Introduction to gamma-ray blazars Deirdre HORAN LLR / Ecole Polytechnique deirdre@llr.in2p3.fr



1. Active galactic nuclei 2. Blazars 3. Models of blazar emission 4. The Fermi blazars 5. The TeV blazars 6. The GeV-TeV Connection 7. The extra-galactic background light 8. Intergalactic magnetic fields 9. Characterising variability



# CHARACTERISTICS

- central nucleus outshines the rest of the galaxy
- high luminosity (normally)
- emission across entire spectrum ... radio to keV, MeV, TeV
  non-thermal
- strong variability
- radio-loud sources:
  - relativistic jets ... superluminal motion















### CHARACTERISTICS

- <5% of all AGN</li>
- jet points "at" us
- flat radio spectrum
- radio loud AGN
- ·large amplitude variability
- optical polarisation
- spectral energy distribution



#### Superluminal motion

FIRST PROPOSED BY REES IN 1966 (NATURE, 211, 468) YEARS BEFORE IT WAS FIRST OBSERVED WHEN VLBI TECHNIQUES WERE DEVELOPED



first pulse travels to observer in time D/c; the second, emitted time  $\Delta t$  later, has a shorter distance to travel: D- $\Delta y$ . Difference in arrival time is:  $\Delta t_{obs} = \left[\Delta t + \frac{(D - \Delta y)}{C}\right] - \left[\frac{D}{C}\right]$ 

#### Superluminal motion



Difference in arrival time is:  $\Delta t_{obs} = \left[ \Delta t + \frac{(D - \Delta y)}{C} \right] - \left| \frac{D}{C} \right|$ substitute for  $\Delta y$  & rearrange:  $\Delta t_{ab} = \Delta t (1 - \beta \cos \theta)$ measured transverse velocity, vobs  $v_{obs} = \frac{\Delta x}{\Delta t_{obs}} = \beta c \sin \theta$  $(1 - \beta \cos \theta)$ therefore:  $\beta_{abs} = \frac{\beta sin\theta}{(1 - \beta cos\theta)}$ 

#### Superluminal motion

So, if the plasma velocity was 0.95c and the angle to the observer was 5deg, the apparent velocity would be 1.5c. The effect is maximised when  $\cos\theta=\beta$ 



#### Relativistic beaming

 Another relativistic effect occurs because the knots of plasma are moving at velocities close to that of light

• When an emitting plasma has a bulk relativistic motion relative to a fixed observer, its emission is beamed in the forward direction in the fixed frame

 The flux density is thus changed by relativistic time dilation so an observer sees much more intense emission than if the plasma were at rest

• The observed emission,  $S_{obs}$  is boosted in energy over that emitted in the rest frame, S

#### Definitions:

The Doppler factor is a measure of the strength of the beaming:  $\delta = [\Gamma (1 - \beta \cos \theta)]^{-1}$ 

The Lorenz factor:  $\Gamma = \frac{1}{1-6}$ 

where 
$$\beta = \frac{v}{c}$$

$$S_{obs} = S \left[ \Gamma \left( 1 - \beta \cos \theta \right) \right]^{-3}$$

If plasma velocity is 0.95c and 
$$\Theta$$
 is 5deg, the boosting factor will be ~198

#### Pair-production optical depth

\*Dondi & Ghisellini (1995) MNRAS, 273, 583 \*\* URRY & PADOVANI (1995) PASP, 107, 803

- High energy gamma rays collide with softer radiation to produce ete- pairs
- For gammas to escape from a source, the optical depth for this process  $\tau_e$  must be sufficiently low
- The cross section for this process is maximized for collisions between gamma rays of energy ...

 $X_{\gamma} = \frac{hv_{\gamma}}{mc^2}$  and target photons of energy ...  $X_{target} = \frac{1}{X_{\gamma}}$ 

• The optical depth is then defined as:  $\gamma_{r} = \frac{\gamma_{r}}{5} N X_{tarret} R$ 

... where N is the number of soft photons, R is the radius of the plasma "blob" (assumed shape) and  $\sigma_{T}$  is the Thompson-scattering cross section\*

• A useful parameter that can then be derived \*\* is the compactness of the source - it is a direct measure of the importance of the pair-production process  $\mathcal{L} = \frac{L}{R} \frac{\sigma_{T}}{m_{e}c^{3}}$ 

• The criterion for gammas to escape from a source is ...  $\mathcal{T}_{e^{\pm}} \sim \frac{\mathcal{L}}{\mathcal{L}_{e}} \ll 1$ 



#### SPECTRAL ENERGY DISTRIBUTION





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Lower energy emission due to synchrotron emission from relativistic e-s in the jet Two fundamentally different approaches to explain the higher energy emission

•Leptonic & Hadronic



Lower energy emission due to synchrotron emission from relativistic e<sup>-</sup>s in the jet Two fundamentally different approaches to explain the higher energy emission

•Leptonic 🗧 Hadronic

- radiative output dominated by  $e^{-}/e^{+}$
- high-energy photons most likely the result of inverse Compton scattering by the same e<sup>-</sup>s that produced the synch
  - upscatter the low-energy photons responsible for first bump
    - synchrotron self-Compton
  - upscatter photons from the broad-line region, disc, torus ...
    - → external Compton



#### Lower energy emission due to synchrotron emission from relativistic e<sup>-</sup>s in the jet

BOETTCHER, M. (2012) FERMI & JANSKY, ASTRO-PH/1205.0539 BOETTCHER M. (2013), APJ (IN PRESS), ASTRO-PH/1304.0605 Two fundamentally different approaches to explain the higher energy emission

• Leptonic & Hadronic

- both e<sup>-</sup>/e<sup>+</sup> and p accelerated to ultrarelativistic energies
- p's exceed threshold for py photo-pion production on soft photon field in emission region
- high energy emission dominated by
  - proton synchrotron
  - $\pi^0$  decay products
  - synchrotron and Compton emission from
    - secondary products of charged pions
    - → external Compton

leptonic models provide good fits to many blazars



leptonic models provide good fits to many blazars

Modified Julian Date

• X-ray and gamma-ray emission often correlated - a fact naturally explained by SSC models



FORTSON ET AL. (2012) GAMMA 2012

Deirdre HORAN --- 2013 Fermi Summer School --- Lewes, Delaware

· leptonic models provide good fits to many blazars

#### • X-ray and gamma-ray emission often correlated - a fact naturally explained by SSC models

• in hadronic models, the cooling times are longer, which makes it more difficult to explain the rapid variability often seen in blazars

⇒ proton synchrotron can produce rapid variability with very high energy protons in extremely magnetised, compact regions HOLDER, J. (2012), ASTROPART. PHYS., 39, 61





"orphan" flare from TeV blazar, 1ES 1959+650

hadronic models have been invoked to explain this behaviour

E.G., BOETTCHER (2005), APJ, 621, 176 SAHU ET AL. (2013), PHYS. REV. D ... (IN PRESS) ASTRO-PH/1305.4985

#### H.E.S.S. COLLABORATION ET AL. (2013 - IN PRESS) DAVID SANCHEZ & PASCAL FORTIN

#### • VHE emission discovered by H.E.S.S. (ATel July 2010)

- Hard Fermi spectrum (2.1) and nearby (z=0.049)
  - ⇒ excellent candidate for

Tev emission

• X-ray observations taken simultaneously with VHE reveal onset of HE component at unusually low energies for a TeV BL Lac





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#### Difficult to model with a single zone homogeneous SSC model due to the unprecedented width of the HE component compared to the LE component









Table 5 Census of Sources AGN Type Entire 2LAC 2LAC Clean Sample<sup>a</sup> Low-lat Sample All 1017 886 104 FSRO 360 19 310 221 7 LSP 246 ISP 3 2 4 HSP 2 0 0 No classification 108 86 10 BL Lac 423 395 16 LSP 65 61 3 3 ISP 82 81 HSP 174 160 5 No classification 5 102 93 Blazar of unknown type 204 157 67 LSP 19 24 10 ISP 3 13 11 HSP 65 53 13 No classification 102 74 41 Other AGNs 30 24 2

Note. <sup>a</sup> Sources with single counterparts and without analysis flags. See Section 5 for the definitions of this sample.

ACKERMANN ET. AL 2011, APJ 743, 171



Figure 18. Theoretical contribution (W(E) of Equation (A3)) to TS per Ms and per log(E) interval as a function of energy for a power-law source over the average background at  $|b| > 10^{\circ}$ . The assumed photon spectral index is 2.2. The dashed line is for an isolated source. The full line includes approximately the effect of source confusion.

#### Abdo et. al 2010, Apjs 188, 405

\*8 misaligned blazars 4 NLSyls 10 AGN of other type 2 starbursts



"If the seed photon source for external Compton scattering is the broad line region (BLK), and the BLK strength is correlated with the power injected into electrons in the jet, one would expect that more luminous jets have stronger broad emission lines and greater Compton cooling, and thus a lower  $v_{sy}$ . As the power injected in electrons is reduced, the broad line luminosity decreases, there are fewer seed photons for Compton scattering, and consequently the peak synchrotron frequency moves to higher frequencies. This is also reflected in the lower luminosity of the Compton-scattered component relative to the synchrotron component as  $v_{sy}$  moves to higher frequencies." FINKE, J. (2012), FERMI-JANSKY (ASTRO-PH/1301.6081)



ACKERMANN ET. AL 2011, APJ 743, 171

#### The Hard Source List - a catalog above 10 GeV D. PANEQUE, FERMI SYMPOSIUM (2012)

• Work is underway to publish the 1st Fermi-LAT catalog of sources > 10 GeV

• Shape of the spectrum at > 10 GeV might not be well characterized if we use a single fit in the energy range 0.1 GeV - 100 GeV



which may radiate from the same/different location

- → Are sources more variable at HE than at LE ? or the other way around ?
- Understand better the population of sources emitting above 10 GeV
  - What are the sources dominating the highest LAT energies?

#### •IDENTIFY PROMISING CANDIDATES FOR IACTS ... NEW VHE DISCOVERIES



# The TeV blazars



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As of writing this talk, there are 145 sources in TeVCat - 57 Extragalactic



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