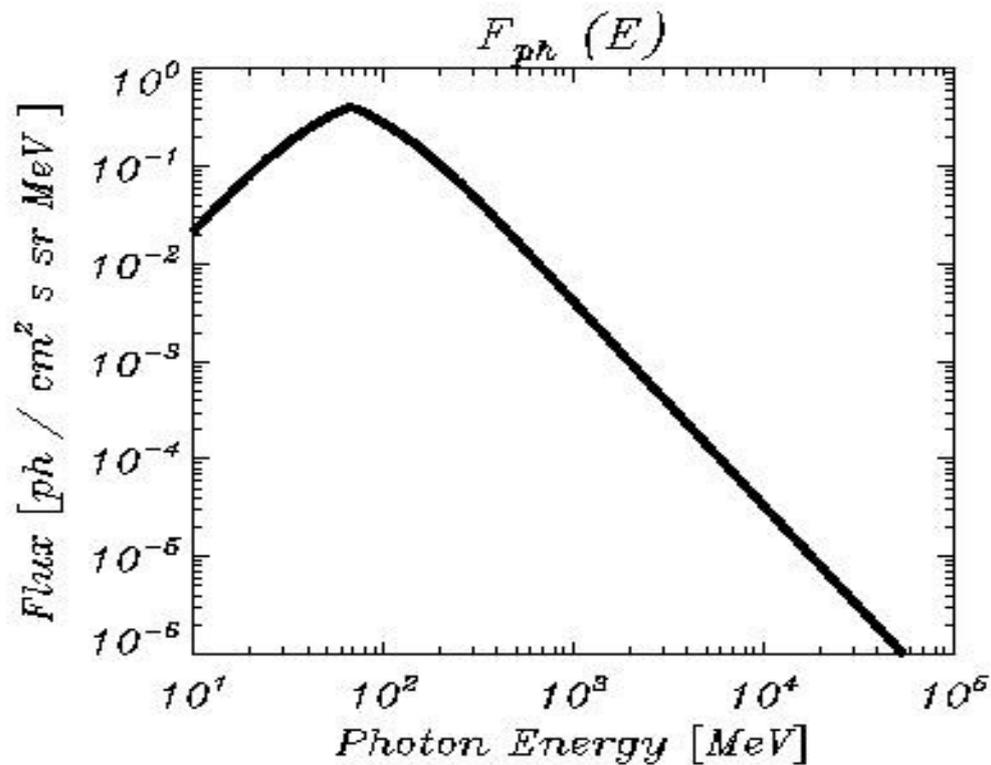
Electron Bremsstrahlung:
electrons + ISM nuclei
 $\rightarrow \gamma$ rays

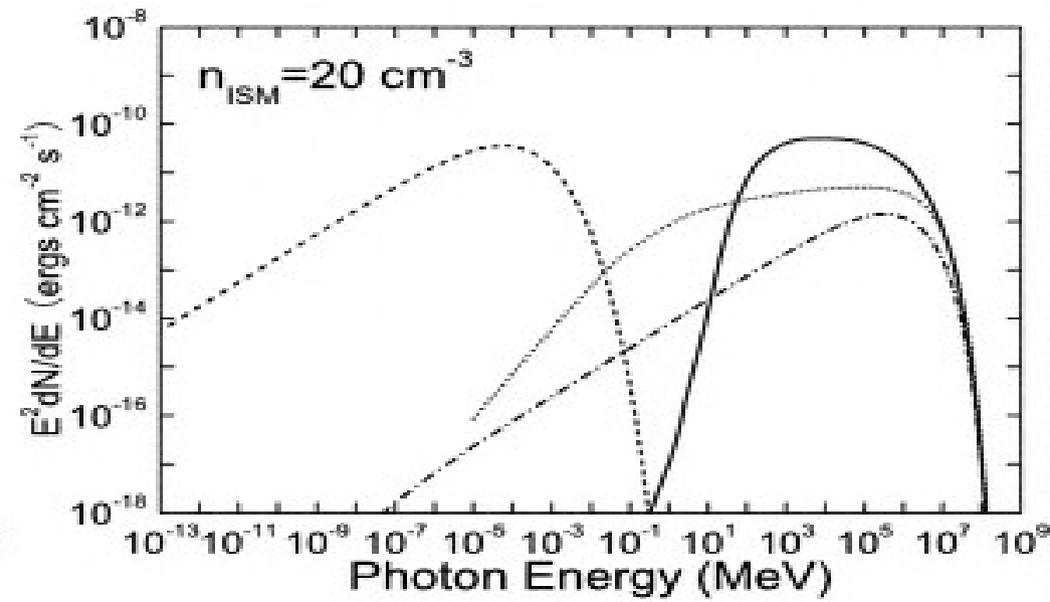
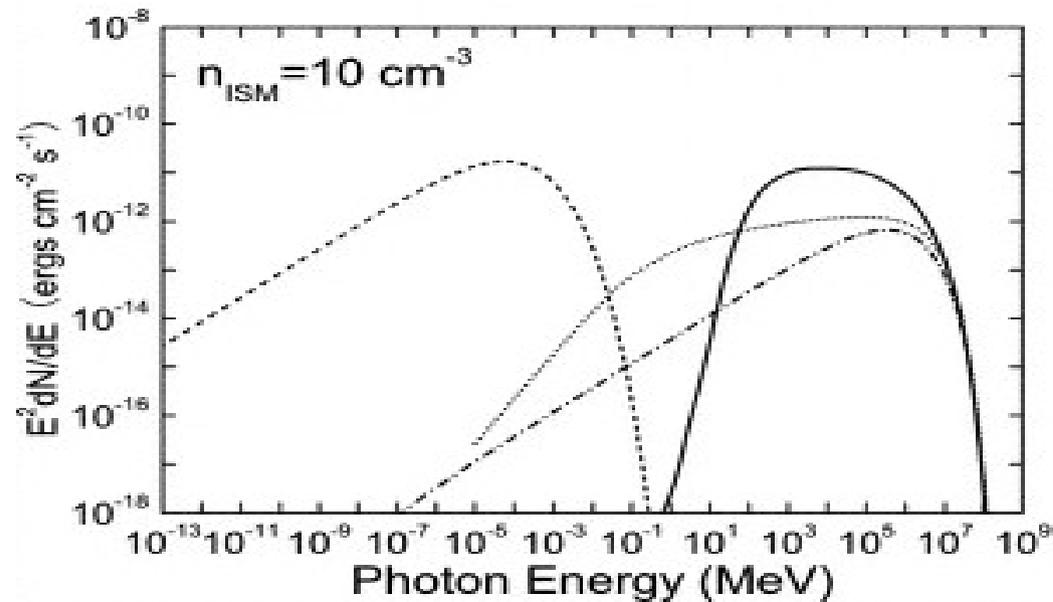
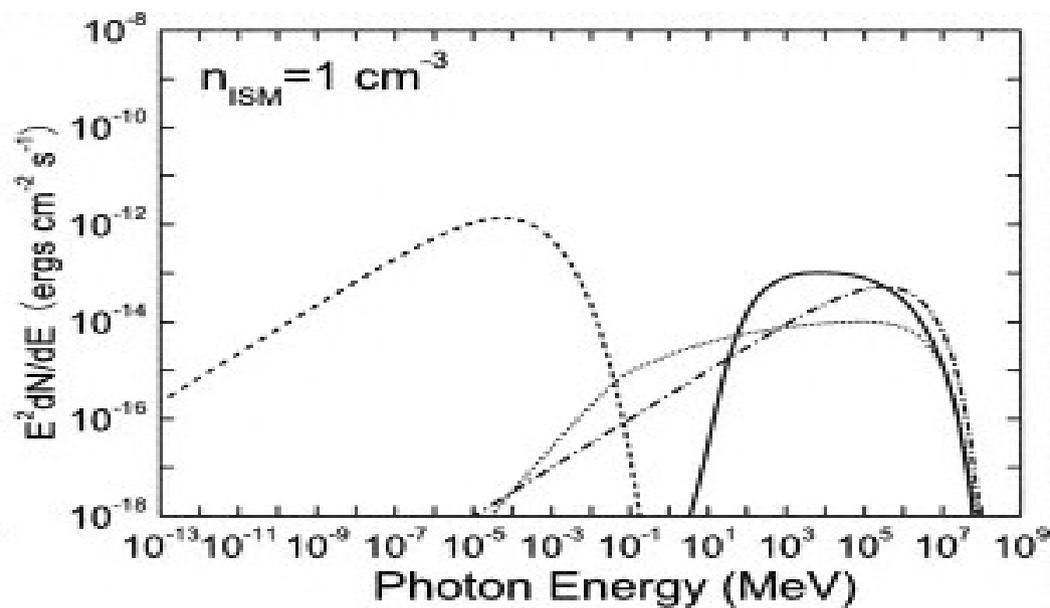
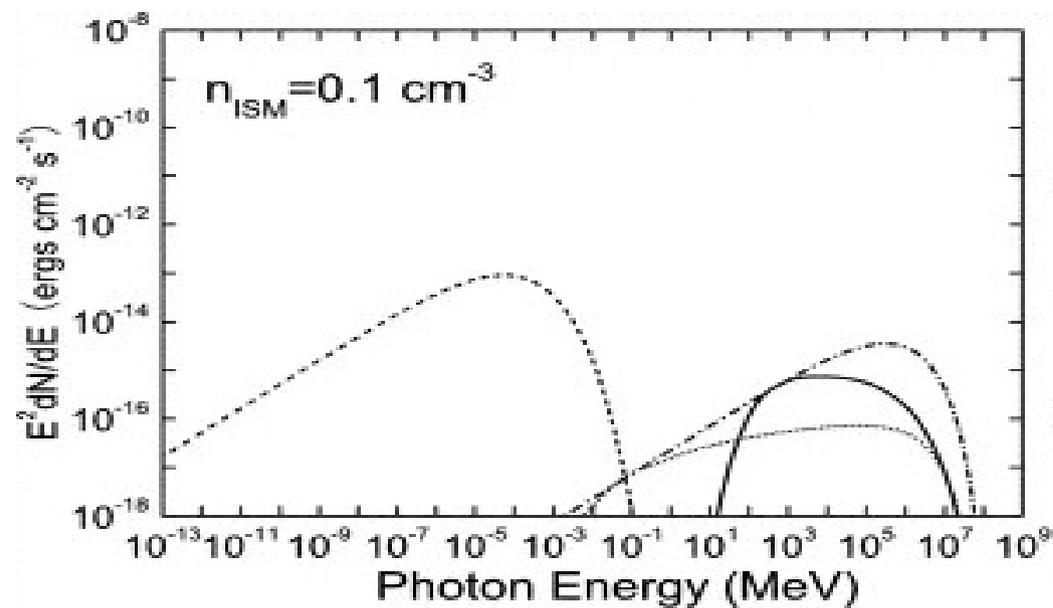
What about protons ?

protons + ISM nuclei $\rightarrow \pi^0 \rightarrow \gamma$ rays

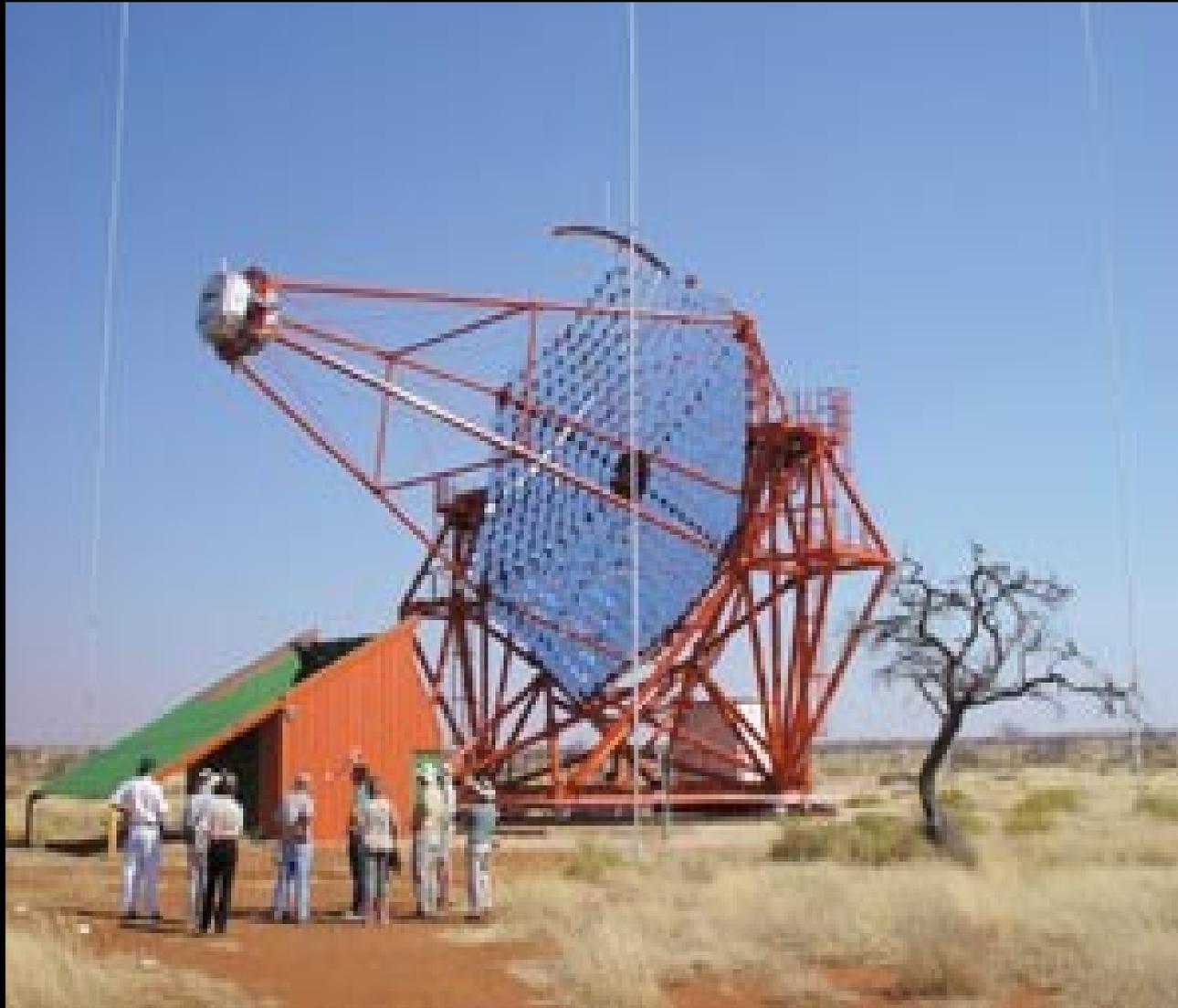
Emission in the hard gamma ray band

$F_{\text{gamma}} \sim f_p * n$ (density of ISM nuclei)





New generation of Cherenkov telescopes (HESS, MAGIC, VERITAS)

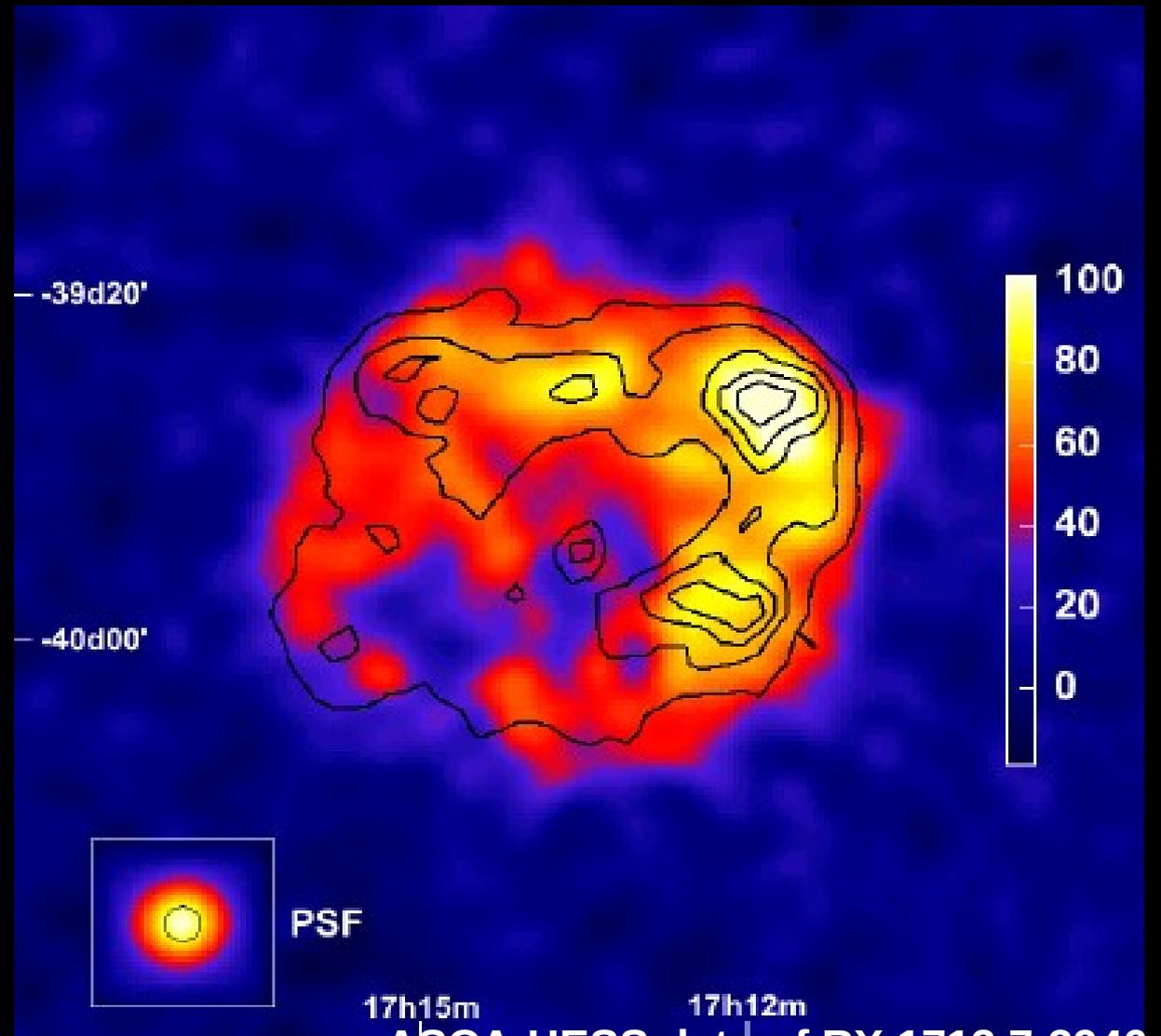


Energy Range : 100 GeV – 10 TeV

RX J1713.7-3946

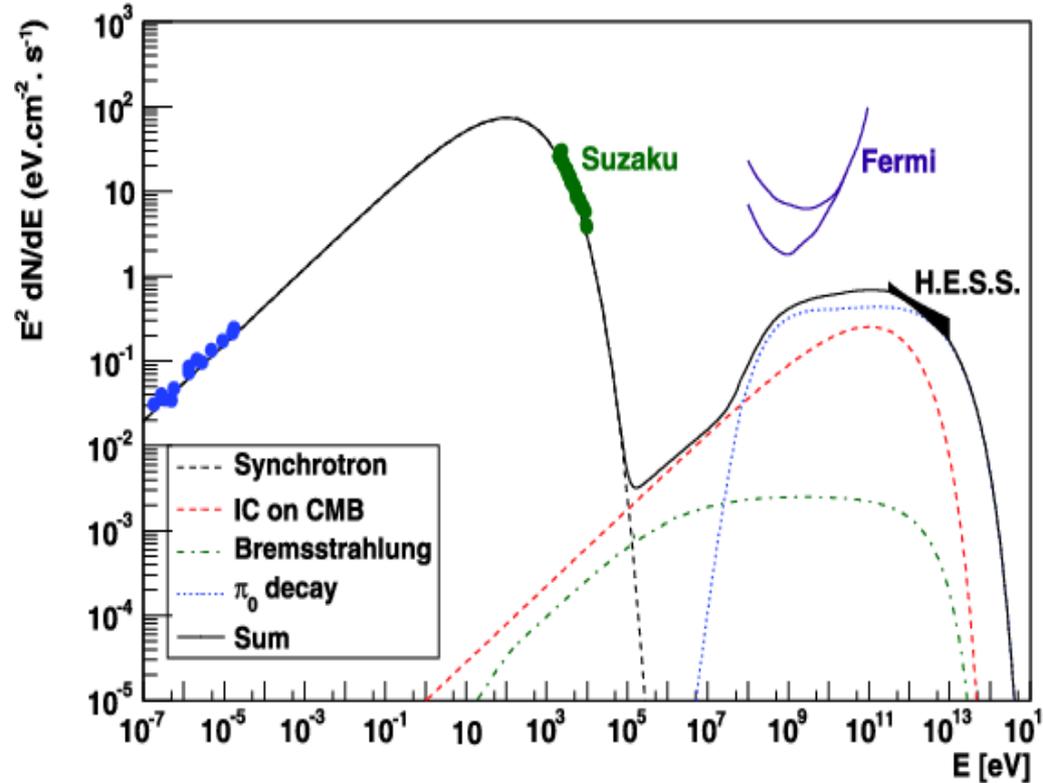
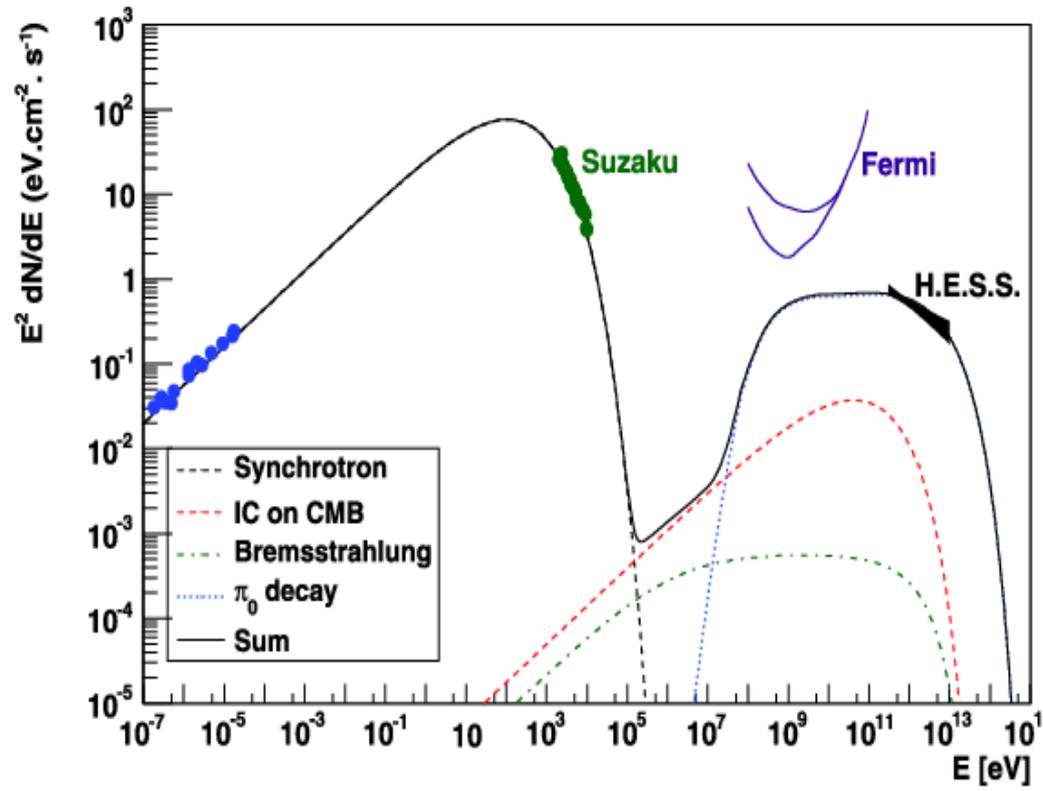
Distance : 1.2 kpc
Age : 2 kyrs
Size : 65'

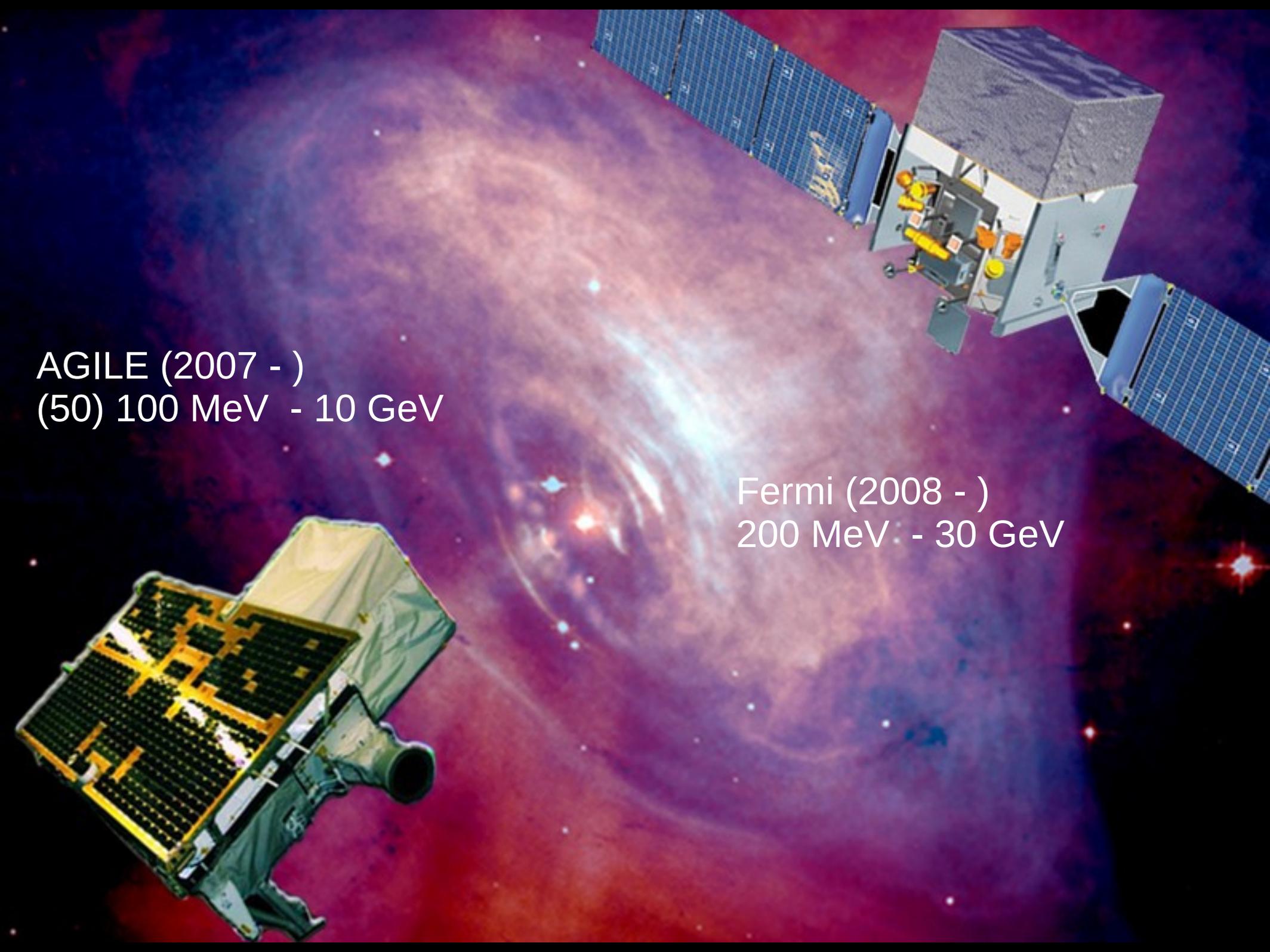
Good X-ray-TeV
correlation



ASCA-HESS data of RX 1713.7-3946
(Goumard et al. 2006)

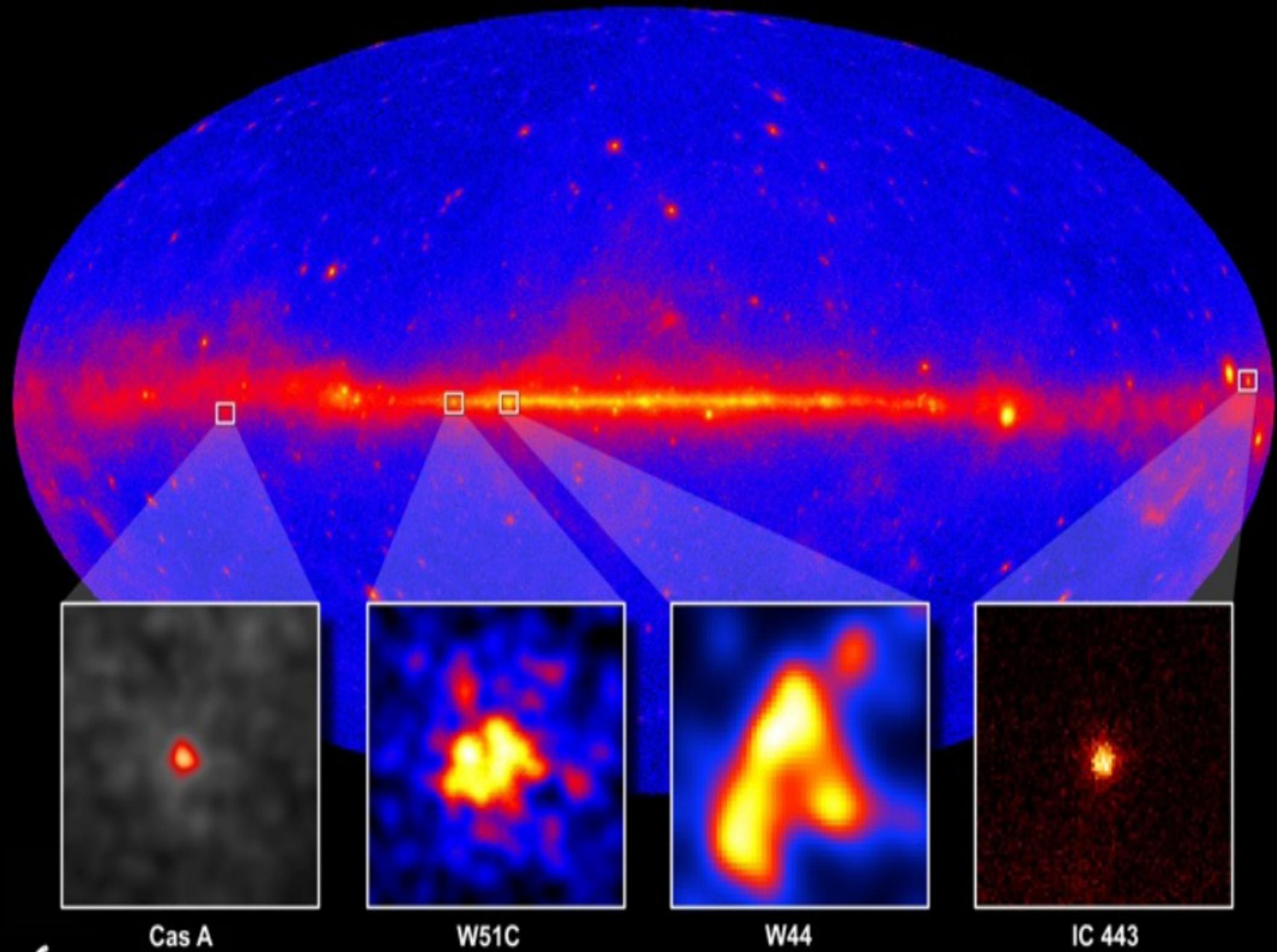
Non-thermal spectrum well fitted by both leptonic and hadronic models

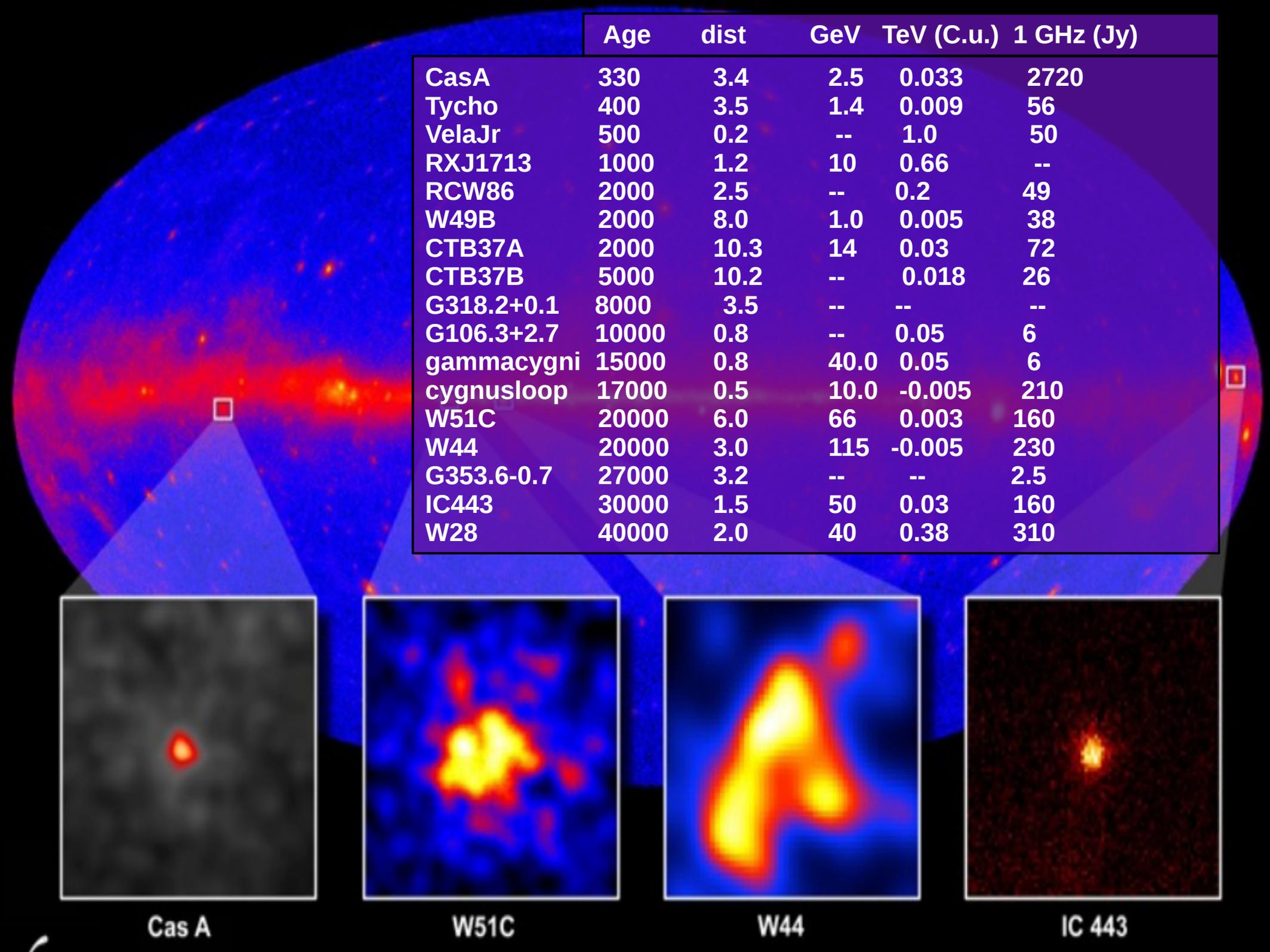


The image shows two gamma-ray satellites, AGILE and Fermi, in space. The AGILE satellite is in the lower-left, and the Fermi satellite is in the upper-right. The background is a colorful, swirling nebula. The Fermi satellite has a large, grey, rectangular instrument package and two large blue solar panel arrays. The AGILE satellite has a green and yellow instrument package and a large green solar panel array.

AGILE (2007 -)
(50) 100 MeV - 10 GeV

Fermi (2008 -)
200 MeV - 30 GeV





Cas A

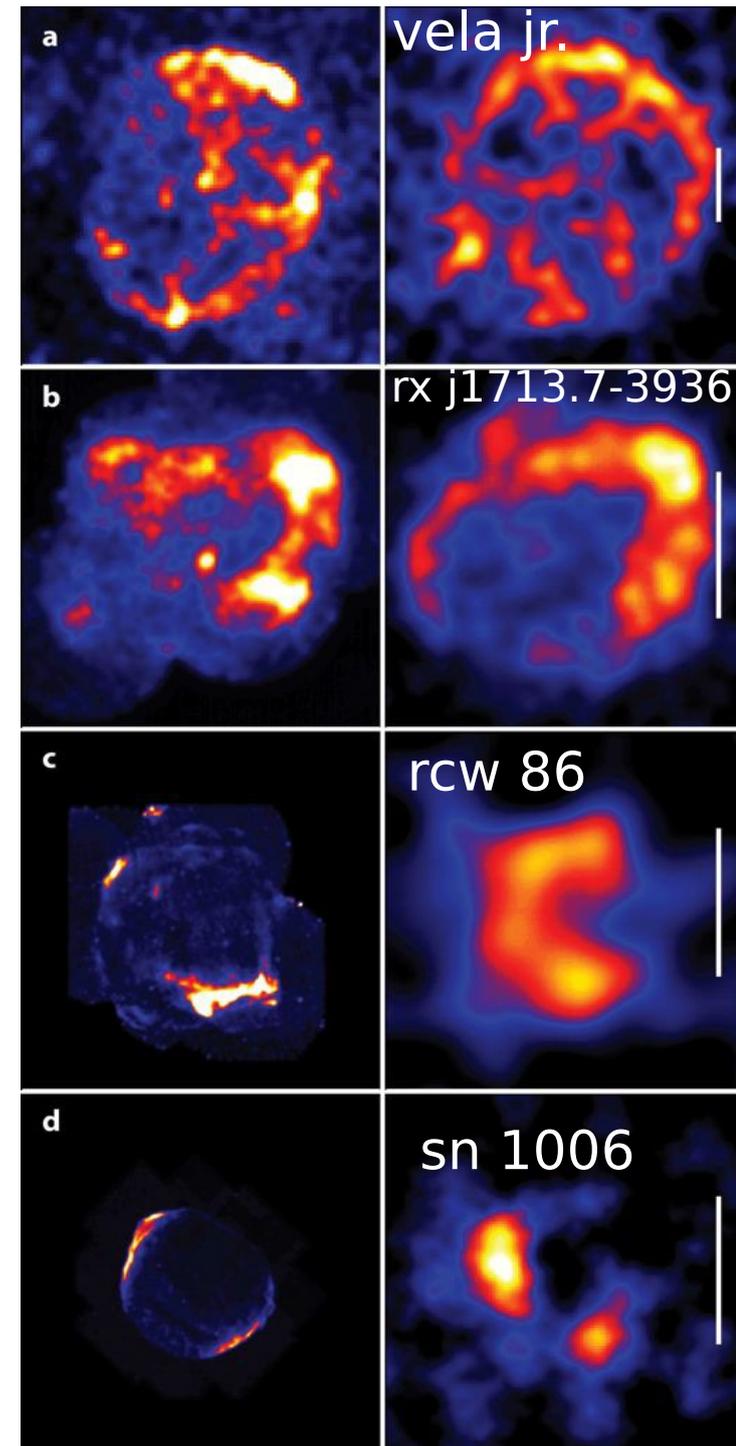
W51C

W44

IC 443

2 classes of gamma-rays SNRs

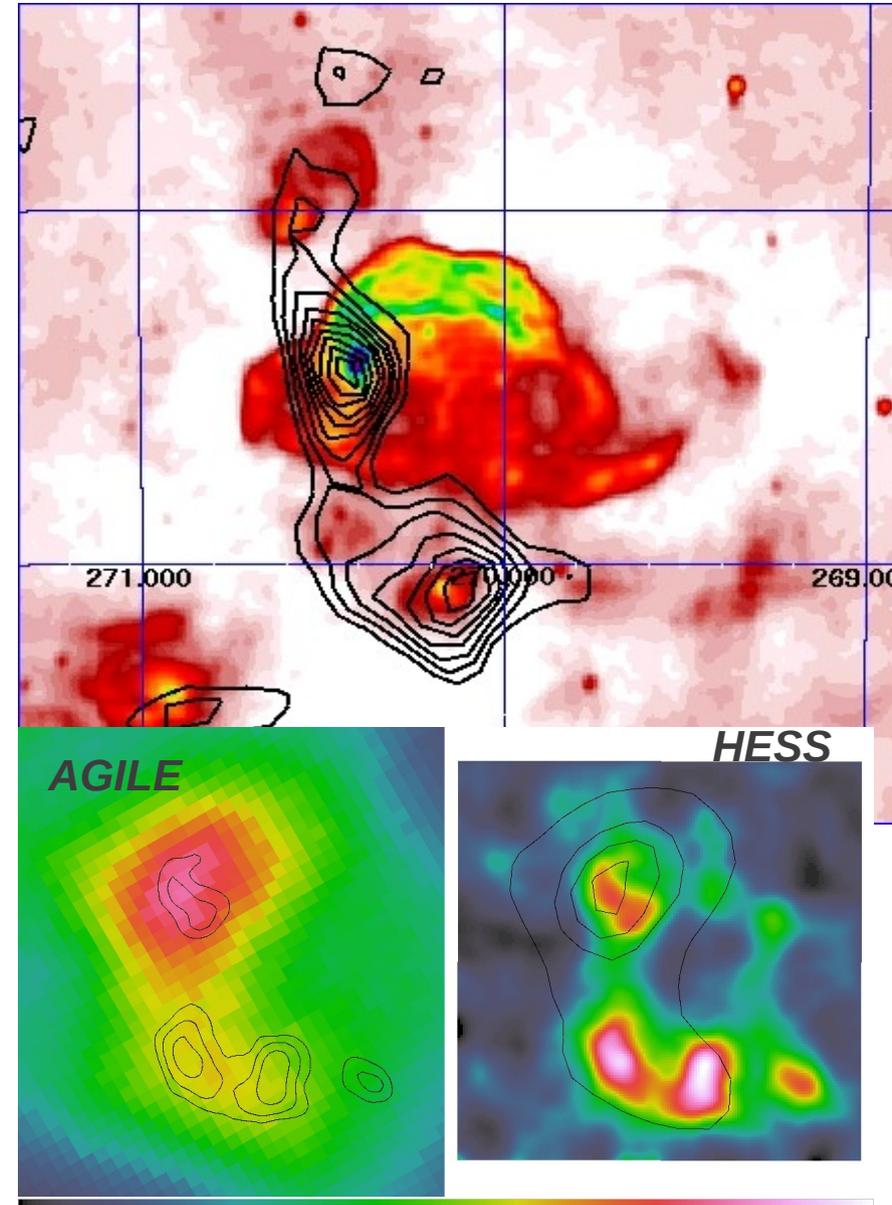
1) the young SNRs ($10^2 - 10^3$ yrs) are shell-like object, expanding in a relatively low density medium, with gamma emission morphology typically very nicely correlated with the radio (and often X) shell



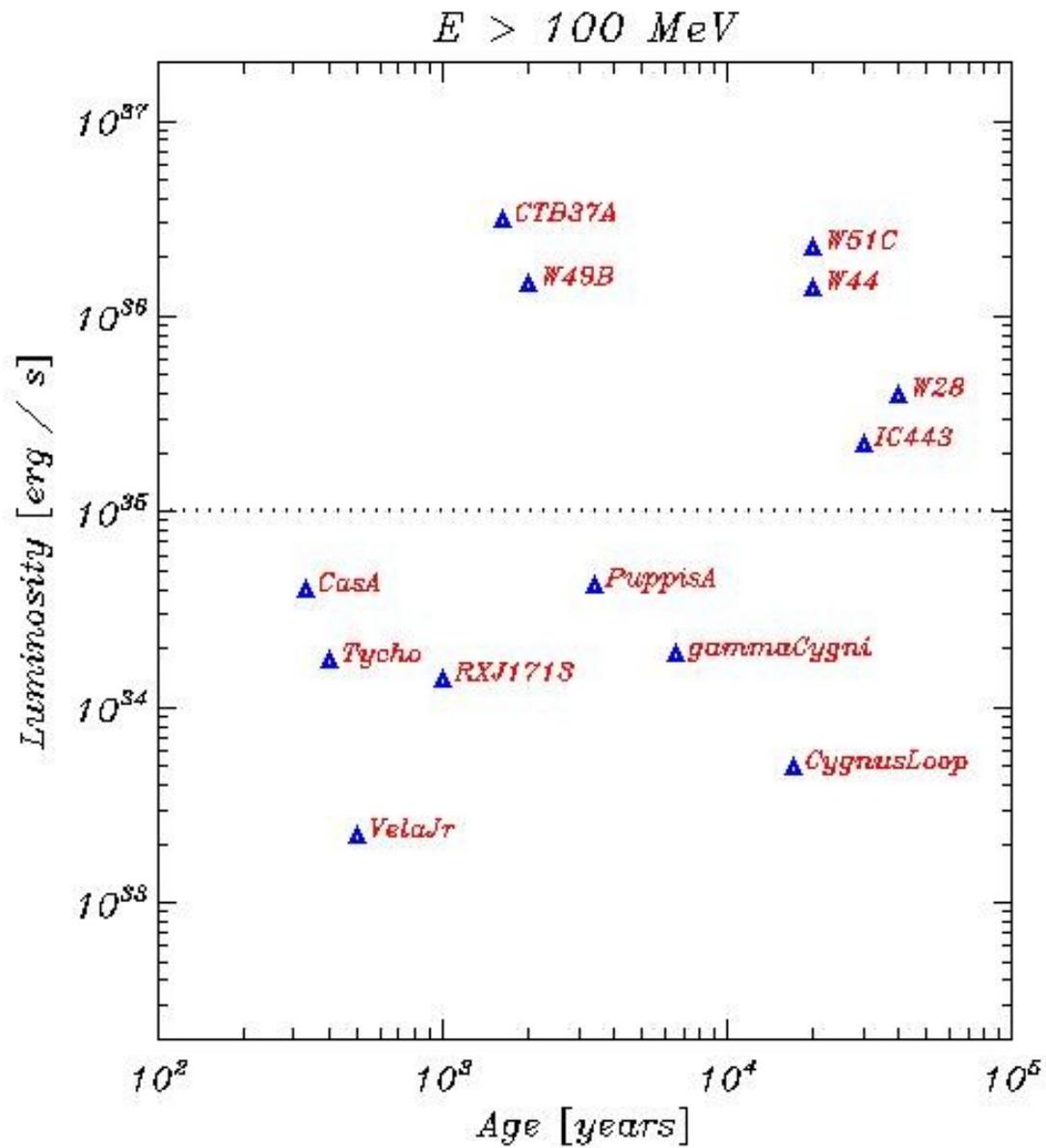
2 classes of gamma-rays SNRs

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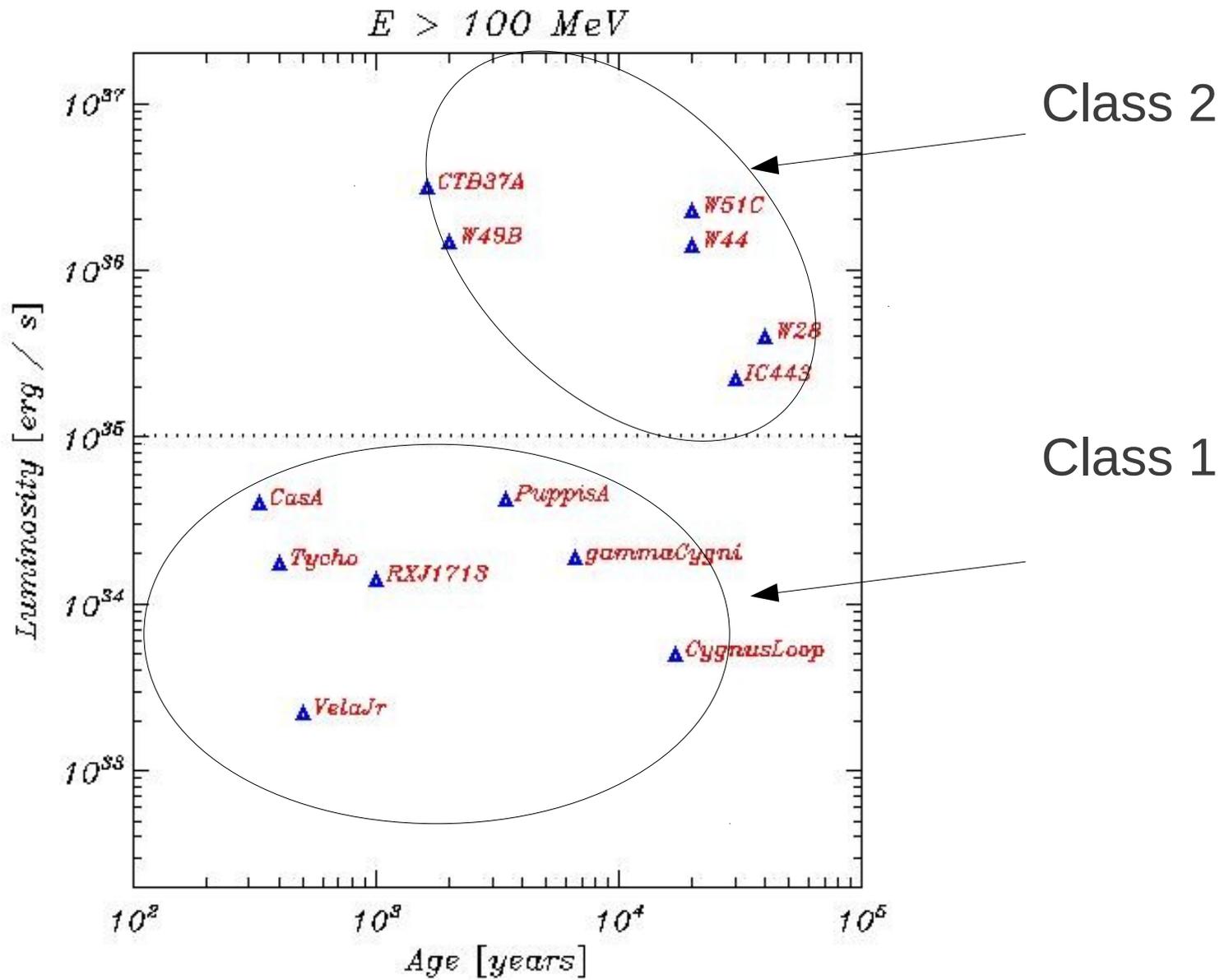
2) the middle-aged SNRs ($10^3 - 10^4$ yrs) are mixed-morphology objects, interacting with giant molecular clouds and with a gamma morphology that correlates with M.C. better than with the radio shell.



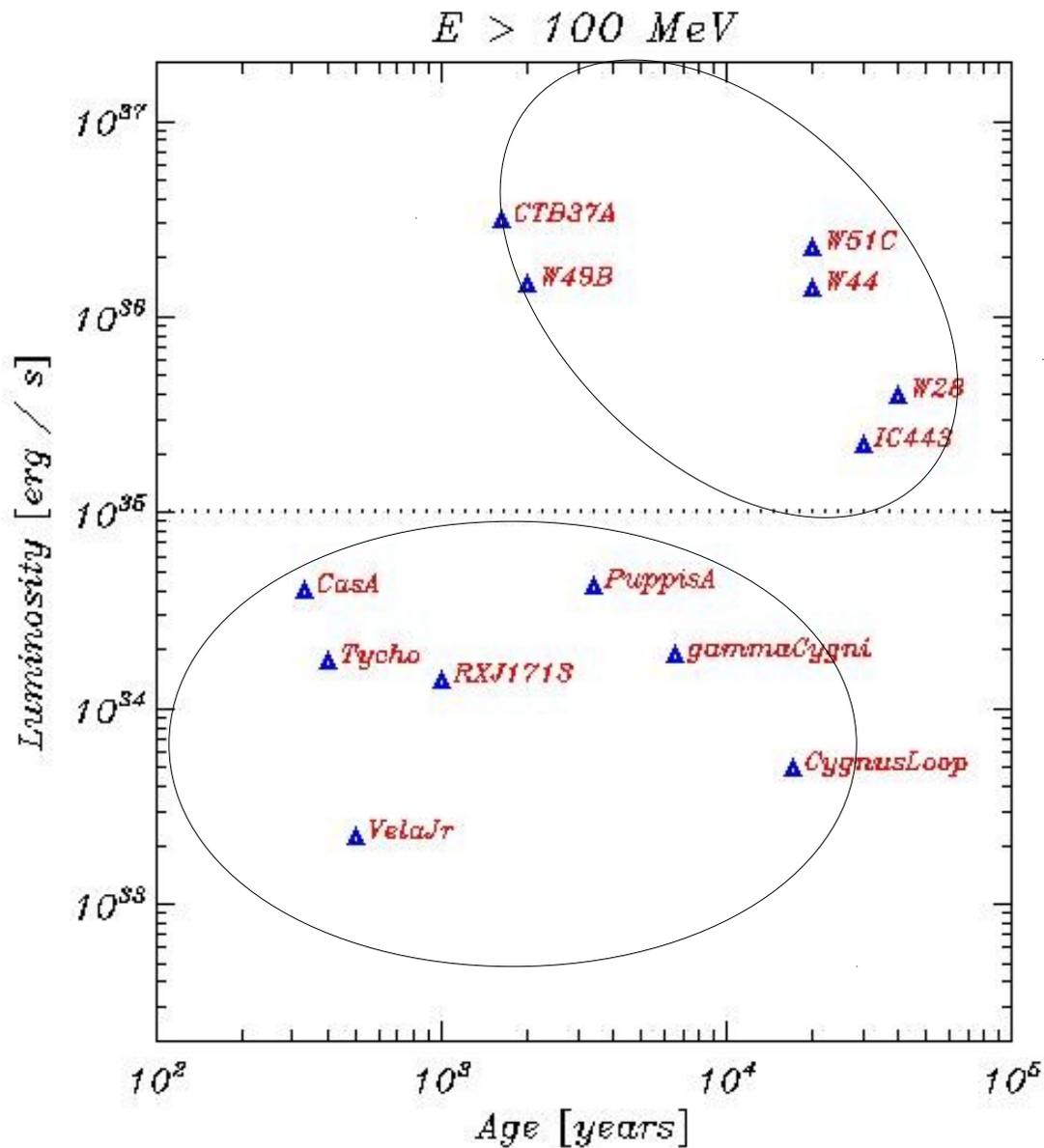
2 classes of gamma-rays SNRs



2 classes of gamma-rays SNRs



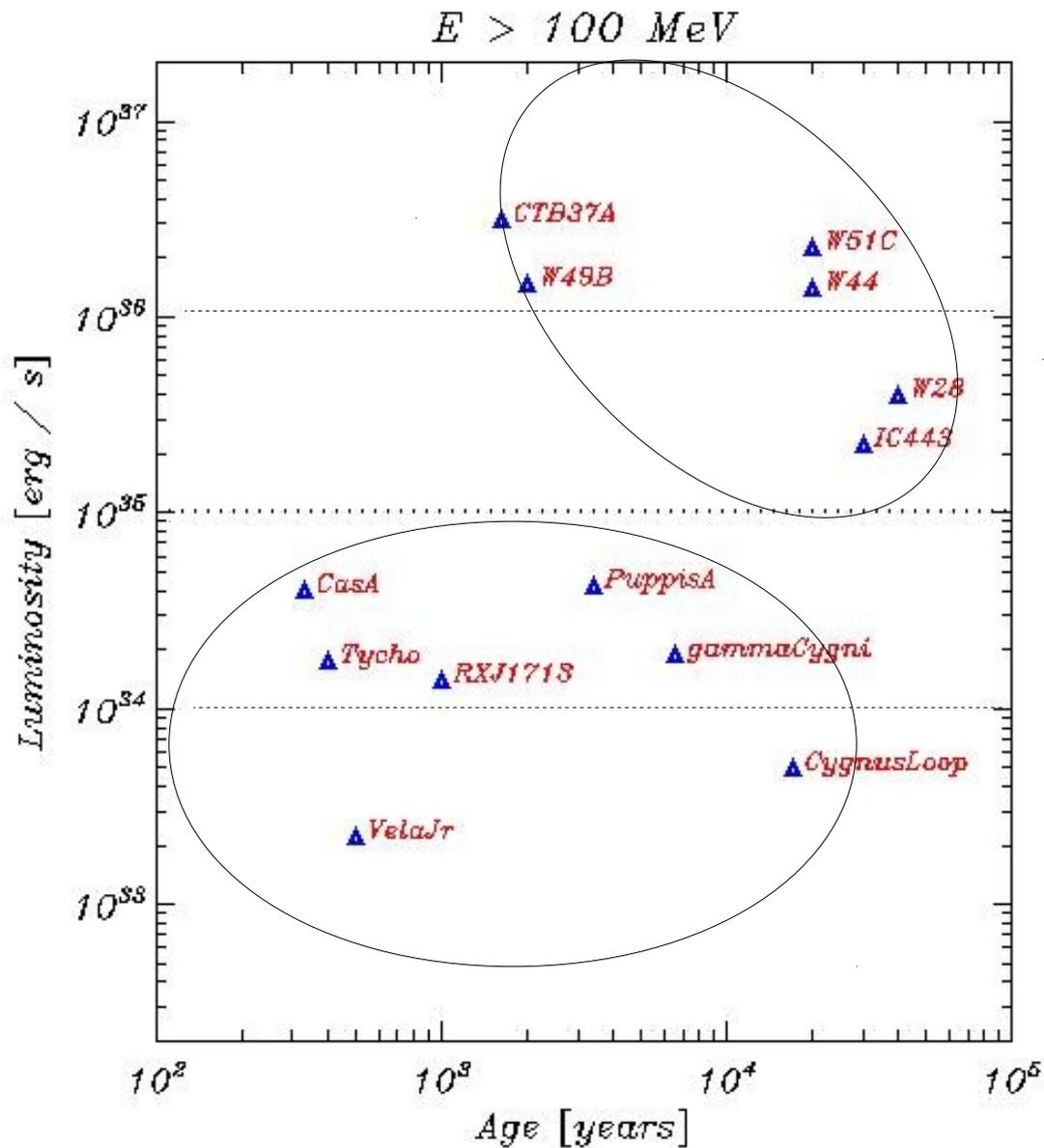
2 classes of gamma-rays SNRs



$$L_{\text{gamma}}(E) = \sigma(E) c N_p(E)$$

n

2 classes of gamma-rays SNRs



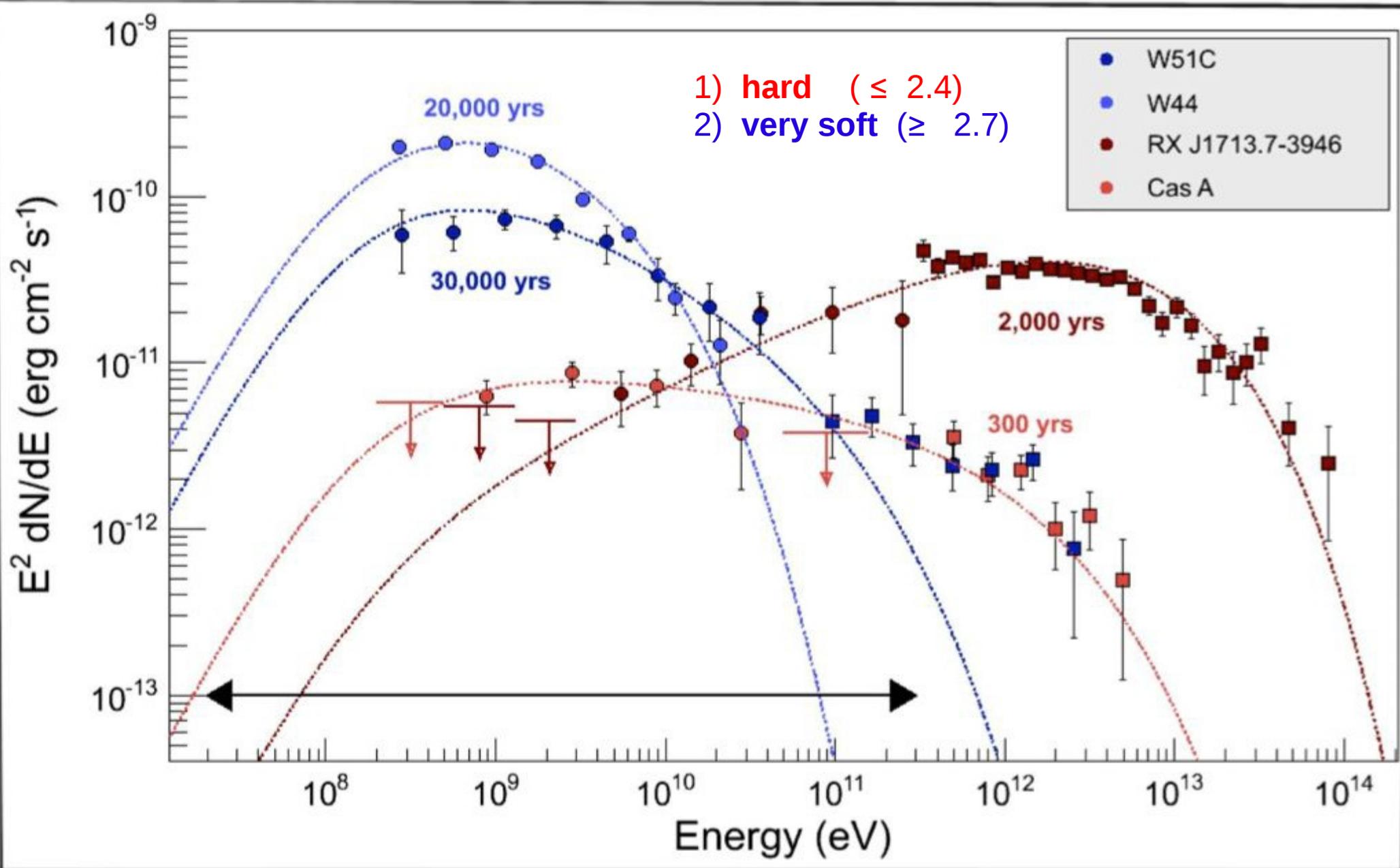
$$L_{\text{gamma}}(E) = \sigma(E) c N_p(E)$$

n

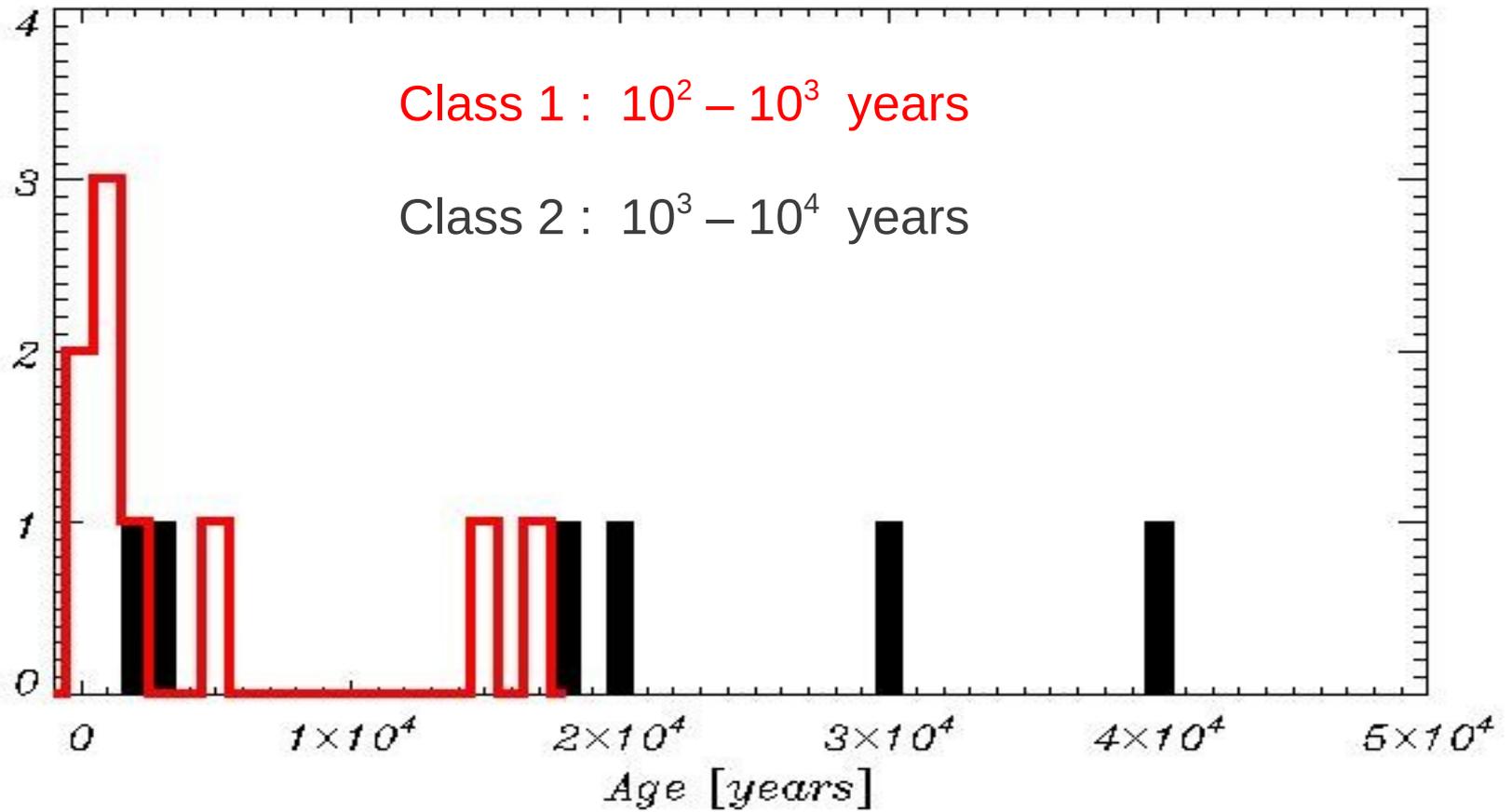
100 cm^{-3}

1 cm^{-3}

Different spectrum



Different Ages



Different frequency

Class 1 : Very common
(All the historical SNRs emit gamma-rays!)

Class 2 : Quite rare
(~ 1 %, only those SNRs close to a GMC)

Diffusion of CR in the ISM

$$\frac{dn(E, r)}{dt} = D(E)\nabla^2 n(E, r) - \frac{\partial}{\partial E} n(E, r)b(E) + Q(E, r)$$



Diff. in physical space



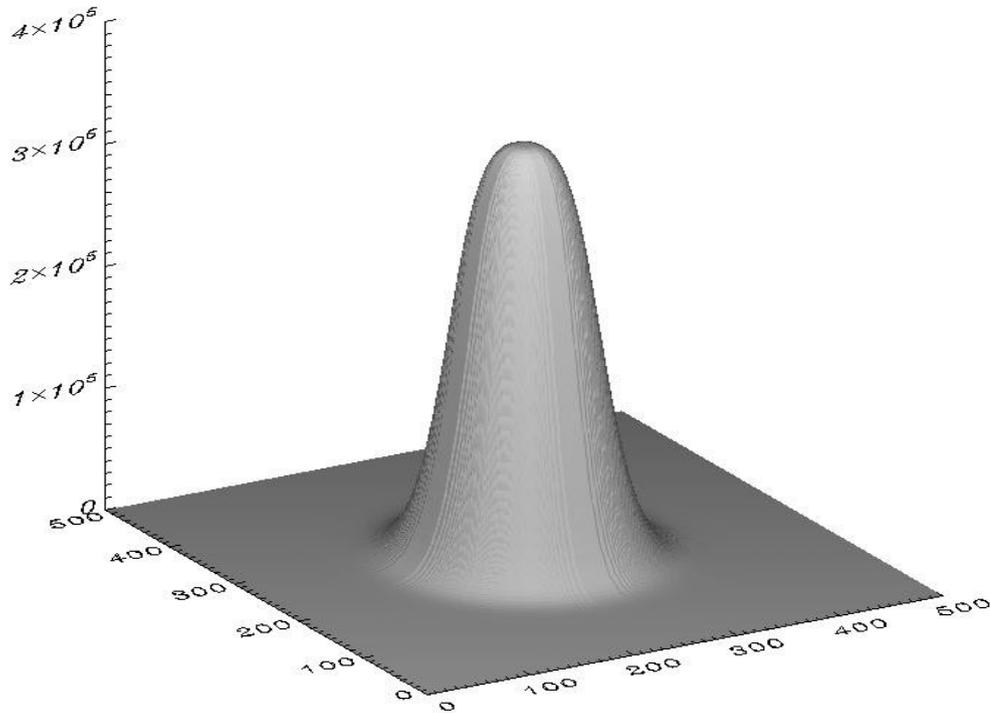
Energy losses



Source

Diffusion of CR in the ISM

$$\frac{dn(E, r)}{dt} = D(E)\nabla^2 n(E, r) - \frac{\partial}{\partial E} n(E, r)b(E) + Q(E, r)$$



For an impulsive source
and ignoring E losses

$$R_{diff}(E, t) = 2\sqrt{D(E)t}$$

see Aharonian & Atoyan, A&A, 309, 1996

Diffusion of CR in the ISM

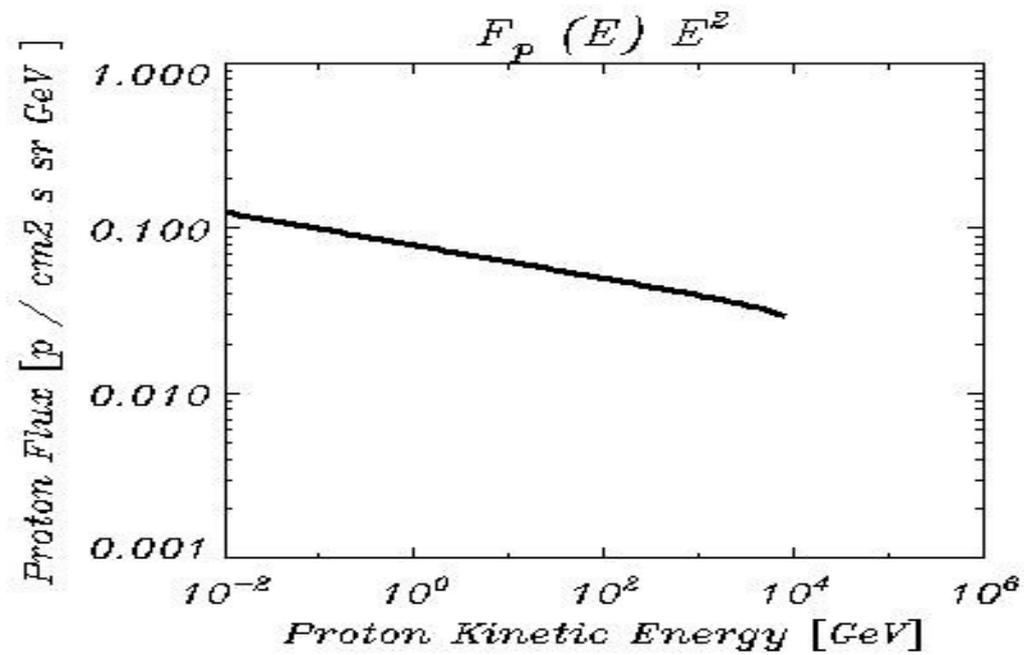
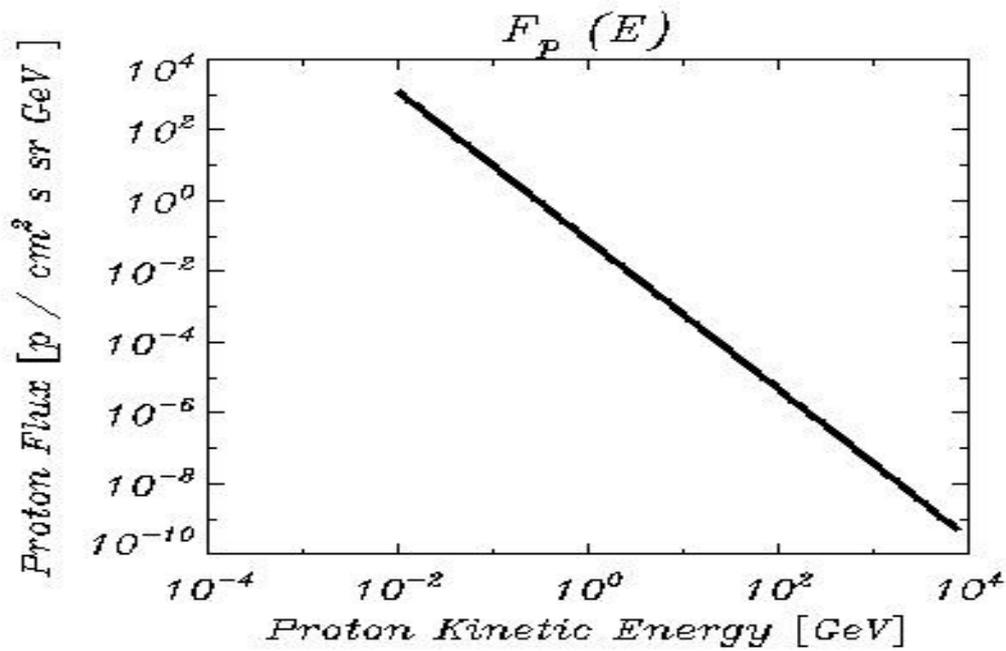
$$R_{diff}(E, t) = 2\sqrt{D(E)t}$$

$$D(E) = D_0 E^\delta \simeq 10^{26} \left(\frac{E}{10 \text{ GeV}} \right)^{0.5} \text{ cm}^2 \text{ s}$$

--> Slow !

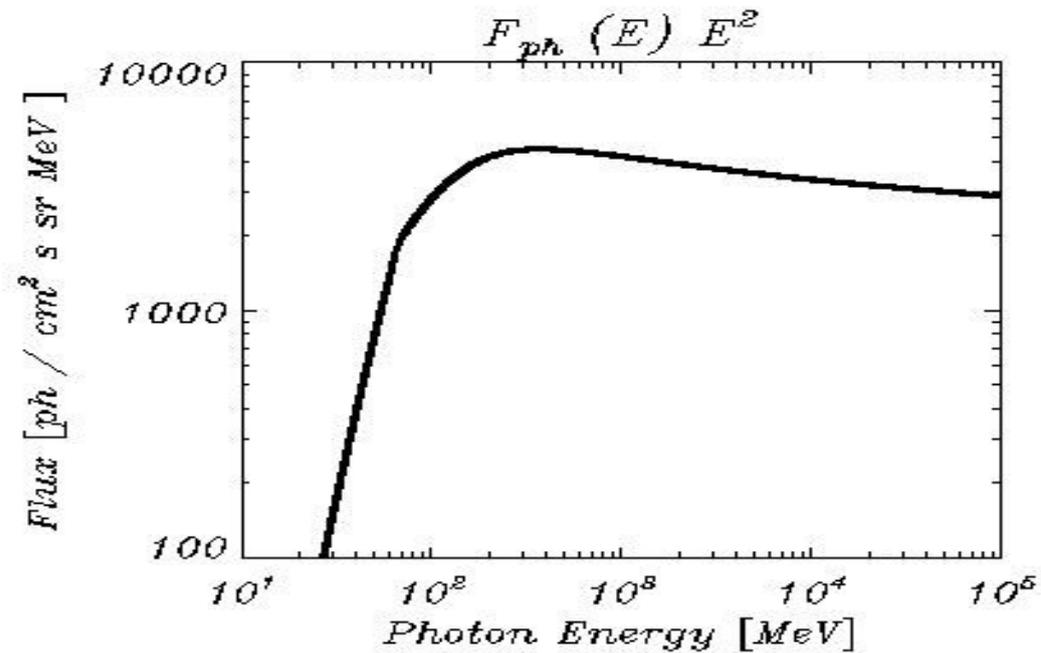
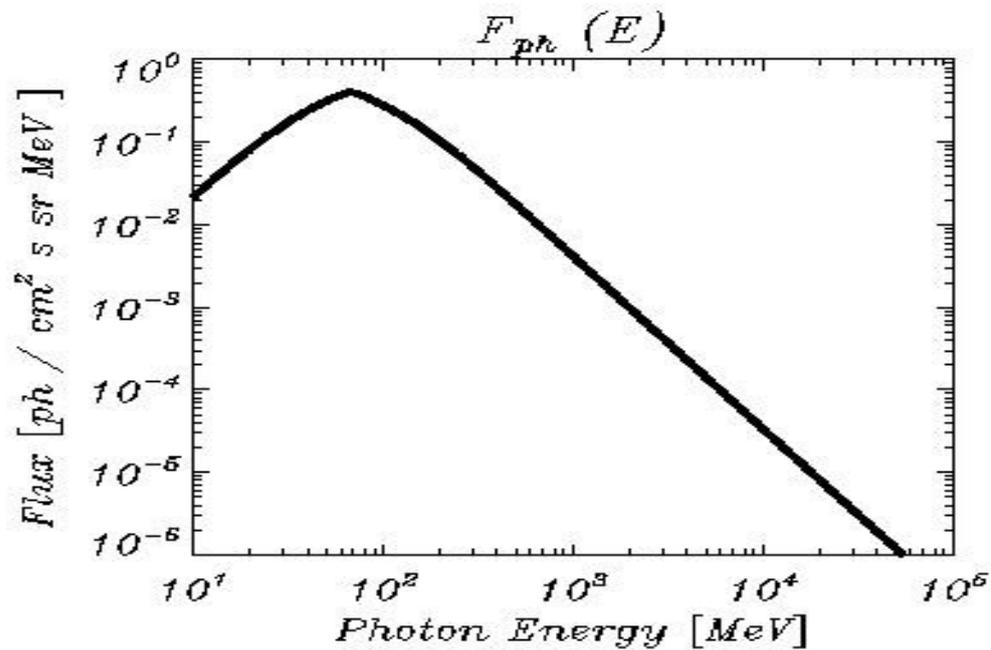
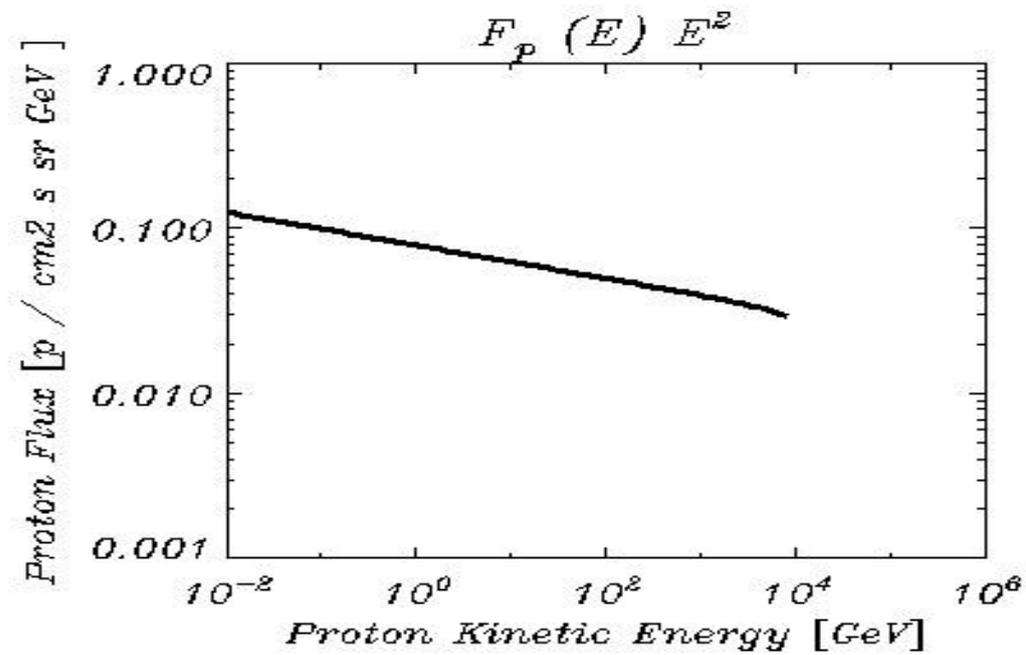
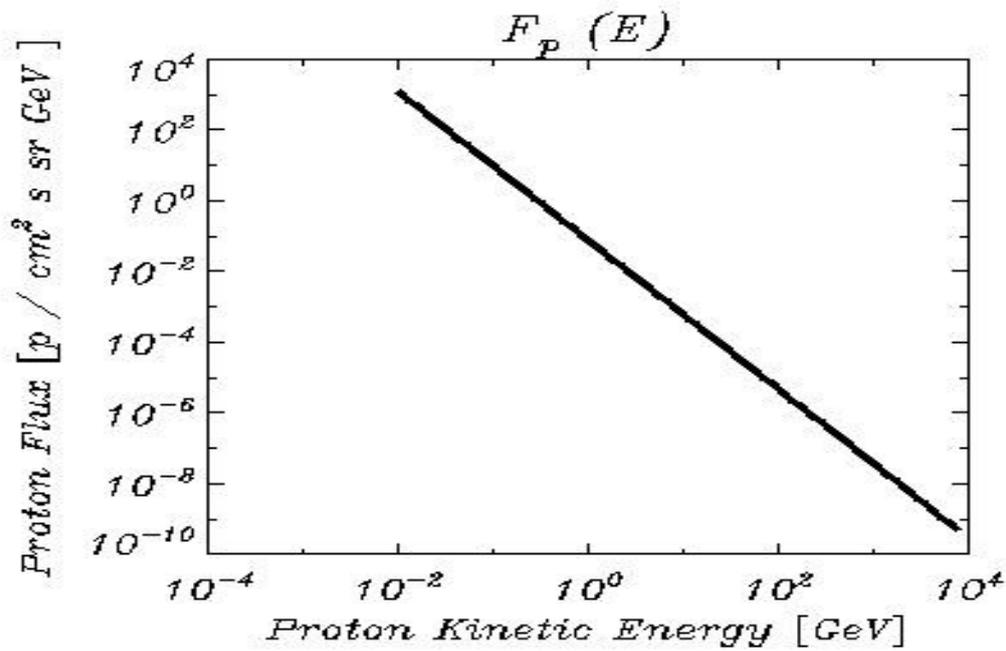
--> Faster diffusion for high energy CR

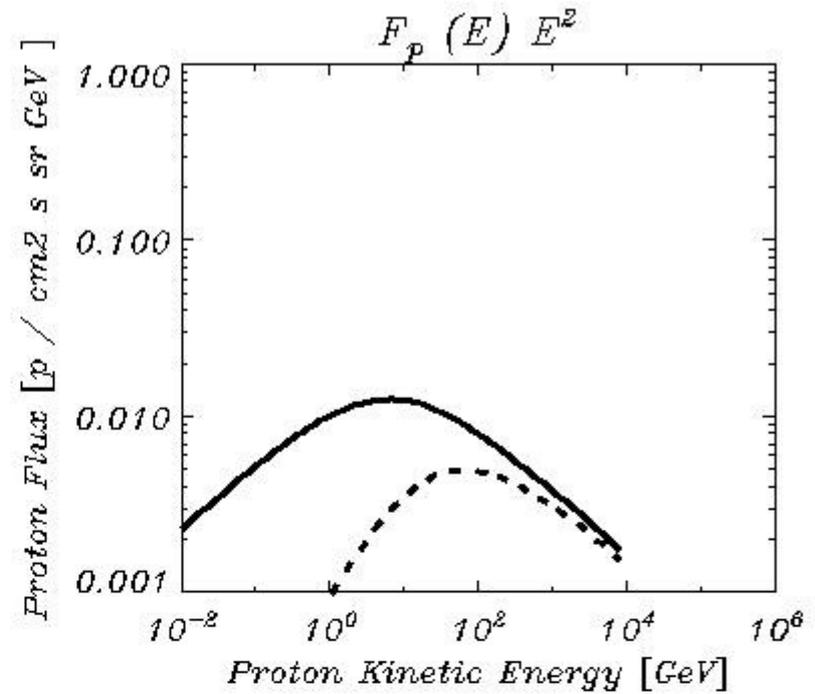
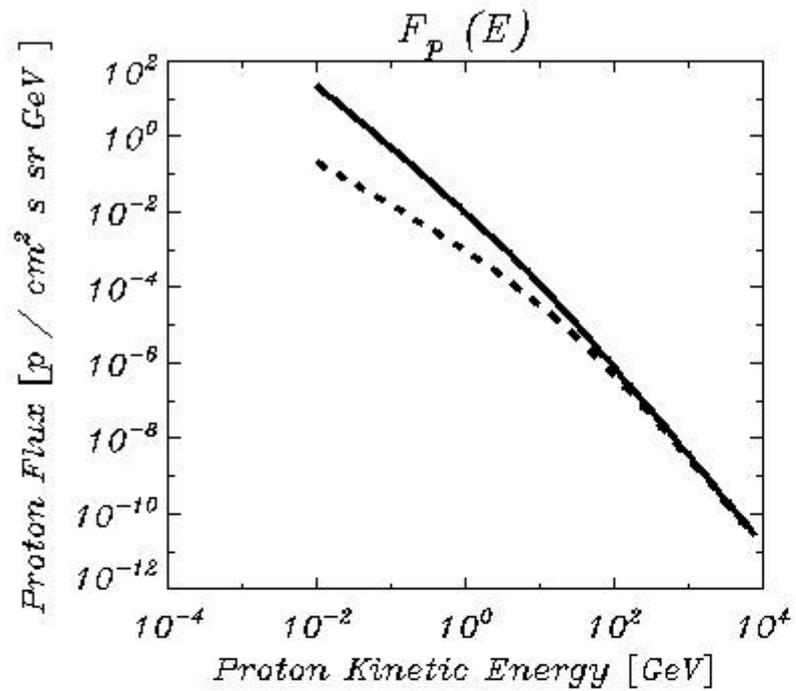
: protons + ISM nuclei $\rightarrow \pi^0 \rightarrow \gamma$ rays



$$F(E) = K E^{-p} \quad [\text{particles} / \text{cm}^2 \text{ s sr GeV}]$$

: protons + ISM nuclei $\rightarrow \pi^0 \rightarrow \gamma$ rays





For a given R and D :

$$E_{cut} = \left(\frac{R^2}{4tD_0} \right)^{\frac{1}{\delta}}$$

